The Silicon Vertex Detector of the Belle II Experiment

Y. Uematsu^q, K. Adamczyk^t, L. Aggarwalⁱ, H. Aihara^q, T. Aziz^j, S. Bacher^t, 2 S. Bahinipati^f, G. Batignani^{k,l}, J. Baudot^e, P. K. Behera^g, S. Bettarini^{k,l} 3 T. Bilka^c, A. Bozek^t, F. Buchsteiner^b, G. Casarosa^{k,l}, L. Corona^{k,l}, T. Czank^p, S. B. Das^h, G. Dujany^e, C. Finck^e, F. Forti^{k,l}, M. Friedl^b, A. Gabrielli^{m,n}, 5 E. Ganiev^{m,n}, B. Gobboⁿ, S. Halder^j, K. Hara^{r,o}, S. Hazra^j, T. Higuchi^p, C. Irmler^b, A. Ishikawa^{r,o}, H. B. Jeon^s, Y. Jin^{m,n}, C. Joo^p, M. Kaleta^t, 7 A. B. Kaliyar^j, J. Kandra^c, K. H. Kang^s, P. Kapusta^t, P. Kodyš^c, T. Kohriki^r, 8 M. Kumar^h, R. Kumarⁱ, C. La Licata^p, K. Lalwani^h, R. Leboucher^d, q S. C. Lee^s, J. Libby^g, L. Martel^e, L. Massacesi^{k,1}, S. N. Mayekar^j, 10 G. B. Mohanty^j, T. Morii^p, K. R. Nakamura^{r,o}, Z. Natkaniec^t, Y. Onuki^q, W. Ostrowicz^t, A. Paladino^{k,l}, E. Paoloni^{k,l}, H. Park^s, L. Polat^d, K. K. Rao^j, 11 12 I. Ripp-Baudot^e, G. Rizzo^{k,1}, D. Sahoo^j, C. Schwanda^b, J. Serrano^d, 13 J. Suzuki^r, S. Tanaka^{r,o}, H. Tanigawa^q, R. Thalmeier^b, R. Tiwari^j, 14 T. Tsuboyama^{r,o}, O. Verbycka^t, L. Vitale^{m,n}, K. Wan^q, Z. Wang^q, J. Webb^a, 15 J. Wiechczynski^l, H. Yin^b, L. Zani^d, 16 (Belle-II SVD Collaboration) 17 ^aSchool of Physics, University of Melbourne, Melbourne, Victoria 3010, Australia 18 ^bInstitute of High Energy Physics, Austrian Academy of Sciences, 1050 Vienna, Austria 19 ^c Faculty of Mathematics and Physics, Charles University, 121 16 Prague, Czech Republic 20 ^dAix Marseille Université, CNRS/IN2P3, CPPM, 13288 Marseille, France 21 ^eIPHC, UMR 7178, Université de Strasbourg, CNRS, 67037 Strasbourg, France 22 ^fIndian Institute of Technology Bhubaneswar, Satya Nagar, India 23 ⁹Indian Institute of Technology Madras, Chennai 600036, India 24 ^hMalaviya National Institute of Technology Jaipur, Jaipur 302017, India 25 ⁱPunjab Agricultural University, Ludhiana 141004, India 26 27 ^jTata Institute of Fundamental Research, Mumbai 400005, India ^kDipartimento di Fisica, Università di Pisa, I-56127 Pisa, Italy 28 ¹INFN Sezione di Pisa, I-56127 Pisa, Italy 29 ^mDipartimento di Fisica, Università di Trieste, I-34127 Trieste, Italy 30 ⁿINFN Sezione di Trieste, I-34127 Trieste, Italy 31 ^o The Graduate University for Advanced Studies (SOKENDAI), Hayama 240-0193, Japan 32 ^pKavli Institute for the Physics and Mathematics of the Universe (WPI), University of 33 Tokyo, Kashiwa 277-8583, Japan 34 ^qDepartment of Physics, University of Tokyo, Tokyo 113-0033, Japan 35 ^rHigh Energy Accelerator Research Organization (KEK), Tsukuba 305-0801, Japan 36 ^sDepartment of Physics, Kyungpook National University, Daegu 41566, Korea 37 ^tH. Niewodniczanski Institute of Nuclear Physics, Krakow 31-342, Poland 38

39 Abstract

The Silicon Vertex Detector (SVD) is a part of the vertex detector in the Belle II experiment at the SuperKEKB collider (KEK, Japan). Since the start of data taking in spring 2019, the SVD has been operating stably and reliably

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with a high signal-to-noise ratio and hit efficiency, achieving good spatial resolu-43 tion and high track reconstruction efficiency. The hit occupancy, which mostly 44 comes from the beam-related background, is currently about 0.5% in the in-45 nermost layer, causing no impact on the SVD performance. In anticipation of 46 the operation at higher luminosity in the next years, two strategies to sustain 47 the tracking performance in future high beam background conditions have been 48 developed and tested on data. One is to reduce the number of signal waveform 49 samples to decrease dead time, data size, and occupancy. The other is to utilize 50 the good hit-time resolution to reject the beam background hits. We also mea-51 sured the radiation effects on the sensor current, strip noise, and full depletion 52 voltage caused during the first two and a half years of operation. The results 53 show no detrimental effect on the SVD performance. 54

55 Keywords: Silicon strip detector, Vertex detector, Tracking detector, Belle II

56 1. Introduction

The Belle II experiment [1] aims to probe new physics beyond the Standard 57 Model in high-luminosity e⁺e⁻ collisions at the SuperKEKB collider (KEK, 58 Japan) [2]. The main collision energy in the center-of-mass system is 10.58 GeV 59 on the $\Upsilon(4S)$ resonance, which enables various physics programs based on the 60 large samples of B mesons, τ leptons, and D mesons. Also, the asymmetric 61 energy of the 7 GeV electron beam and 4 GeV positron beam is adopted for 62 time-dependent CP violation measurements. The target of SuperKEKB is to 63 accumulate an integrated luminosity of 50 ab^{-1} with peak luminosity of about 64 $6 \times 10^{35} \text{ cm}^{-2} \text{s}^{-1}$. In June 2021, SuperKEKB recorded the world's highest 65 instantaneous luminosity of 3.1×10^{34} cm⁻²s⁻¹. The data accumulated before 66 July 2021 corresponds to an integrated luminosity of 213 fb^{-1} . 67

The Vertex Detector (VXD) is the innermost detector in the Belle II detector system. The VXD has six layers: the inner two layers (layers 1 and 2) are the Pixel Detector (PXD), and the outer four layers (layers 3 to 6) are the Silicon Vertex Detector (SVD). The schematic cross-sectional view of the VXD is shown

- ⁷² in Fig. 1. The PXD consists of DEPFET pixel sensors, and its innermost radius
- $_{73}\,$ is 1.4 cm from the beam interaction point (IP). A detailed description of the
- ⁷⁴ SVD appears in Sec. 2.

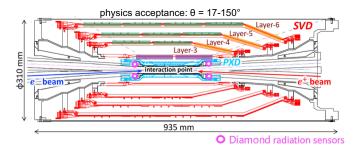


Figure 1: Schematic cross-sectional view of the VXD. The SVD is red, the PXD is light blue, and the IP beam pipe diamonds are pink circles. In the upper half of the VXD the locations of the three types of SVD DSSDs are indicated by boxes in three colors: purