

Fragmentation and Monte Carlo generators at Belle II

Work in progress.

Ami Rostomyan
(for the Belle II collaboration)

Excited QCD 2017

7-13 May 2017

Sintra, Portugal

Belle II @ SuperKEKB

New facility to search for physics beyond the SM by studying B, D and τ decays and the major player for precise studies of fragmentation

Tsukuba, Japan

SuperKEKB – major upgrade of the KEKB B-factory at KEK



Belle II detector – upgraded Belle detector

Belle II @ SuperKEKB

New facility to search for physics beyond the SM by studying B, D and τ decays and the major player for precise studies of fragmentation

Tsukuba, Japan

SuperKEKB – major upgrade of the KEKB B-factory at KEK

- ➔ smaller interaction point
- ➔ increased currents
- ➔ collisions at Y(nS)
 - ➔ $E(e^+) = 4 \text{ GeV}$, $E(e^-) = 7 \text{ GeV}$

First collisions in 2016

Belle II detector – upgraded Belle detector



Belle II @ SuperKEKB

New facility to search for physics beyond the SM by studying B, D and τ decays and the major player for precise studies of fragmentation

Tsukuba, Japan

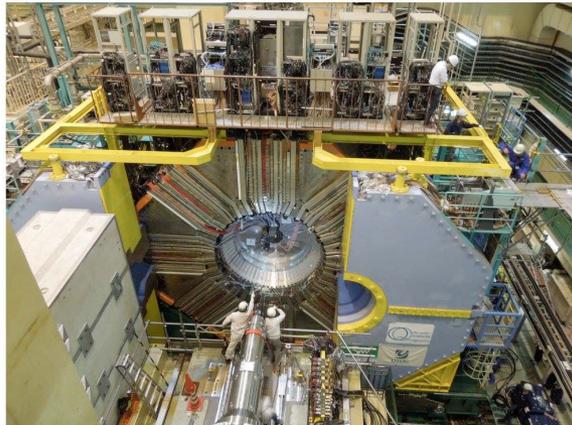
SuperKEKB – major upgrade of the KEKB B-factory at KEK

- ➔ smaller interaction point
- ➔ increased currents
- ➔ collisions at Y(nS)
 - ➔ $E(e^+) = 4 \text{ GeV}$, $E(e^-) = 7 \text{ GeV}$

First collisions in 2016

Belle II detector – upgraded Belle detector

Rolled in in April 2017



Belle II @ SuperKEKB

New facility to search for physics beyond the SM by studying B, D and τ decays and the major player for precise studies of fragmentation

Tsukuba, Japan

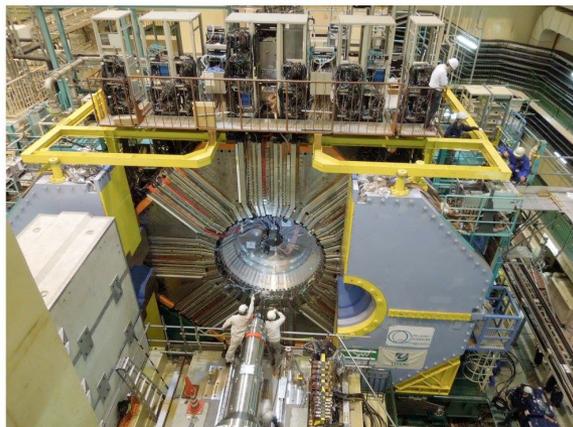
SuperKEKB – major upgrade of the KEKB B-factory at KEK

- ➔ smaller interaction point
- ➔ increased currents
- ➔ collisions at Y(nS)
 - ➔ $E(e^+) = 4 \text{ GeV}$, $E(e^-) = 7 \text{ GeV}$

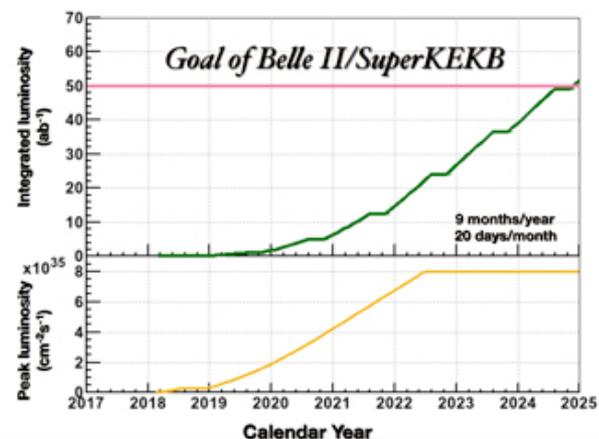
First collisions in 2016

Belle II detector – upgraded Belle detector

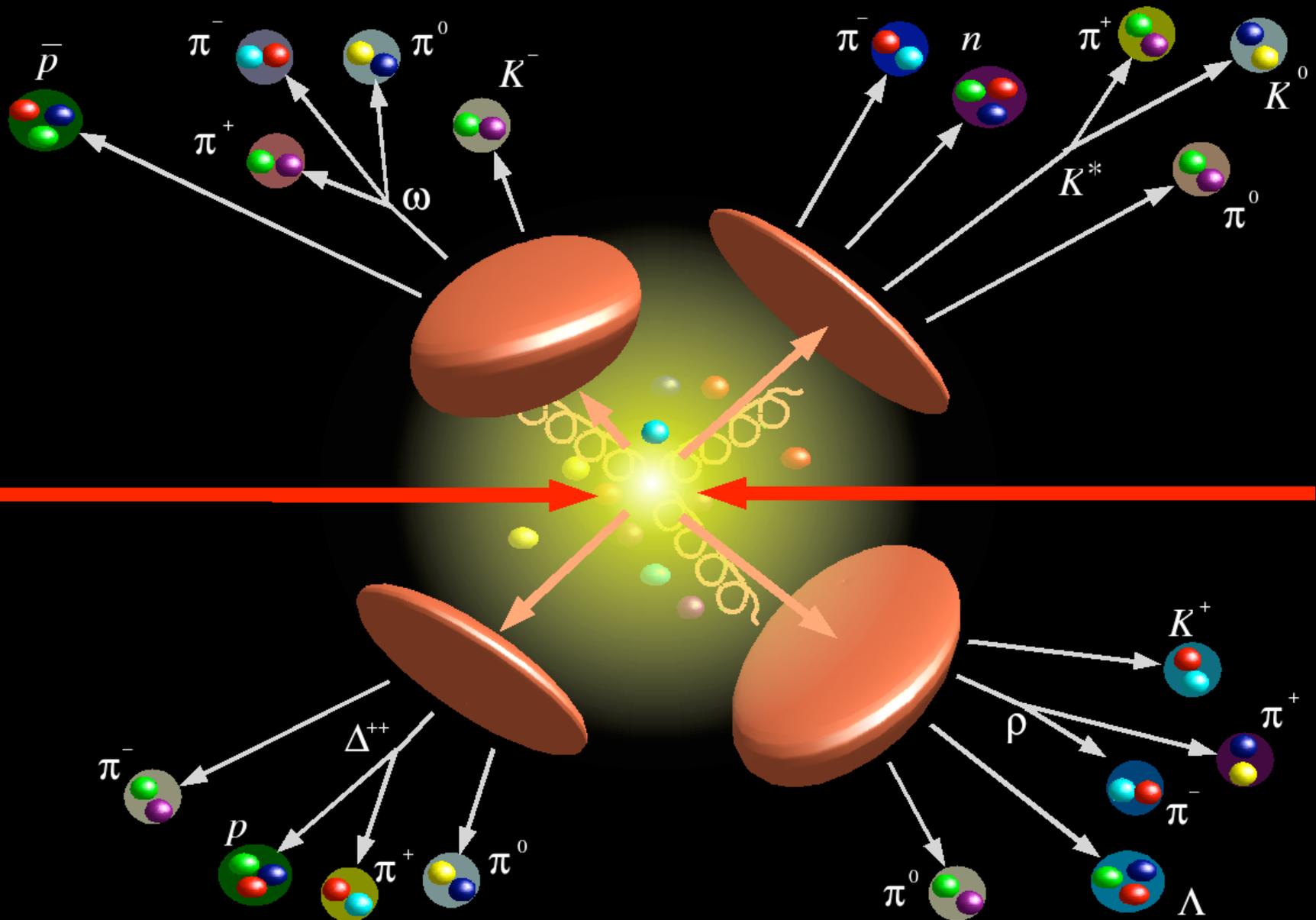
Rolled in in April 2017



SuperKEKB luminosity projection



Hadronisation or fragmentation



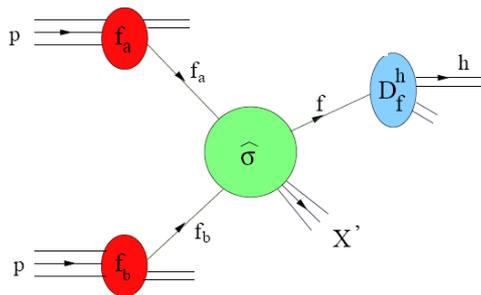
The fragmentation functions

Field, Feynman (1977): Fragmentation functions describe the process of hadronization

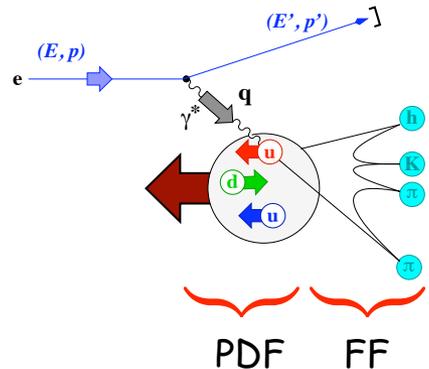
→ non-perturbative but universal objects of QCD

→ $D_q^h(z)$ is the probability that the parton q fragments into a hadron h carrying a fraction z of the parton's momentum

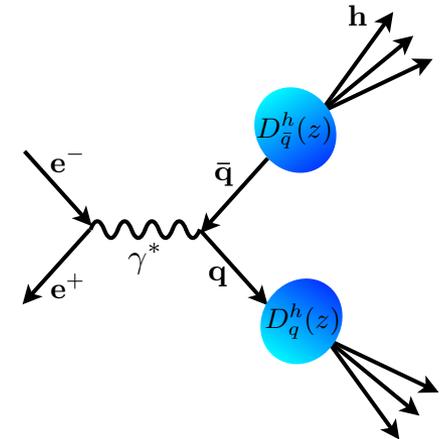
→ proton-proton scattering



→ lepton-hadron scattering



→ e^+e^- annihilation



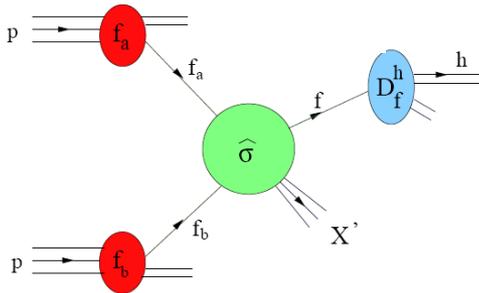
The fragmentation functions

Field, Feynman (1977): Fragmentation functions describe the process of hadronization

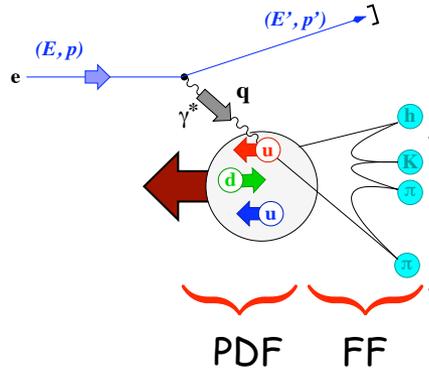
→ non-perturbative but universal objects of QCD

→ $D_q^h(z)$ is the probability that the parton q fragments into a hadron h carrying a fraction z of the parton's momentum

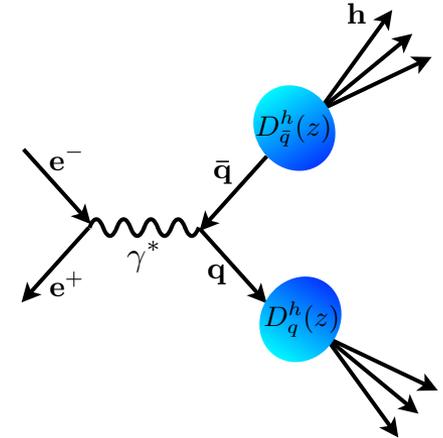
→ proton-proton scattering



→ lepton-hadron scattering



→ e^+e^- annihilation



Measurements of fragmentation functions in lepton-hadron, proton-proton scattering and e^+e^- annihilation are complementary

→ e^+e^- annihilation:

→ direct probe in clean environment

→ stringent constraints on the combinations $D_q^h + D_{\bar{q}}^h$

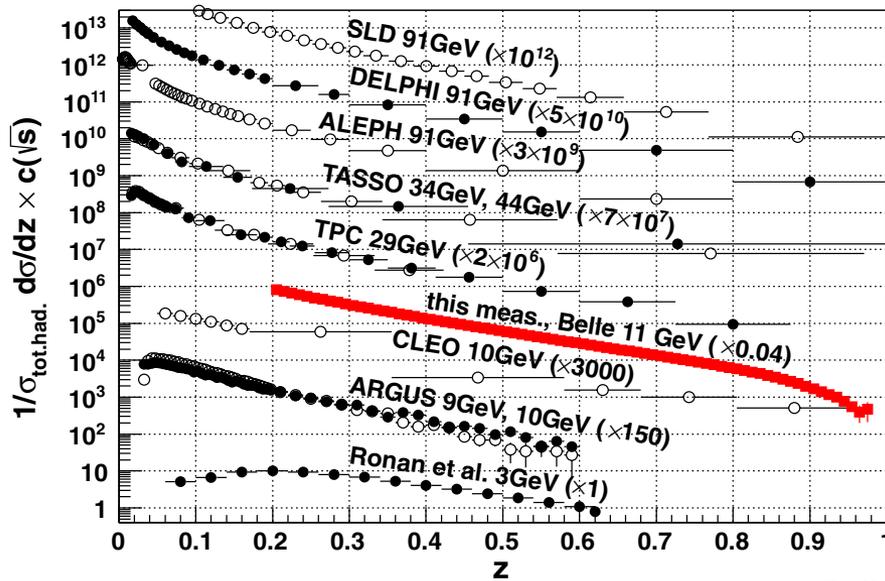
→ insensitive to quark-antiquark separation

→ far less sensitive to D_g^h

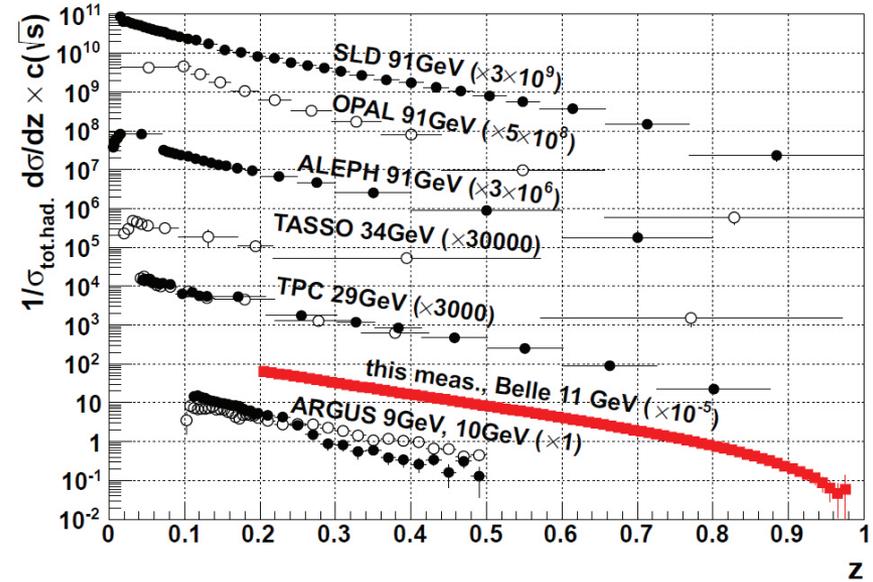
$$\sigma^{e^+e^- \rightarrow hX} \propto \sum_q \sigma^{e^+e^- \rightarrow q\bar{q}} \otimes (D_q^h + D_{\bar{q}}^h)$$

Measurements from e^+e^- annihilation

World Data (Sel.) for $e^+e^- \rightarrow \pi^\pm + X$ Production

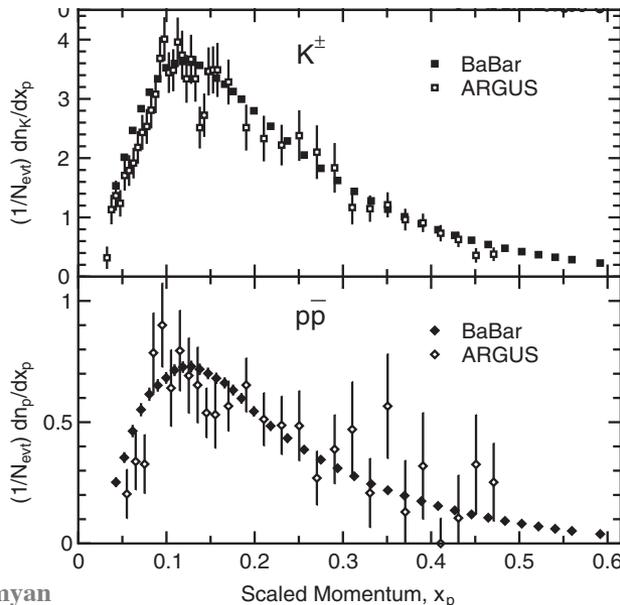


World Data (Sel.) for $e^+e^- \rightarrow K^\pm + X$ Production



- Belle collaboration - Phys Rev D 111 062002 (2013) -

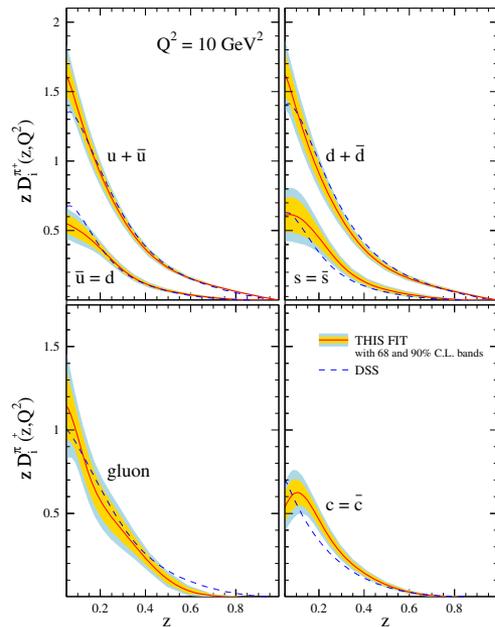
- ➔ large amount of cross section measurements at high \sqrt{s} , near Z^0 mass
- ➔ the region around $\sqrt{s} \sim 10$ GeV remained until recently poorly investigated
- ➔ precise measurements from Belle and BaBar
 - ➔ differential cross-section for π^\pm , K^\pm and p/\bar{p}



- BaBar collaboration - Phys Rev D 88, 032011 (2013) -

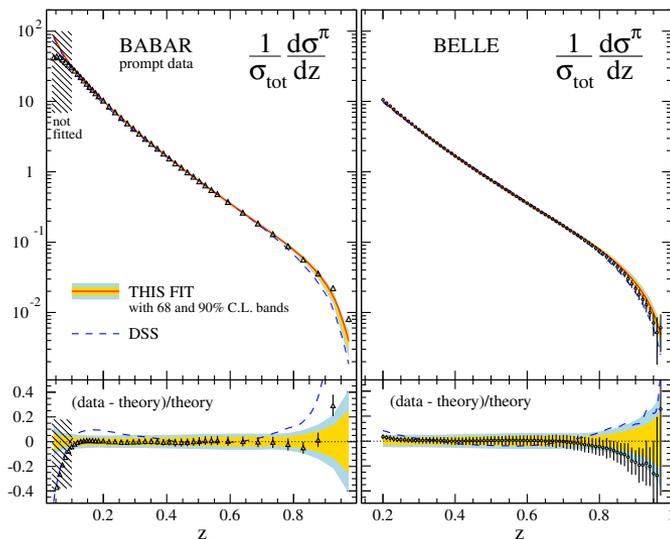
Extraction of fragmentation functions

→ global fit



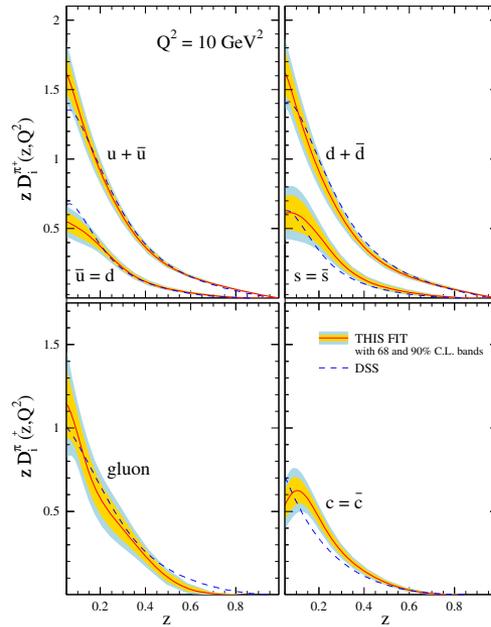
→ good agreement with data

–Phys Rev D91 014035 (2015)–



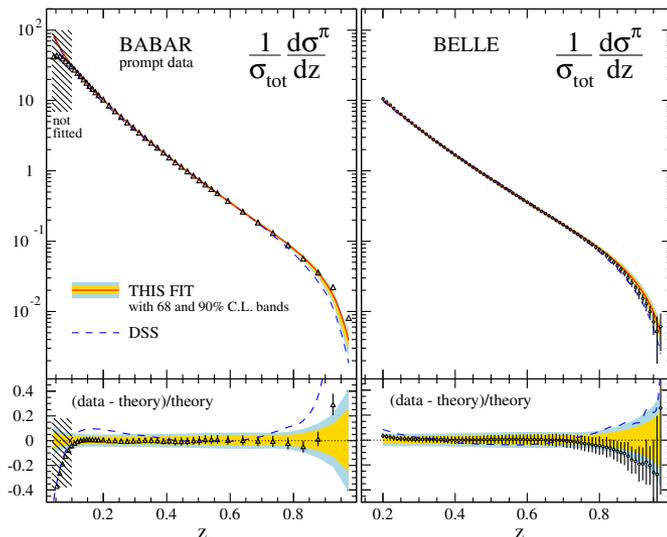
Extraction of fragmentation functions

→ global fit



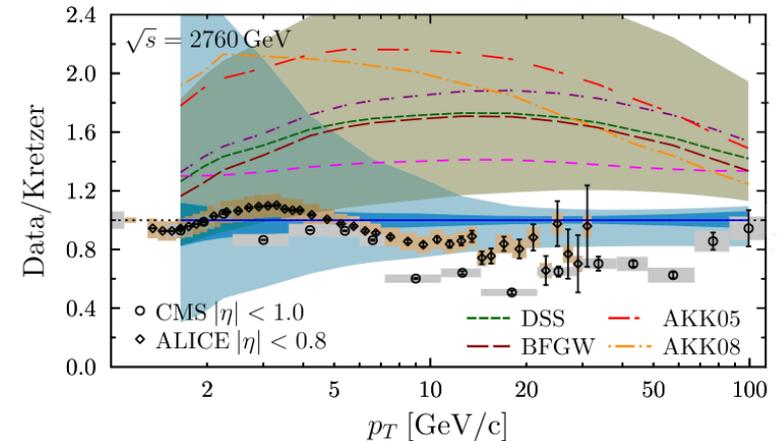
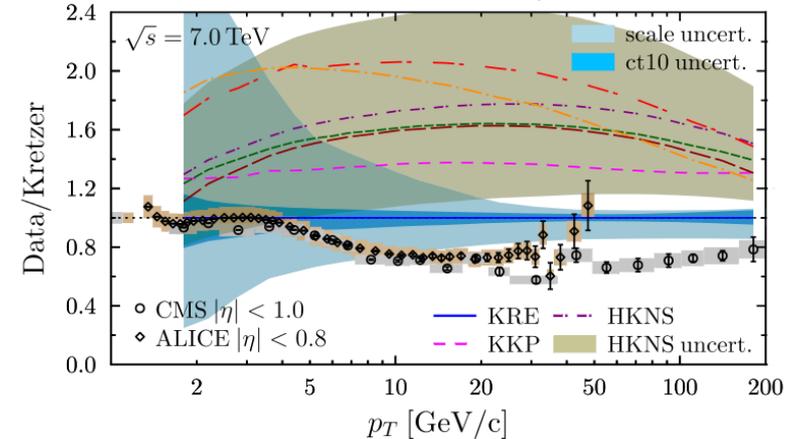
→ good agreement with data

–*Phys Rev D*91 014035 (2015)–



Surprises come with new data

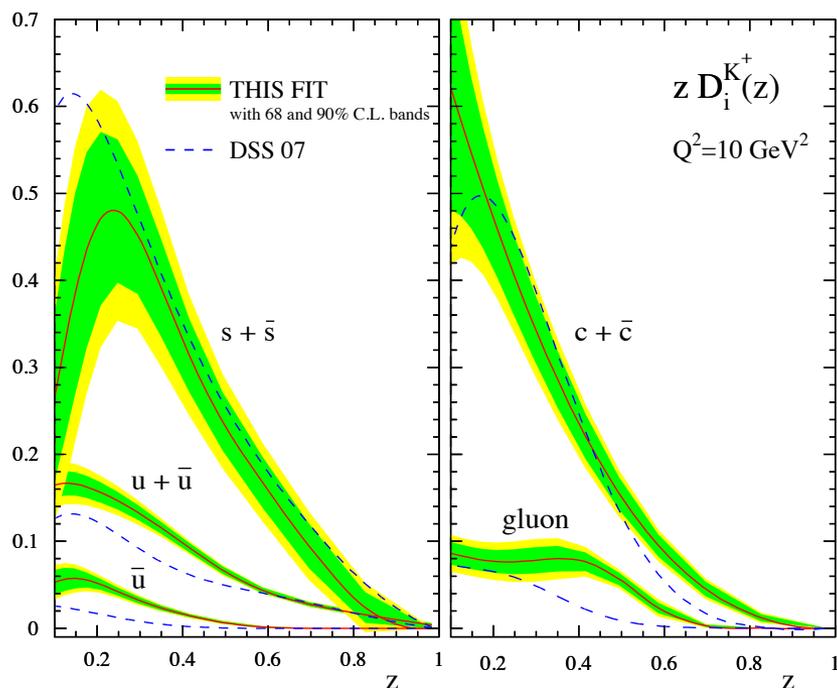
–*Nucl Phys B* 883 (2014) 615–



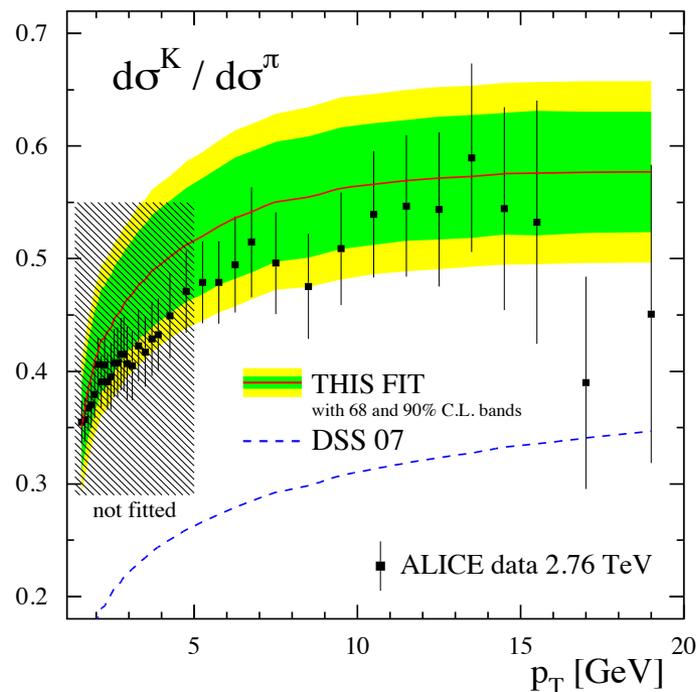
- predictions from all available FF sets are not compatible with CMS and ALICE inclusive charged hadron spectra
- reason might be the gluon-to-hadron FFs

New global fits including the ALICE data

- ➔ the main constraints on FFs come from e^+e^- data
 - ➔ gluon-to-hadron fragmentation functions largely unconstrained
- ➔ gluon FFs constraints from pp data
 - ➔ data from RHIC at small values of $p_T < 5$ GeV
- ➔ inclusion of new data from LHC is mandatory
 - ➔ $d\sigma^K/d\sigma^\pi$ from Alice



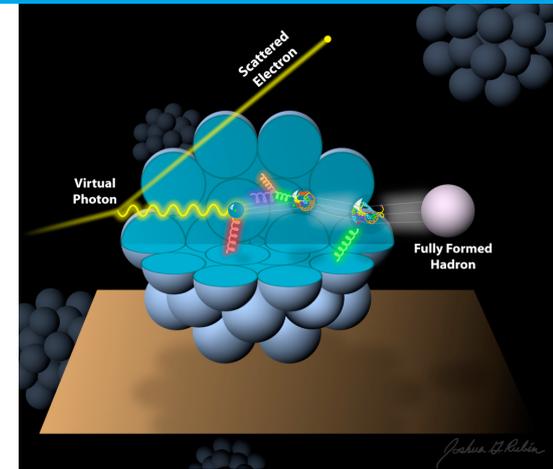
– de Florian et al, 2017 –



- ➔ substantial differences between new and old fits
- ➔ reduced gluon-to-pion fragmentation functions
- ➔ larger gluon-to-kaon fragmentation functions

The Lund string model

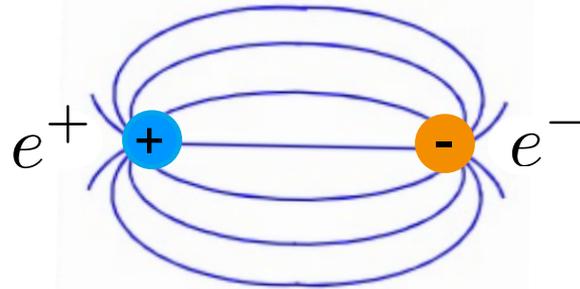
- The fragmentation process has yet to be understood from first principles
- Commonly used phenomenological models of hadronisation:
 - **String model**
 - Cluster model
 - Independent model
 - ...
- All models have a probabilistic nature → probabilistic rules are given for the production of new flavours, and for the sharing of energy and momentum between the products
- String fragmentation model was proposed by X. Artru and G. Mennessier as early as 1974
- Later was developed by **Lund** group
- And still is widely in use in Monte Carlo generators
 - **PYTHIA**
 - Lepto
 - ...
- ... is the basic part of many computer simulations which are necessary to obtain accurate predictions for production of hadrons



The basic picture

- **Asymptotic freedom:** at short distances the force between the quark-antiquark pair becomes increasingly QED-like

QED:



- At long distances, gluon self-interaction makes field lines attract each other
 - gluon interactions described by flux tubes – strings

QCD:



- Approximately constant energy density in the string

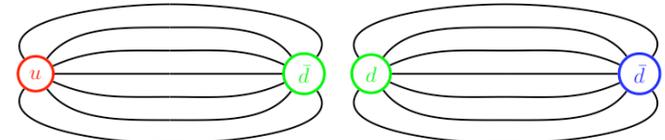
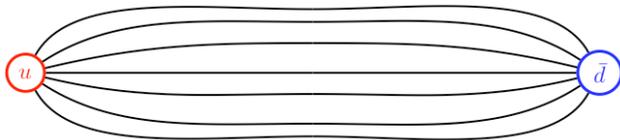
$$F(r) \approx \text{const} = \kappa \approx 1 \text{ GeV/fm}$$

- **Confinement:** linear potential at long distances → needs an infinite energy to drag the quarks apart

$$V(r) \approx \kappa r$$

String fragmentation

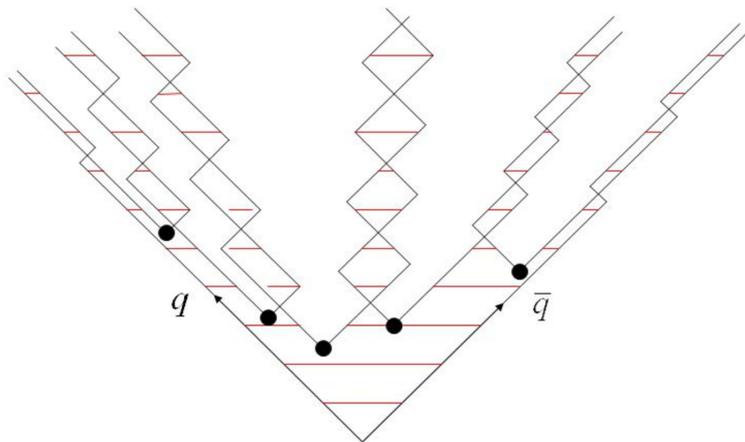
Stretch enough and the strings break apart!



➔ String break by tunnelling:

$$\mathcal{P} \propto \exp\left(\frac{-\pi m_T^2}{\kappa}\right) = \exp\left(\frac{-\pi p_T^2}{\kappa}\right) \exp\left(\frac{-\pi m^2}{\kappa}\right)$$

- ➔ common Gaussian p_T spectrum
- ➔ suppression of heavy-quark production $\rightarrow u : d : s : c \approx 1 : 1 : 0.3 : 10^{-11}$
- ➔ diquark \sim antiquark \Rightarrow simple model for baryon production
- ➔ Motion of quark and antiquark pairs with adjacent pairs forming hadrons

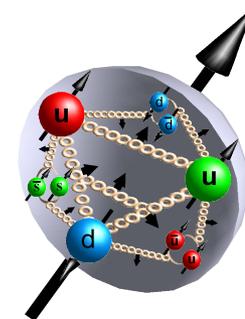


- ➔ generate quark-antiquark pair
- ➔ Gaussian transverse momentum
- ➔ select the hadron state ($L=0$, $S=0$ or 1)
- ➔ determine the longitudinal momentum

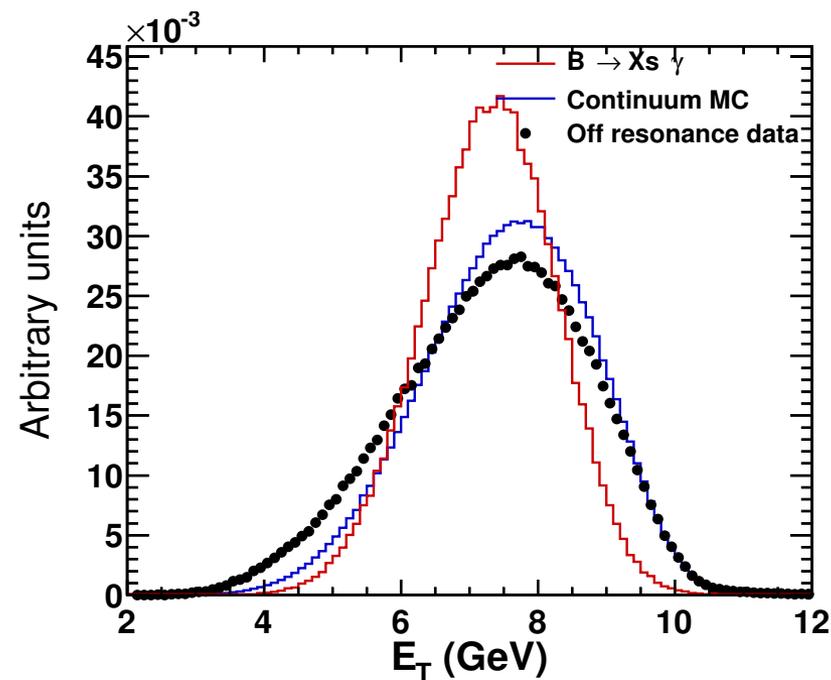
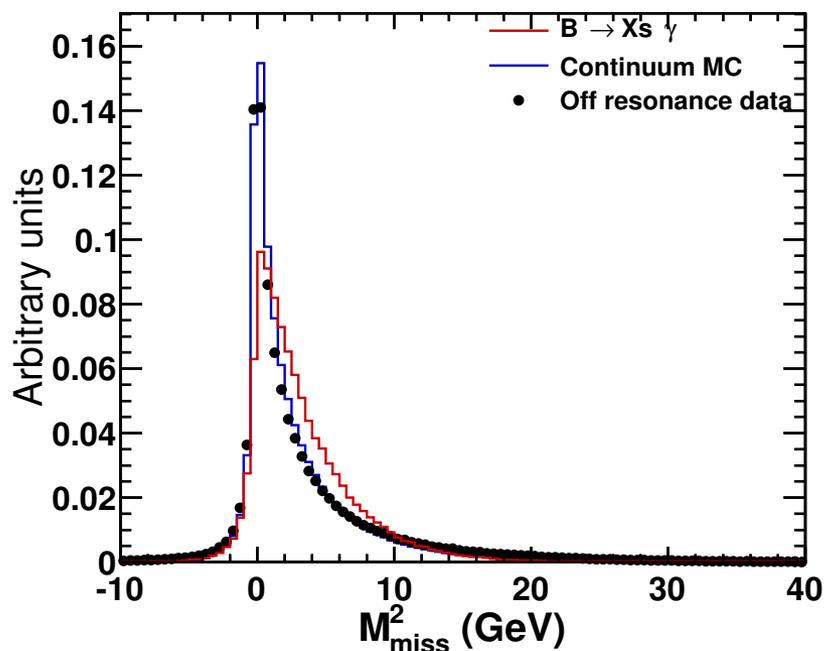
Continuum spectrum

Understanding the hadronization and the continuum spectrum

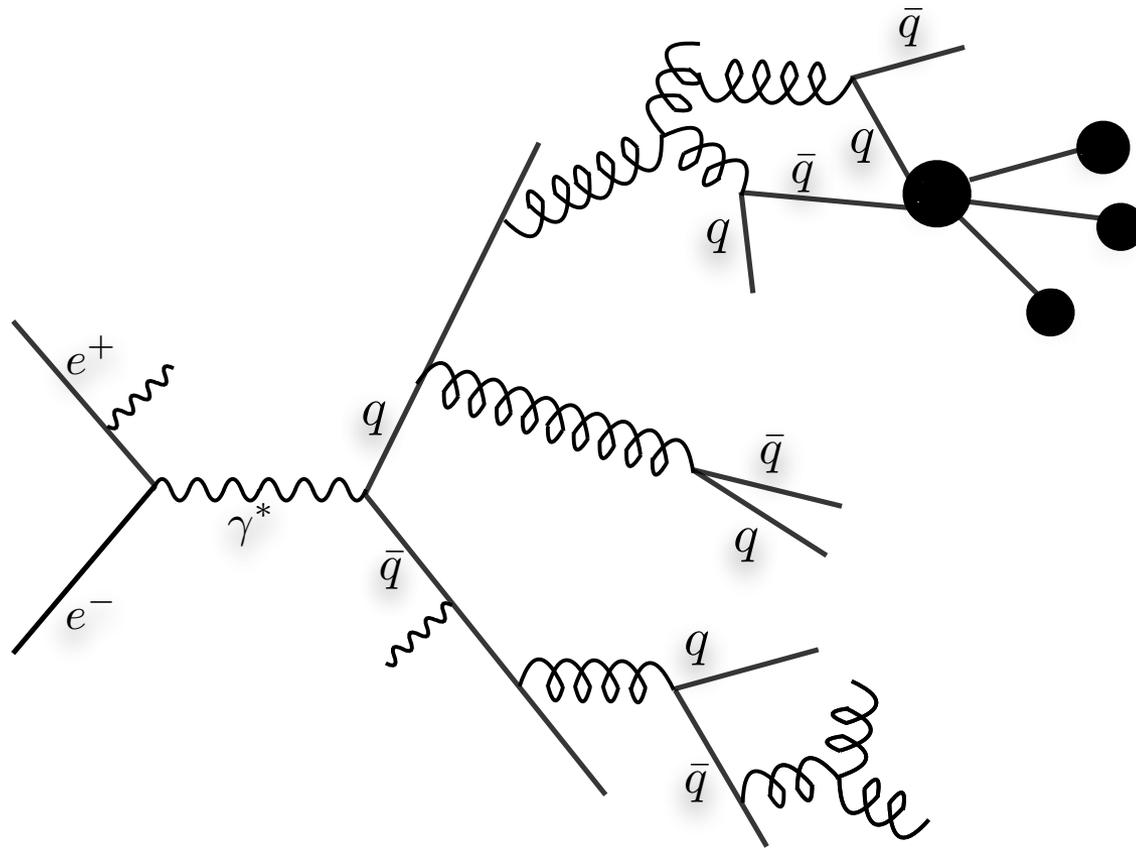
- contribute to all processes with hadrons in the final state
- **particularly important for studying the spin structure of nucleons**
- a tool to study the QCD vacuum structure
- background in NP searches



– Belle collaboration – PRL 114, 151601 (2015) –

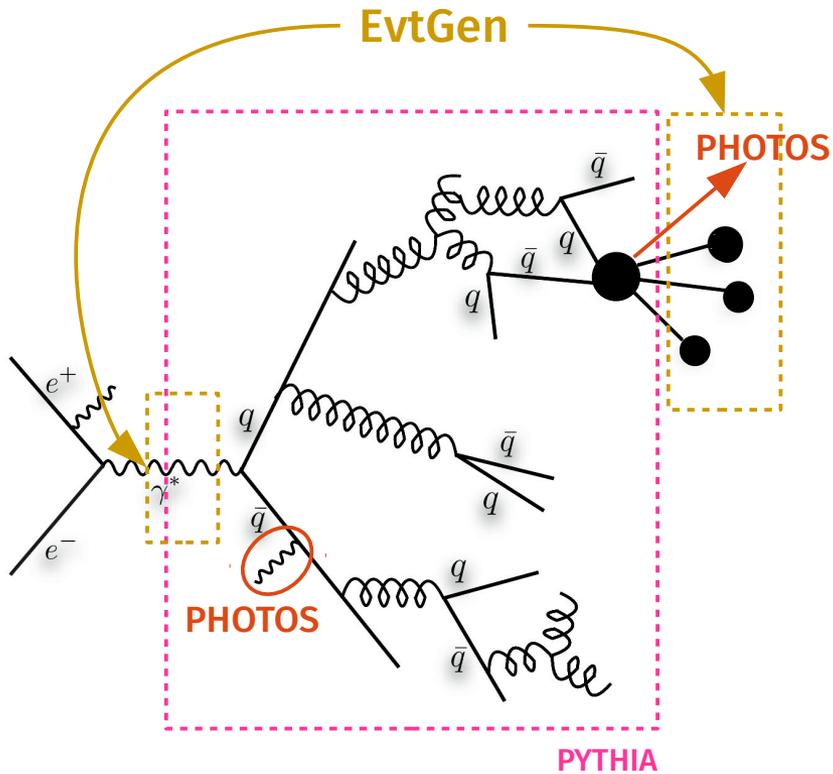


Belle → Belle II: Monte Carlo generators



Belle → Belle II: Monte Carlo generators

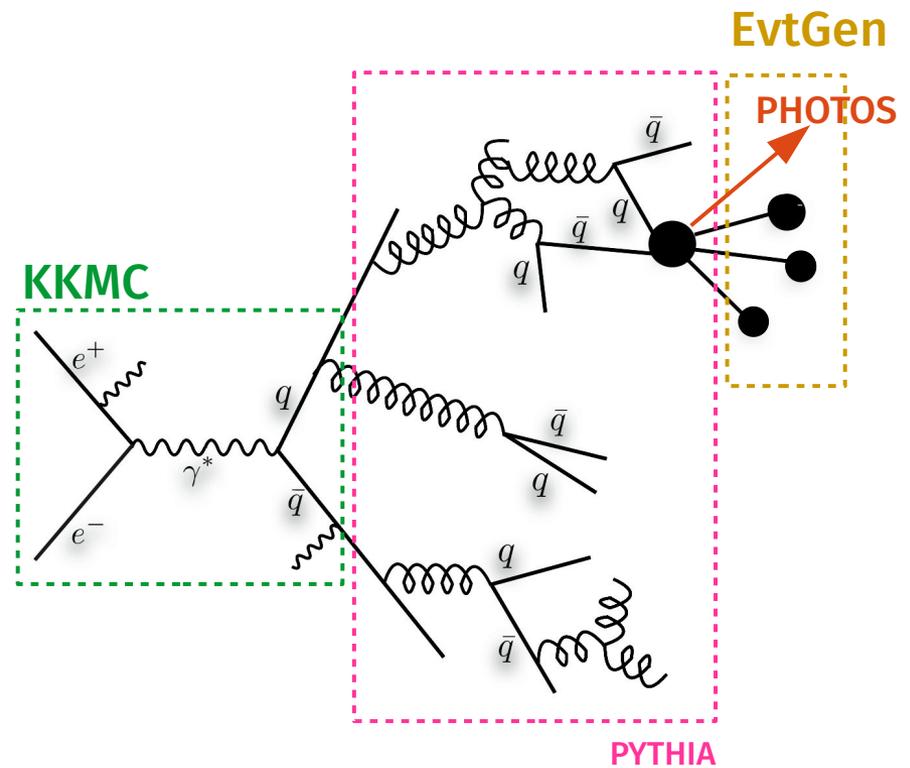
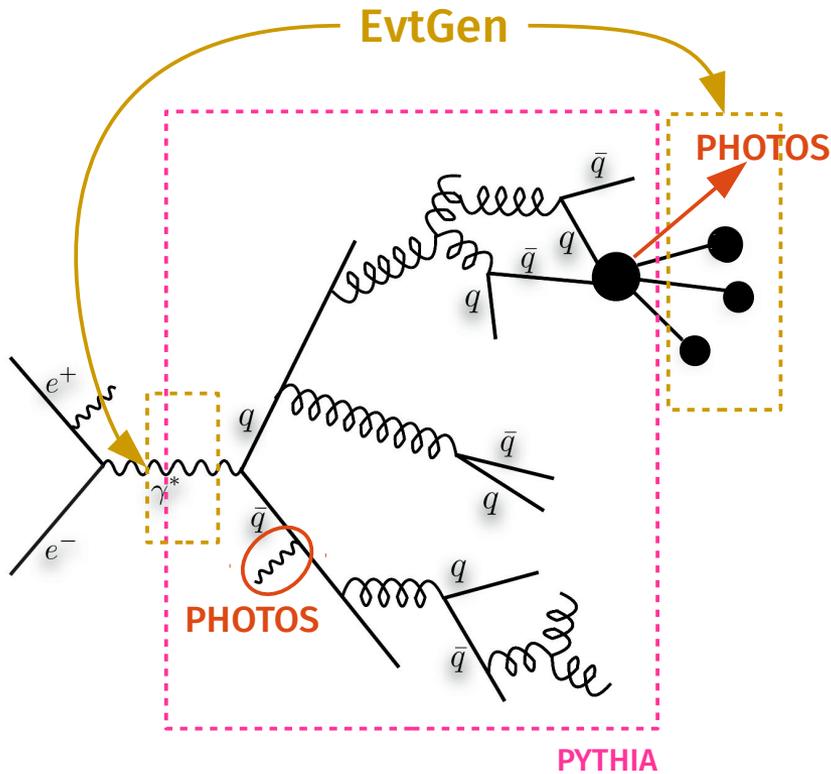
Belle



Belle → Belle II: Monte Carlo generators

Belle

Belle II

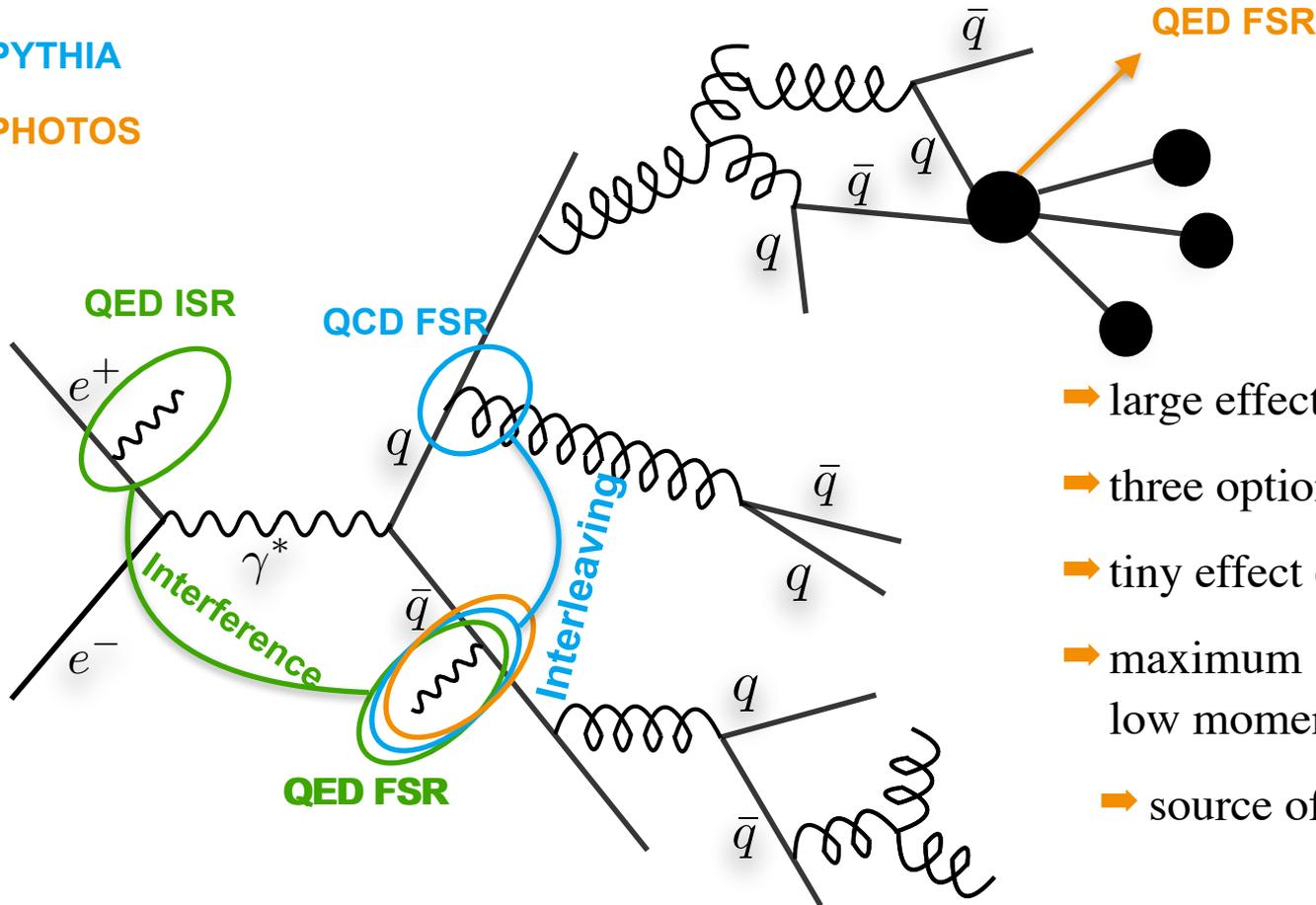


Belle II Monte Carlo generators

KKMC

PYTHIA

PHOTOS



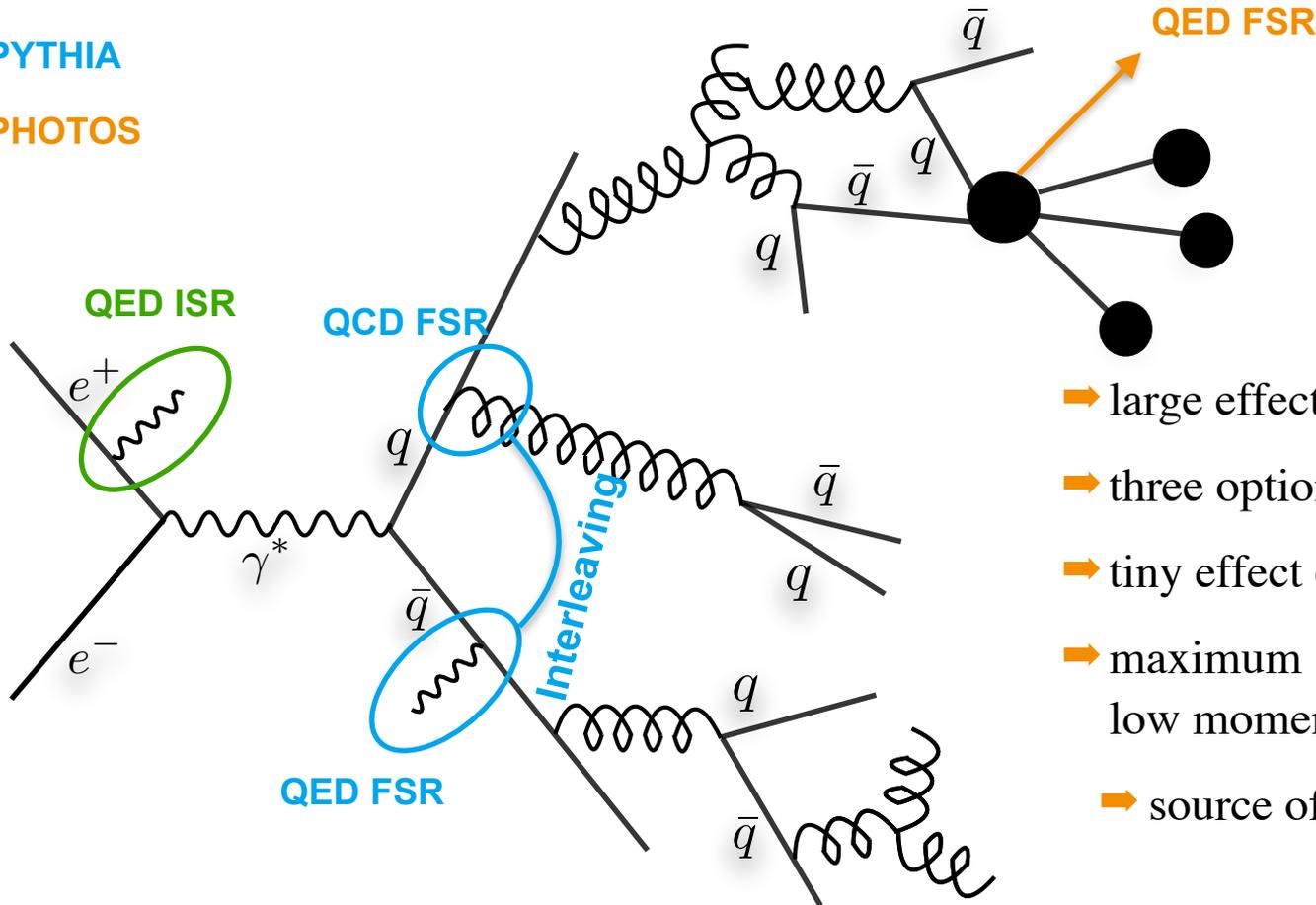
- large effect of ISR
- three options to generate the FSR
- tiny effect of ISR/FSR interference
- maximum 10-20% effect of FSR at low momenta
- source of systematic uncertainty

Belle II Monte Carlo generators

KKMC

PYTHIA

PHOTOS

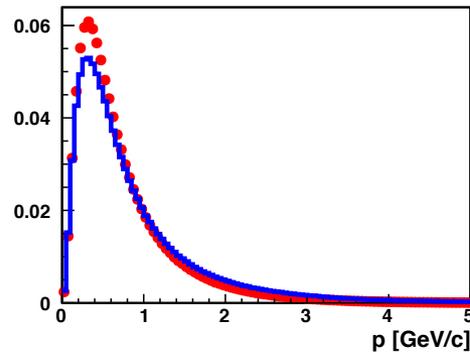
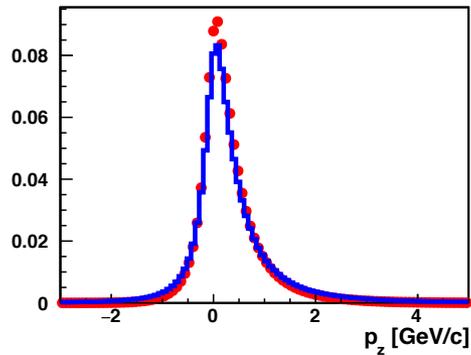
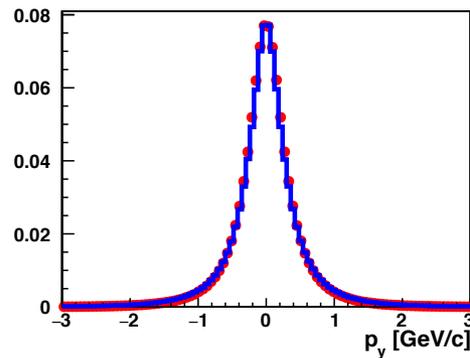
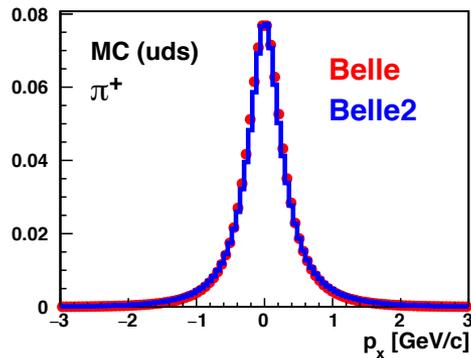


- ➔ large effect of ISR
- ➔ three options to generate the FSR
- ➔ tiny effect of ISR/FSR interference
- ➔ maximum 10-20% effect of FSR at low momenta
- ➔ source of systematic uncertainty

Belle MC % Belle2 MC

Pythia 8 is a clean new start, to provide a successor to Pythia 6.

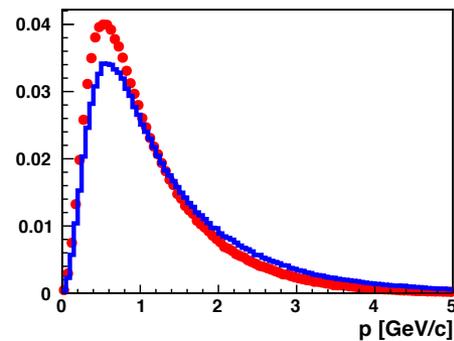
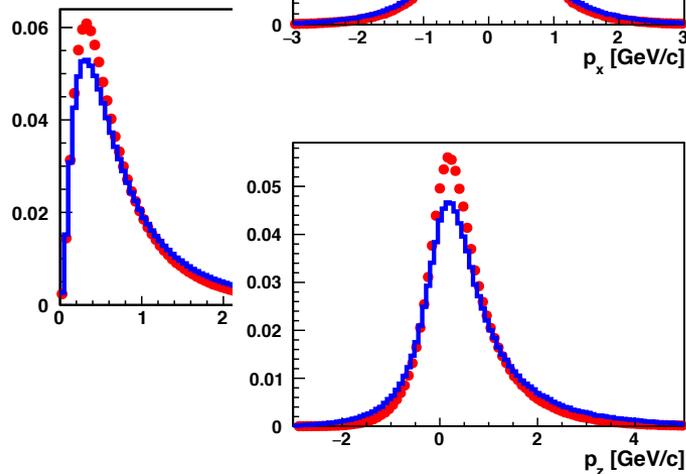
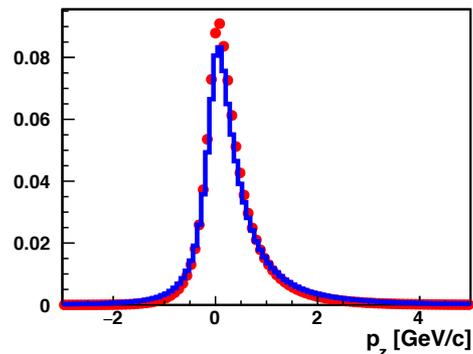
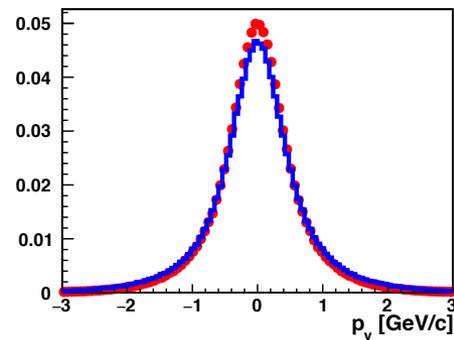
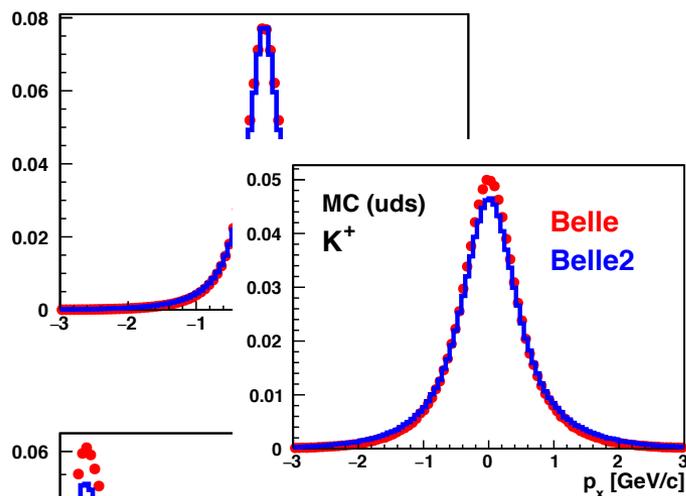
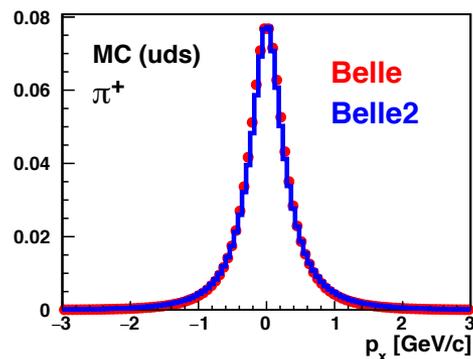
- in general PYTHIA 8 is different from PYTHIA 6
- not possible to "port" a PYTHIA 6 tune to PYTHIA 8



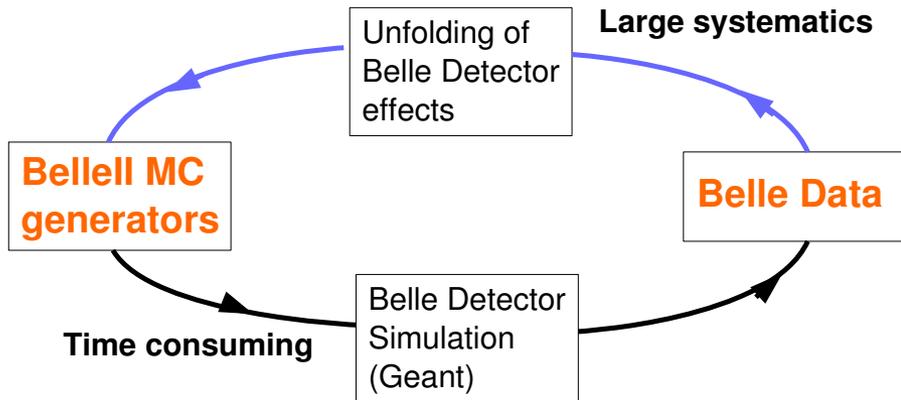
Belle MC % Belle2 MC

Pythia 8 is a clean new start, to provide a successor to Pythia 6.

- in general PYTHIA 8 is different from PYTHIA 6
- not possible to "port" a PYTHIA 6 tune to PYTHIA 8



Towards PYTHIA8 tuning

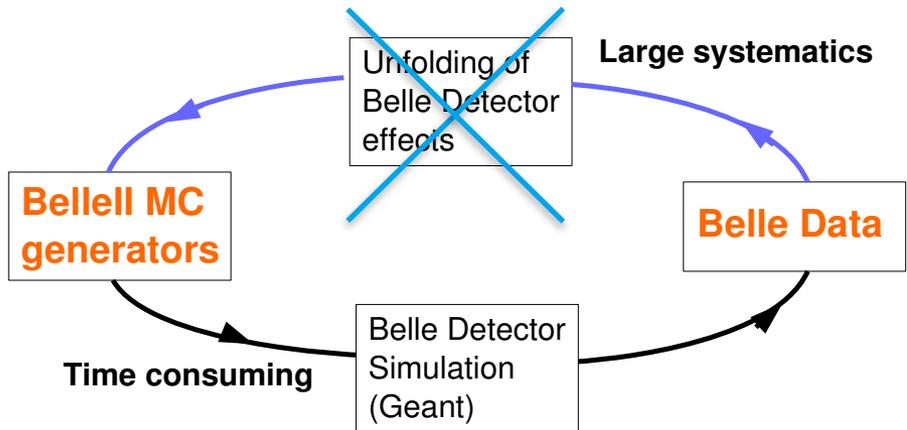


use Belle data

➔ off-resonance data: $q\bar{q}$ events
($q=u, d, s, c$)

use Belle II MC generators

Towards PYTHIA8 tuning

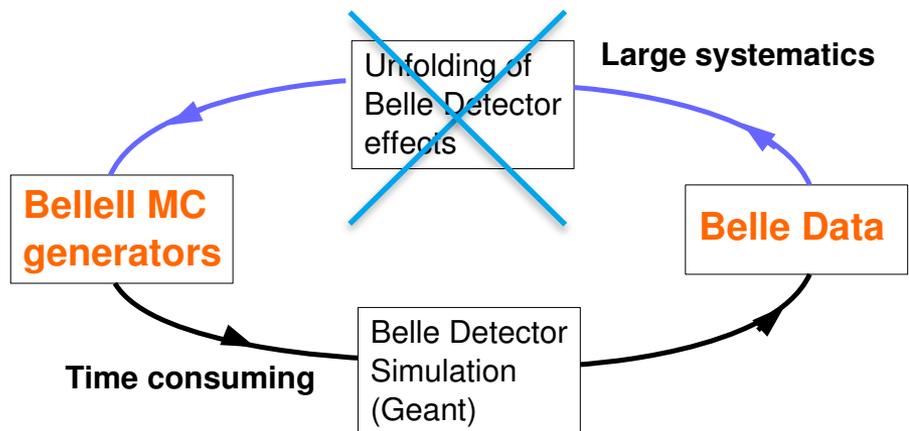


use **Belle data**

➔ off-resonance data: $q\bar{q}$ events
($q=u, d, s, c$)

use **Belle II MC generators**

Towards PYTHIA8 tuning



use Belle data

→ off-resonance data: $q\bar{q}$ events
($q=u, d, s, c$)

use Belle II MC generators

- a tuning tool for Monte Carlo event generators
 - automated tuning approach
 - tune itself is very fast
- professor supplies the parameter grid
- generate Monte Carlo for a given set of parameter values
- calculate observables
- build interpolations
 - parametrise the MC in parameter space with a polynomial
- tune polynomial to data
 - determination of minimum in parameters space



Parameter sensitivity

PYTHIA8 fragmentation $\rightarrow \sim 100$ parameters \rightarrow can't tune all

$$res_i = \frac{r_i - R \hat{p}_i}{\sqrt{R \hat{p}_i \left(\frac{1-R}{R+M} \right) (1-\hat{p}_i)}} \quad \text{with} \quad \hat{p}_i = \frac{r_i + m_i}{R + M}$$

r_i \rightarrow number of events in i^{th} bin in the **reference** sample

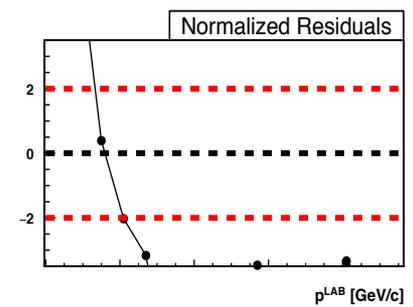
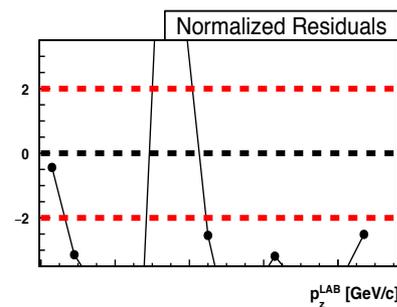
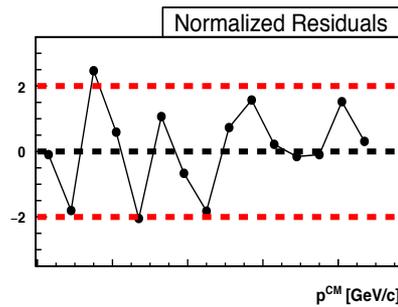
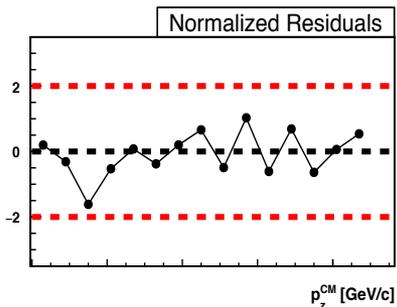
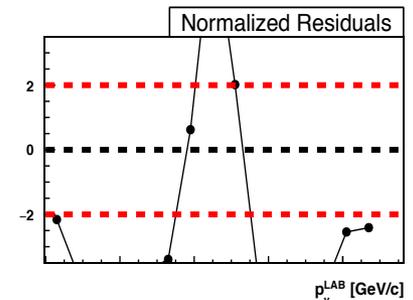
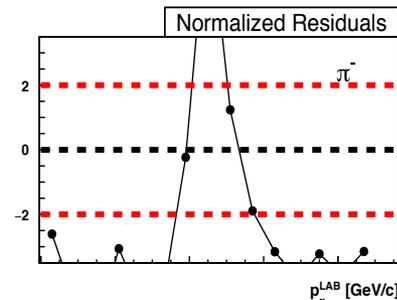
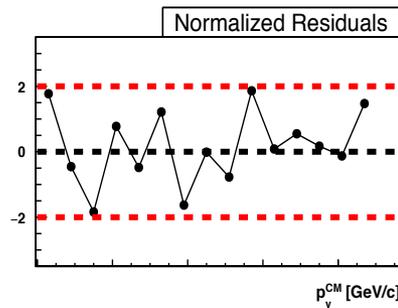
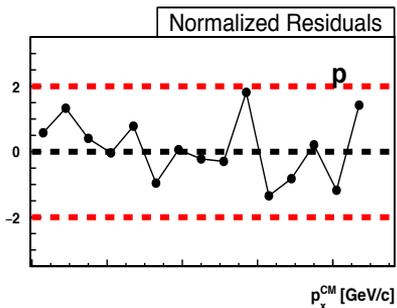
R \rightarrow total number of events in the **reference** sample

m_i \rightarrow number of events in i^{th} bin in the **modified** sample

M \rightarrow total number of events in the **modified** sample

insensitive

sensitive



The parameter list

- Lund symmetric fragmentation functions

$$f(z) = \frac{1}{z} (1 - z)^{\text{StringZ:aLund}} e^{-\text{StringZ:bLund} m_T^2/z}$$

- parameters modifying **s-quark, diquark and c-quark fragmentation**

The parameter list

- Lund symmetric fragmentation functions

$$f(z) = \frac{1}{z} (1 - z)^{\text{StringZ:aLund}} e^{-\text{StringZ:bLund} m_T^2/z}$$

- parameters modifying **s-quark, diquark and c-quark fragmentation**
- the total width in the fragmentation process
$$\langle p_T^2 \rangle_{\text{hadron}} = 2 \text{StringPT : sigma}$$
- parameters to describe the non-Gaussian tail in transverse momentum distribution

The parameter list

- Lund symmetric fragmentation functions

$$f(z) = \frac{1}{z} (1 - z)^{\text{StringZ:aLund}} e^{-\text{StringZ:bLund} m_T^2/z}$$

- parameters modifying **s-quark, diquark and c-quark fragmentation**
- the total width in the fragmentation process
$$\langle p_T^2 \rangle_{\text{hadron}} = 2 \text{StringPT : sigma}$$
- parameters to describe the non-Gaussian tail in transverse momentum distribution
- flavour sector: control the relative rates

The parameter list

- Lund symmetric fragmentation functions

$$f(z) = \frac{1}{z} (1 - z)^{\text{StringZ:aLund}} e^{-\text{StringZ:bLund} m_T^2/z}$$

- parameters modifying **s-quark, diquark and c-quark fragmentation**
- the total width in the fragmentation process
$$\langle p_T^2 \rangle_{\text{hadron}} = 2 \text{StringPT : sigma}$$
- parameters to describe the non-Gaussian tail in transverse momentum distribution
- flavour sector: control the relative rates

→ particles

$\pi^+, \pi^-, \pi^0, \pi^+ \pi^-$

$K^+, K^-, K^+ K^-$

η, η', γ

Λ, \bar{p}

D^0, D_0^*

→ observables

z, p_T, M_{2h}

multiplicities

thrust, R_2

The parameter list

- Lund symmetric fragmentation functions

$$f(z) = \frac{1}{z} (1 - z)^{\text{StringZ:aLund}} e^{-\text{StringZ:bLund} m_T^2/z}$$

- parameters modifying **s-quark, diquark and c-quark fragmentation**
- the total width in the fragmentation process

$$\langle p_T^2 \rangle_{\text{hadron}} = 2 \text{StringPT} : \text{sigma}$$
- parameters to describe the non-Gaussian tail in transverse momentum distribution

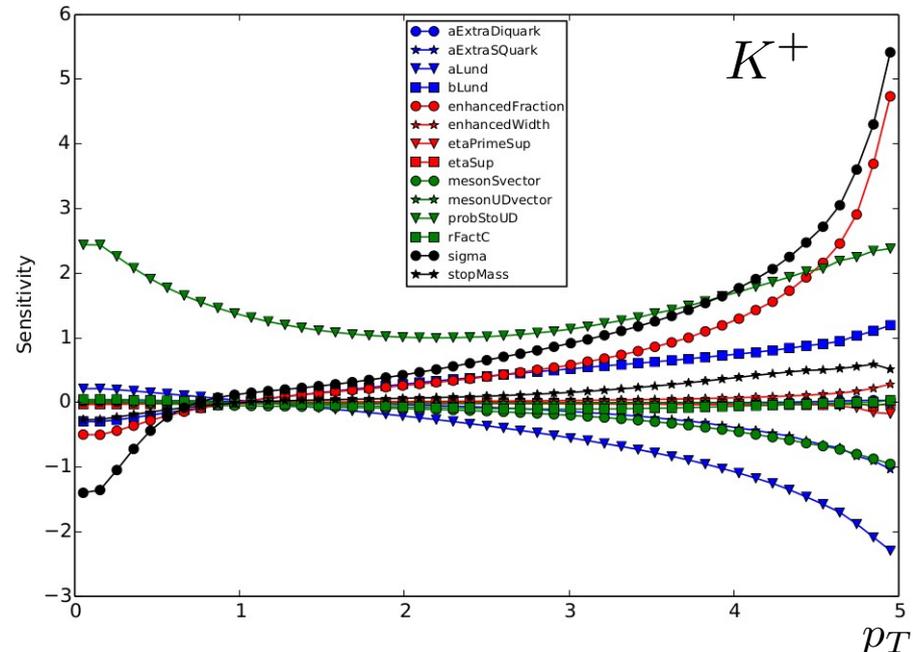
- flavour sector: control the relative rates

→ particles

$\pi^+, \pi^-, \pi^0, \pi^+, \pi^-$
 K^+, K^-, K^+, K^-
 η, η', γ
 Λ, \bar{p}
 D^0, D_0^*

→ observables

z, p_T, M_{2h}
 multiplicities
 thrust, R_2



The parameter list

- Lund symmetric fragmentation functions

$$f(z) = \frac{1}{z} (1 - z)^{\text{StringZ:aLund}} e^{-\text{StringZ:bLund} m_T^2/z}$$

- parameters modifying **s-quark, diquark and c-quark fragmentation**
- the total width in the fragmentation process

$$\langle p_T^2 \rangle_{\text{hadron}} = 2 \text{StringPT} : \text{sigma}$$
- parameters to describe the non-Gaussian tail in transverse momentum distribution

- flavour sector: control the relative rates

→ particles

$\pi^+, \pi^-, \pi^0, \pi^{++}, \pi^-$

K^+, K^-, K^{++}, K^-

η, η', γ

Λ, \bar{p}

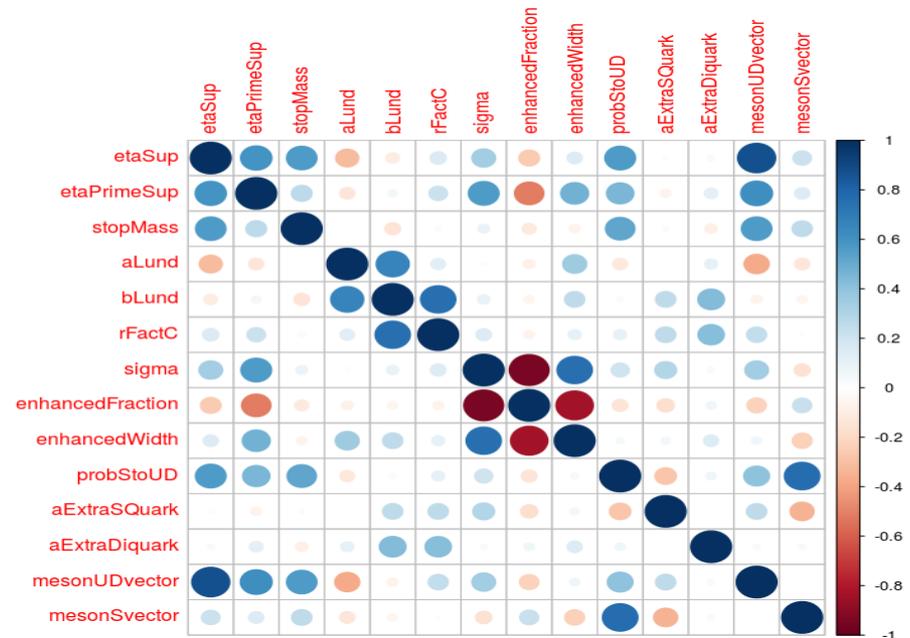
D^0, D_0^*

→ observables

z, p_T, M_{2h}

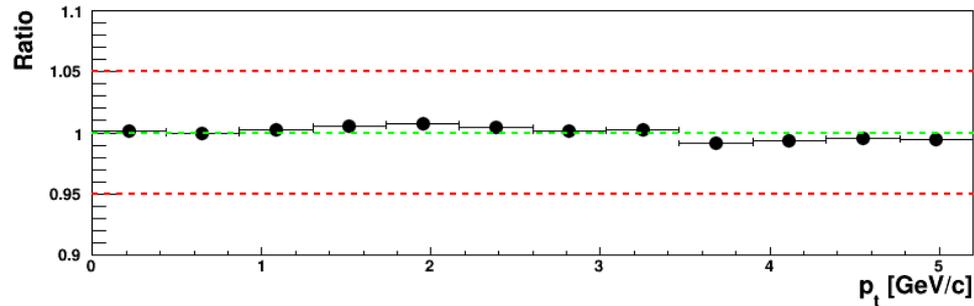
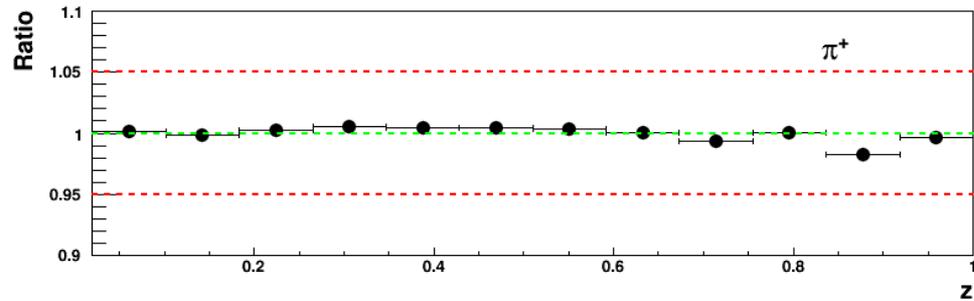
multiplicities

thrust, R_2



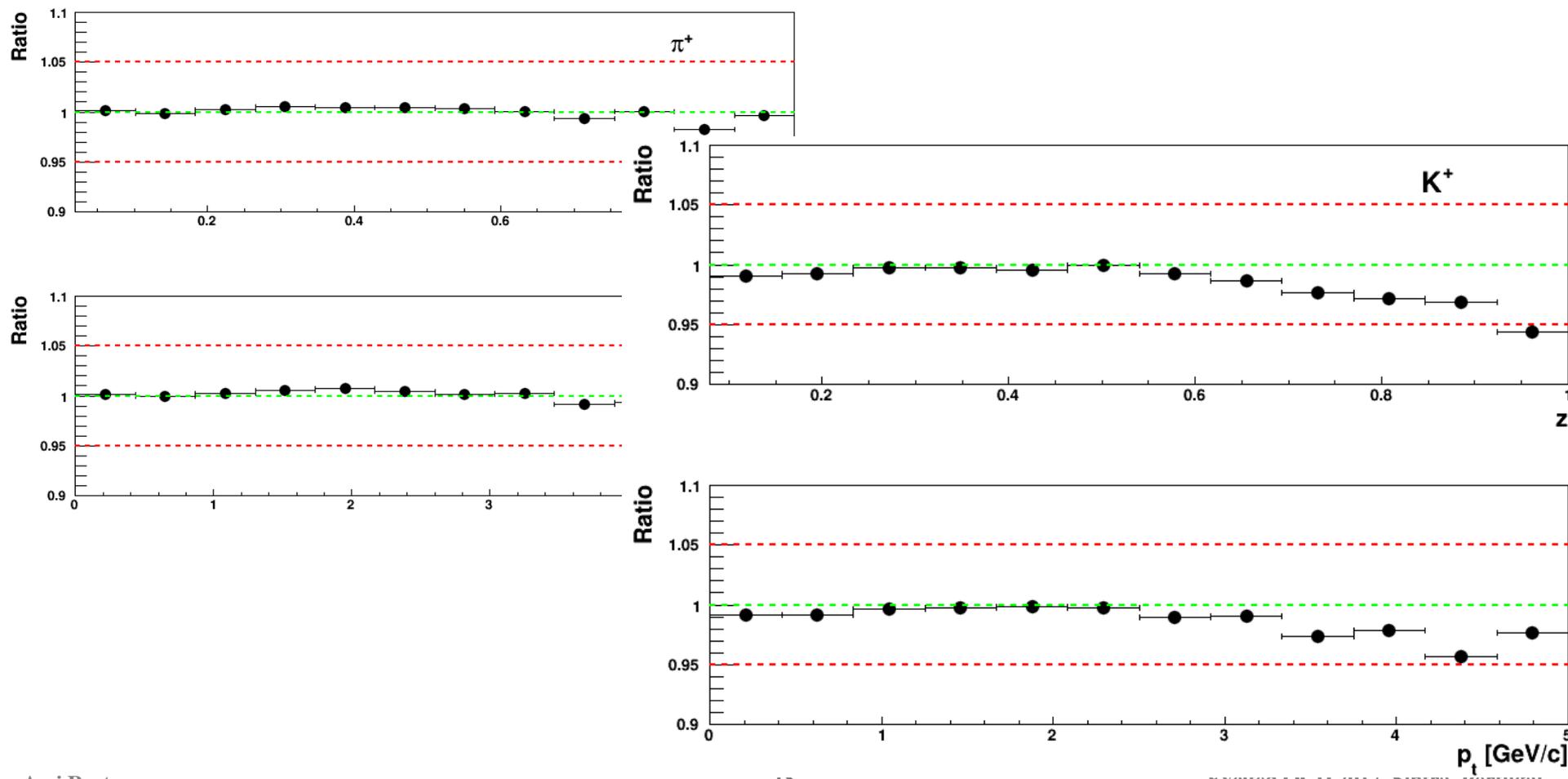
Testing the tuning procedure

- ➔ Simultaneous tuning of all 14 parameters
 - ➔ requires 1362 Monte Carlo samples $\rightarrow \sim 40$ TB
- ➔ Testing the procedure on Monte Carlo “true” kinematics and “true” PID



Testing the tuning procedure

- ➔ Simultaneous tuning of all 14 parameters
 - ➔ requires 1362 Monte Carlo samples $\rightarrow \sim 40$ TB
- ➔ Testing the procedure on Monte Carlo “true” kinematics and “true” PID



Belle II perspectives: fragmentation functions

Rich program to study the fragmentation process

- ➔ release the integration over the transverse-momentum of hadrons
- ➔ extend to different quark or hadron polarisations
- ➔ consider di-hadron production

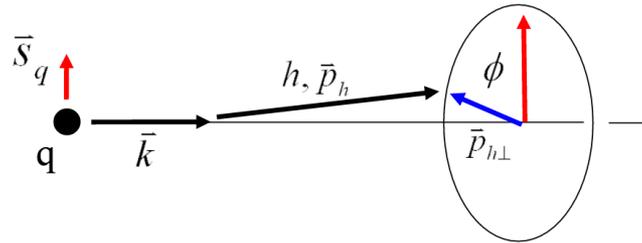
Belle II perspectives: fragmentation functions

Rich program to study the fragmentation process

- release the integration over the transverse-momentum of hadrons
- extend to different quark or hadron polarisations
- consider di-hadron production

1. fragmentation into one spin-0 hadron

$$D_{1,q^\uparrow}^h(z, p_{h\perp}) = D_{1,q}^h(z, p_{h\perp}) + H_{1,q}^\perp(z, p_{h\perp}) \frac{(\hat{\mathbf{k}}_q \times \hat{\mathbf{p}}_{h\perp}) \cdot \mathbf{s}_q}{zM_h}$$



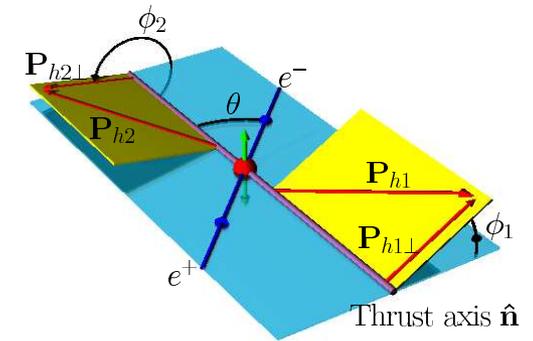
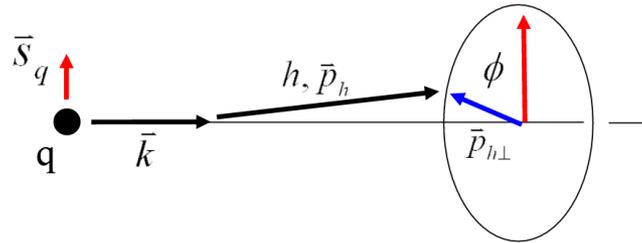
Belle II perspectives: fragmentation functions

Rich program to study the fragmentation process

- ➔ release the integration over the transverse-momentum of hadrons
- ➔ extend to different quark or hadron polarisations
- ➔ consider di-hadron production

1. fragmentation into one spin-0 hadron

$$D_{1,q^\uparrow}^h(z, p_{h\perp}) = D_{1,q}^h(z, p_{h\perp}) + H_{1,q}^\perp(z, p_{h\perp}) \frac{(\hat{\mathbf{k}}_q \times \hat{\mathbf{p}}_{h\perp}) \cdot \mathbf{s}_q}{zM_h}$$



$$\sigma \propto 1 + \frac{\sin^2 \theta}{1 + \cos^2 \theta} \cos(\phi_1 + \phi_2) \frac{H_1^\perp(z_1) \bar{H}_1^\perp(z_2)}{D_1(z_1) \bar{D}_1(z_2)}$$

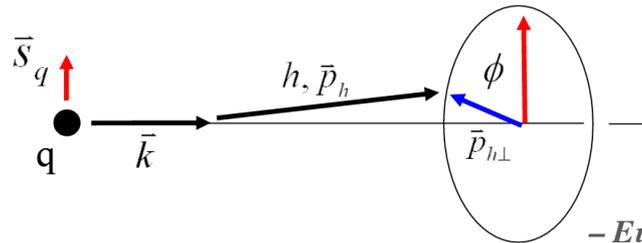
Belle II perspectives: fragmentation functions

Rich program to study the fragmentation process

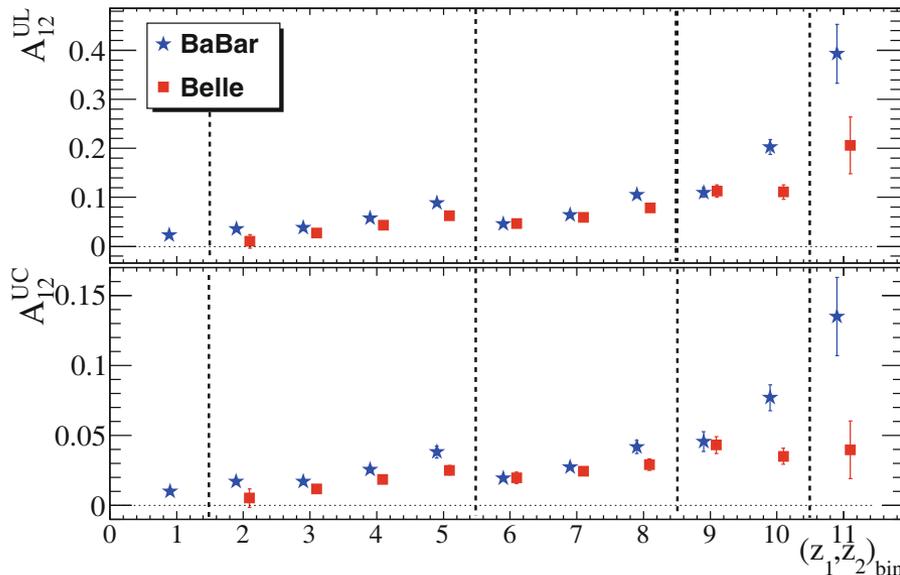
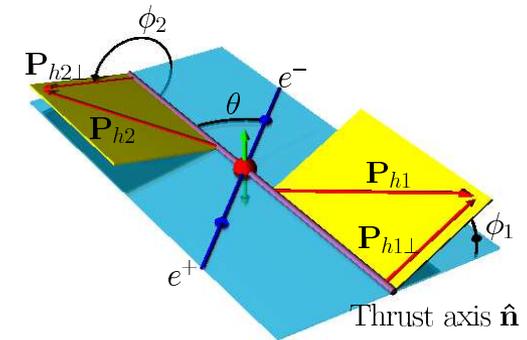
- ➔ release the integration over the transverse-momentum of hadrons
- ➔ extend to different quark or hadron polarisations
- ➔ consider di-hadron production

1. fragmentation into one spin-0 hadron

$$D_{1,q^\uparrow}^h(z, p_{h\perp}) = D_{1,q}^h(z, p_{h\perp}) + H_{1,q}^\perp(z, p_{h\perp}) \frac{(\hat{\mathbf{k}}_q \times \hat{\mathbf{p}}_{h\perp}) \cdot \mathbf{s}_q}{zM_h}$$



– Eur. Phys. J. A (2016) 52: 152 –



$$\sigma \propto 1 + \frac{\sin^2 \theta}{1 + \cos^2 \theta} \cos(\phi_1 + \phi_2) \frac{H_1^\perp(z_1) \bar{H}_1^\perp(z_2)}{D_1(z_1) \bar{D}_1(z_2)}$$

tension between Belle and BaBar results

- ➔ thrust axis corrections
- ➔ background corrections
- ➔ $z < 0.9$ for BaBar, $z < 1$ for Belle

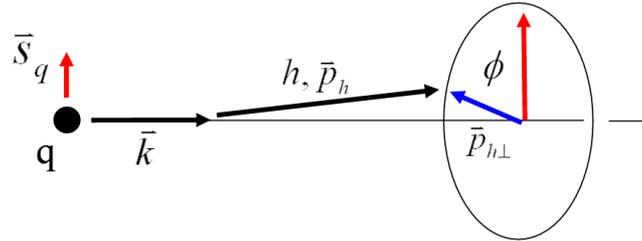
Belle II perspectives: fragmentation functions

Rich program to study the fragmentation process

- release the integration over the transverse-momentum of hadrons
- extend to different quark or hadron polarisations
- consider di-hadron production

1. fragmentation into one spin-0 hadron

$$D_{1,q^\uparrow}^h(z, p_{h\perp}) = D_{1,q}^h(z, p_{h\perp}) + H_{1,q}^\perp(z, p_{h\perp}) \frac{(\hat{\mathbf{k}}_q \times \hat{\mathbf{p}}_{h\perp}) \cdot \mathbf{s}_q}{zM_h}$$



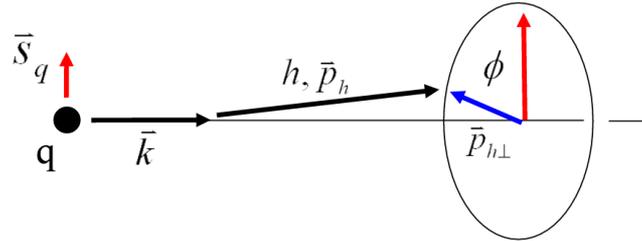
Belle II perspectives: fragmentation functions

Rich program to study the fragmentation process

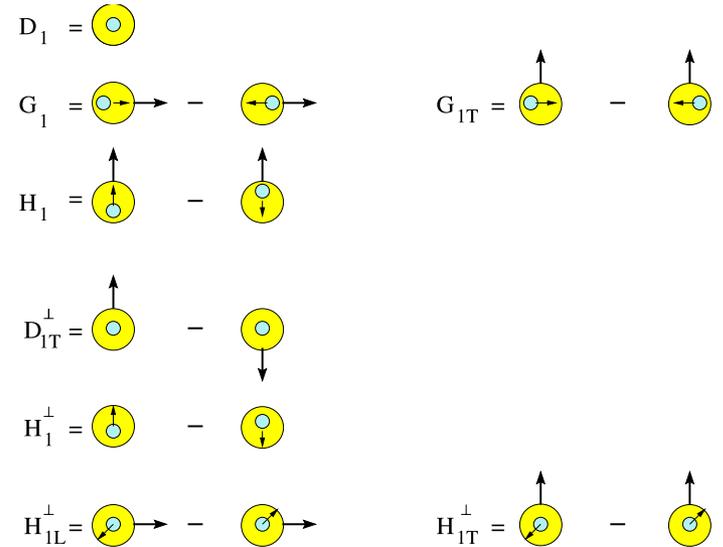
- ➔ release the integration over the transverse-momentum of hadrons
- ➔ extend to different quark or hadron polarisations
- ➔ consider di-hadron production

1. fragmentation into one spin-0 hadron

$$D_{1,q^\uparrow}^h(z, p_{h\perp}) = D_{1,q}^h(z, p_{h\perp}) + H_{1,q}^\perp(z, p_{h\perp}) \frac{(\hat{\mathbf{k}}_q \times \hat{\mathbf{p}}_{h\perp}) \cdot \mathbf{s}_q}{z M_h}$$



2. fragmentation into one spin-1/2 hadron



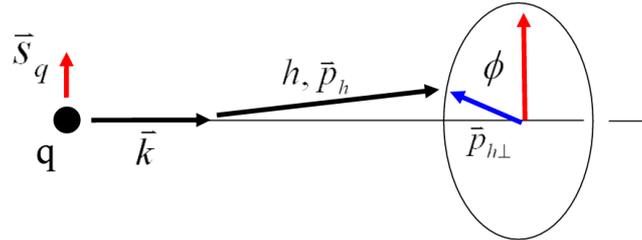
Belle II perspectives: fragmentation functions

Rich program to study the fragmentation process

- ➔ release the integration over the transverse-momentum of hadrons
- ➔ extend to different quark or hadron polarisations
- ➔ consider di-hadron production

1. fragmentation into one spin-0 hadron

$$D_{1,q^\uparrow}^h(z, p_{h\perp}) = D_{1,q}^h(z, p_{h\perp}) + H_{1,q}^\perp(z, p_{h\perp}) \frac{(\hat{\mathbf{k}}_q \times \hat{\mathbf{p}}_{h\perp}) \cdot \mathbf{s}_q}{z M_h}$$



3. fragmentation into two hadrons

$$D_1^{(2h)} = \text{[diagram: blue dot with arrow pointing to two orange dots]} \quad G_1^\perp(2h) = \text{[diagram: blue dot with arrow pointing to two orange dots, minus blue dot with arrow pointing to two orange dots]} \\ H_1^\perp(2h), H_1^\triangleleft = \text{[diagram: blue dot with arrow pointing to two orange dots, minus blue dot with arrow pointing to two orange dots and a red arrow pointing down]}$$

2. fragmentation into one spin-1/2 hadron

$$D_1 = \text{[diagram: yellow circle with blue dot]} \\ G_1 = \text{[diagram: yellow circle with blue dot and arrow pointing right, minus yellow circle with blue dot and arrow pointing left]} \quad G_{1T} = \text{[diagram: yellow circle with blue dot and arrow pointing up, minus yellow circle with blue dot and arrow pointing down]} \\ H_1 = \text{[diagram: yellow circle with blue dot and arrow pointing up, minus yellow circle with blue dot and arrow pointing down]} \\ D_{1T}^\perp = \text{[diagram: yellow circle with blue dot and arrow pointing up, minus yellow circle with blue dot and arrow pointing down]} \\ H_1^\perp = \text{[diagram: yellow circle with blue dot and arrow pointing up, minus yellow circle with blue dot and arrow pointing down]} \\ H_{1L}^\perp = \text{[diagram: yellow circle with blue dot and arrow pointing up and right, minus yellow circle with blue dot and arrow pointing up and left]} \quad H_{1T}^\perp = \text{[diagram: yellow circle with blue dot and arrow pointing up and right, minus yellow circle with blue dot and arrow pointing up and left]}$$

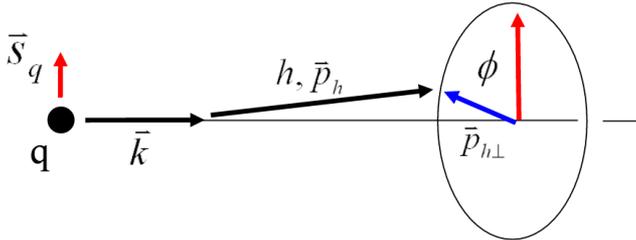
Belle II perspectives: fragmentation functions

Rich program to study the fragmentation process

- ➔ release the integration over the transverse-momentum of hadrons
- ➔ extend to different quark or hadron polarisations
- ➔ consider di-hadron production

1. fragmentation into one spin-0 hadron

$$D_{1,q^\uparrow}^h(z, p_{h\perp}) = D_{1,q}^h(z, p_{h\perp}) + H_{1,q}^\perp(z, p_{h\perp}) \frac{(\hat{\mathbf{k}}_q \times \hat{\mathbf{p}}_{h\perp}) \cdot \mathbf{s}_q}{zM_h}$$



3. fragmentation into two hadrons

$$D_1^{(2h)} = \text{[diagram: two orange dots with a blue dot between them and an arrow pointing right]} \quad G_1^\perp{}^{(2h)} = \text{[diagram: two orange dots with a blue dot between them and an arrow pointing right]} - \text{[diagram: two orange dots with a blue dot between them and an arrow pointing left]}$$

$$H_1^\perp{}^{(2h)}, H_1^\triangleleft = \text{[diagram: two orange dots with a blue dot between them and an arrow pointing up]} - \text{[diagram: two orange dots with a blue dot between them and an arrow pointing down]}$$

2. fragmentation into one spin-1/2 hadron

$$D_1 = \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing up]} \\ G_1 = \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing right]} - \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing left]} \quad G_{1T} = \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing up]} - \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing down]} \\ H_1 = \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing up]} - \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing down]} \\ D_{1T}^\perp = \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing up]} - \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing down]} \\ H_1^\perp = \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing up]} - \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing down]} \\ H_{1L}^\perp = \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing up]} - \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing down]} \quad H_{1T}^\perp = \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing up]} - \text{[diagram: yellow circle with a blue dot in the center and an arrow pointing down]}$$

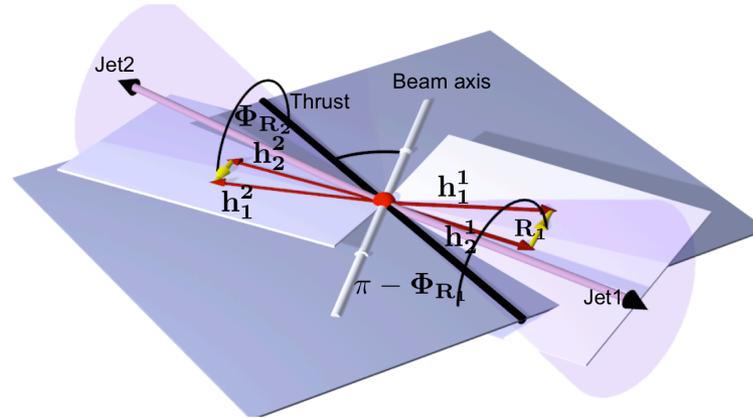
Next generation high precision measurements from Belle II

- ➔ systematic uncertainty under control
- ➔ better vertex reconstruction
- ➔ more precise hadron identification
- ➔ an immense amount of data
- ➔ provide multi-dimensional results (small bin sizes)
- ➔ better constraints on FFs

Belle II perspectives: QCD vacuum structure

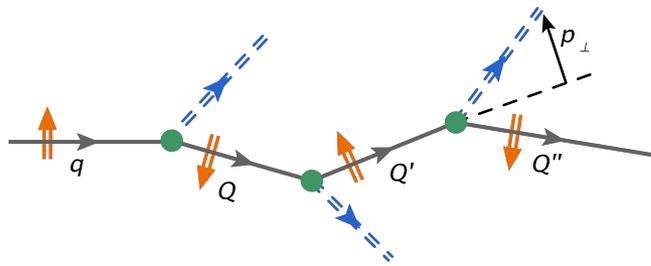
Hadronisation depends on the QCD-vacuum structure!

$$e^+ + e^- \rightarrow \gamma^* \rightarrow q\bar{q} \rightarrow (h_1 h_2) + (\bar{h}_1 \bar{h}_2) + X$$



Quark fragmentation in:

➤ perturbative QCD vacuum

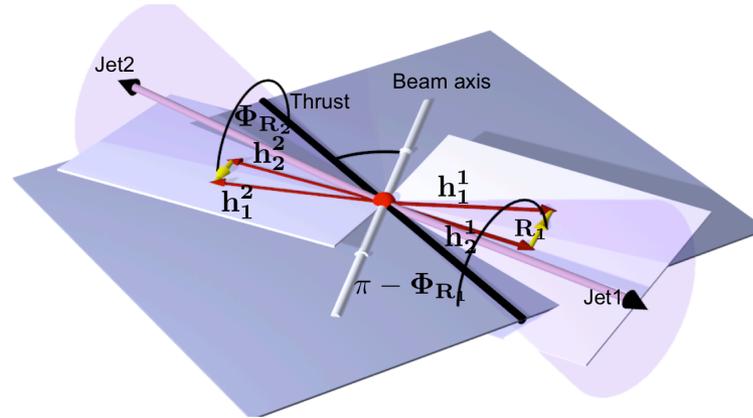


➔ spin-dynamics of hadronisation

Belle II perspectives: QCD vacuum structure

Hadronisation depends on the QCD-vacuum structure!

$$e^+ + e^- \rightarrow \gamma^* \rightarrow q\bar{q} \rightarrow (h_1 h_2) + (\bar{h}_1 \bar{h}_2) + X$$



Quark fragmentation in:

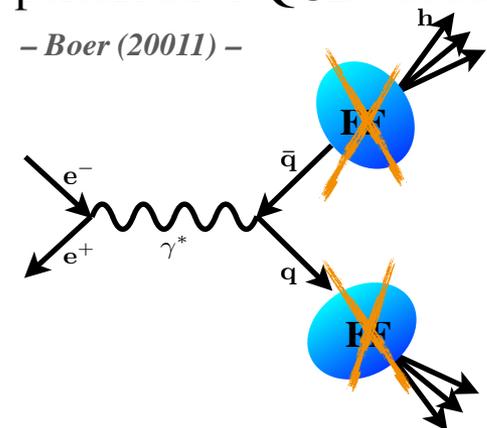
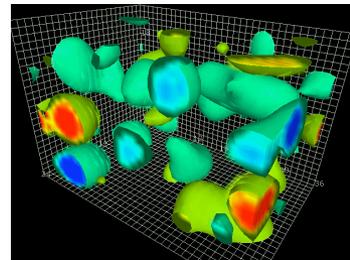
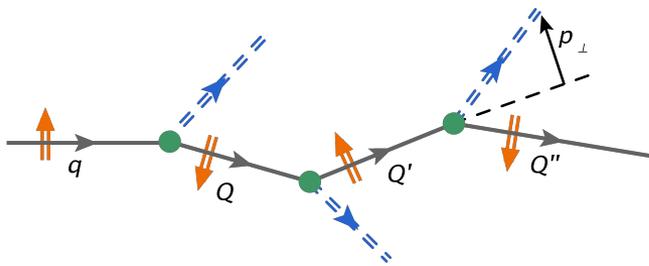
> perturbative QCD vacuum

> θ -vacuum

> non-perturbative QCD vacuum

– Kang, Kharzeev (2011) –

– Boer (2001) –



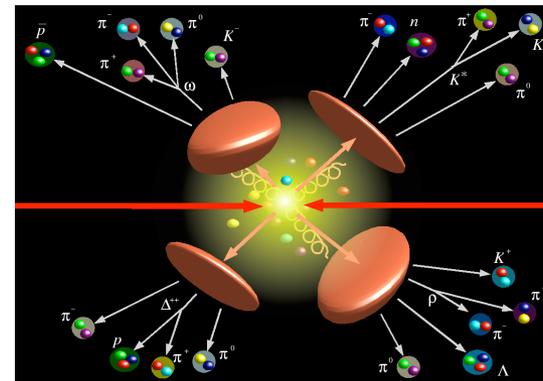
→ spin-dynamics of hadronisation

→ local CP violation

→ factorisation breaking

Outlook

- ➔ Long way ahead before the fragmentation functions will be as precise as the PDFs
- ➔ The spin-dependent and transverse momentum unintegrated fragmentation functions even less constraint
- ➔ Belle II will be a major player in the near future



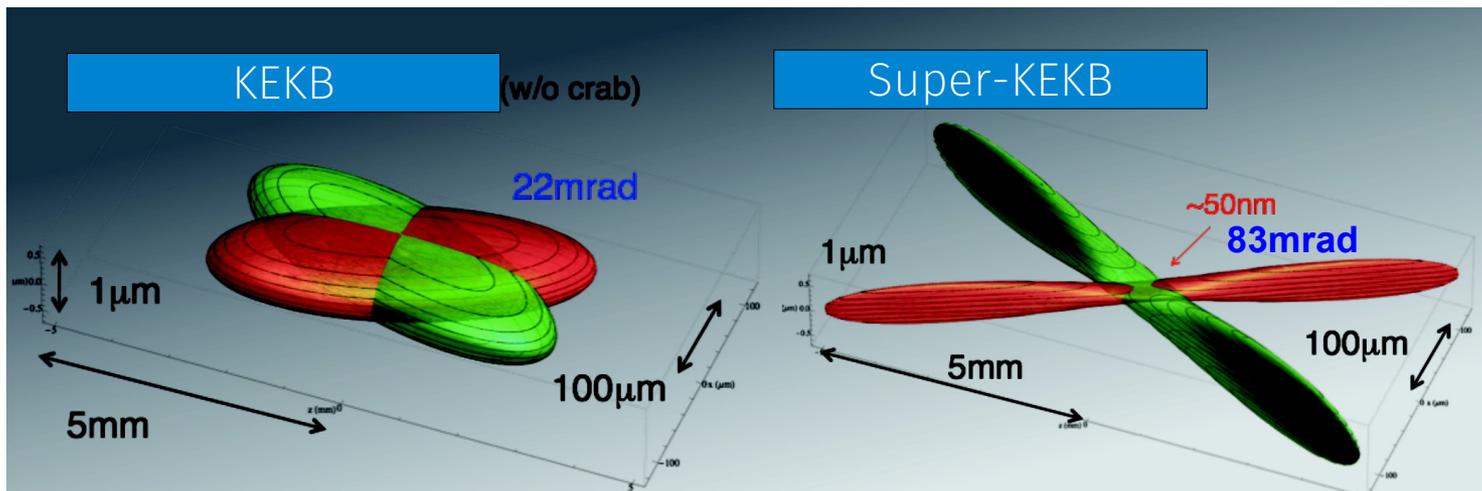
- ➔ The framework for generating the continuum spectrum is ready and validated
- ➔ The strategy for tuning of PYTHIA8 has been settled
- ➔ Tuning using the Belle off-resonance data is in progress
- ➔ As soon as the Belle II data is available, tune using the Belle II data
 - ➔ systematics dominated: both from experimental and phenomenological part

Backups

KEKB → SuperKEKB

$$L = \frac{\gamma_{\pm}}{2er_e} \left(1 + \frac{\sigma_y^*}{\sigma_x^*} \right) \frac{I_{\pm} \xi_{y\pm} R_L}{\beta_{y\pm} R_{\xi_y}}$$

Lorentz factor γ_{\pm}
 classical electron radius er_e
 beam size ratio at IP 1-2% (flat beam) $\frac{\sigma_y^*}{\sigma_x^*}$
 beam current I_{\pm}
 beam-beam parameter $\xi_{y\pm}$
 geometrical reduction factors (crossing angle, hourglass effect) R_L, R_{ξ_y}
 vertical beta function at IP $\beta_{y\pm}$



	E (GeV) LER/HER	β_y^* (mm) LER/HER	β_x^* (cm) LER/HER	ϕ (mrad)	I (A) LER/HER	L (cm ⁻² s ⁻¹)
KEKB	3.5/8.0	5.9/5.9	120/120	11	1.6/1.2	2.1 × 10 ³⁴
SuperKEKB	4.0/7.0	0.27/0.30	3.2/2.5	41.5	3.6/2.6	80 × 10 ³⁴

factor 20

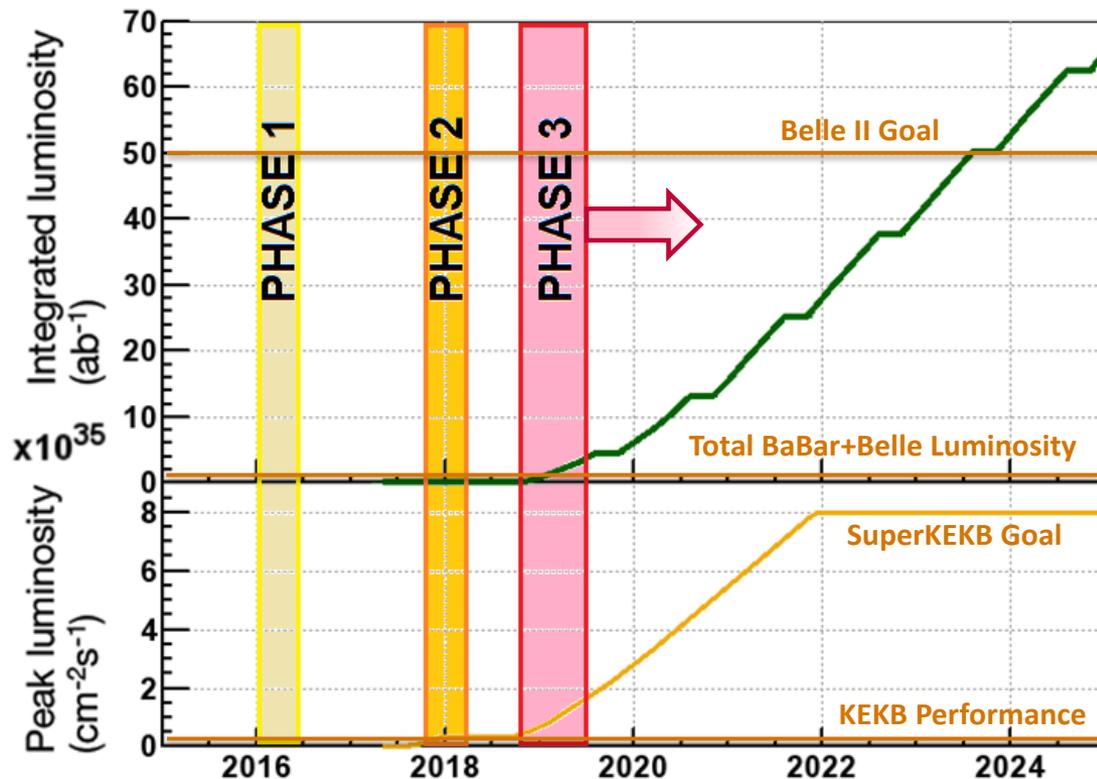
factor 2-3

factor 40



Belle II Schedule

- Phase 1 (2016, complete):
 - Accelerator commissioning
 - No detector
- Phase 2 (start of 2018):
 - Partial detector
 - Background studies
 - First physics
- Phase 3 (end of 2018):
 - Full detector
 - Belle II run



Methodology

$$e^+ + e^- \rightarrow \gamma^* \rightarrow q\bar{q} \rightarrow (h_1 h_2) + (\bar{h}_1 \bar{h}_2) + X$$

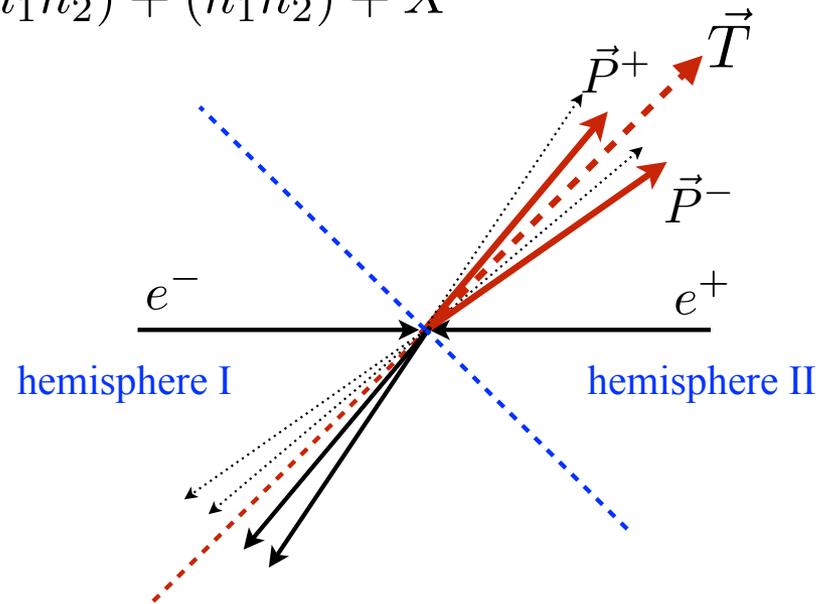
> “left” pairs:

→ $\Omega < 0$

> “right” pairs:

→ $\Omega > 0$

$$\Omega = \frac{(\vec{P}^+ \times \vec{P}^-) \cdot \vec{T}}{|\vec{P}^+| |\vec{P}^-|}$$



The experimental observables for disentangling the hadronization mechanisms :

> jet handedness

$$H = \frac{N_R(\Omega > 0) - N_L(\Omega < 0)}{N_R(\Omega > 0) + N_L(\Omega < 0)}$$

jet handedness correlation

$$C = \frac{\langle \Omega_1 \Omega_2 \rangle - \langle \Omega_1 \rangle \langle \Omega_2 \rangle}{\sqrt{\langle \Omega_1^2 \rangle - \langle \Omega_1 \rangle^2} \sqrt{\langle \Omega_2^2 \rangle - \langle \Omega_2 \rangle^2}}$$

Methodology

$$e^+ + e^- \rightarrow \gamma^* \rightarrow q\bar{q} \rightarrow (h_1 h_2) + (\bar{h}_1 \bar{h}_2) + X$$

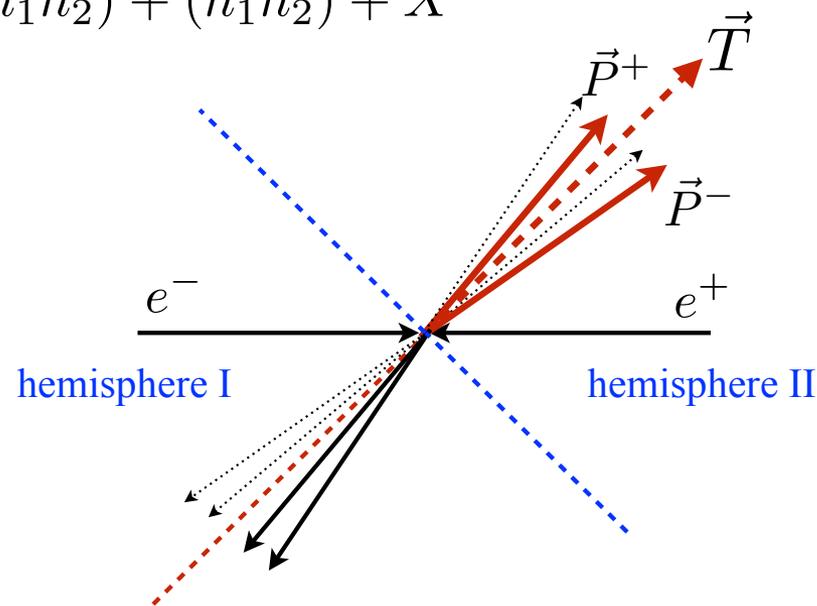
> “left” pairs:

→ $\Omega < 0$

> “right” pairs:

→ $\Omega > 0$

$$\Omega = \frac{(\vec{P}^+ \times \vec{P}^-) \cdot \vec{T}}{|\vec{P}^+| |\vec{P}^-|}$$



The experimental observables for disentangling the hadronization mechanisms :

> jet handedness

$$H = \frac{N_R(\Omega > 0) - N_L(\Omega < 0)}{N_R(\Omega > 0) + N_L(\Omega < 0)}$$

jet handedness correlation

$$C = \frac{\langle \Omega_1 \Omega_2 \rangle - \langle \Omega_1 \rangle \langle \Omega_2 \rangle}{\sqrt{\langle \Omega_1^2 \rangle - \langle \Omega_1 \rangle^2} \sqrt{\langle \Omega_2^2 \rangle - \langle \Omega_2 \rangle^2}}$$

Baryon production

Meson production \approx same colour everywhere.
 Fluctuations with other colour \rightarrow no net force.



i.e. $r + g = \bar{b}$

Baryon production as if diquark
 when only one break
 inside “wrong-colour” region:



Popcorn when several breaks:

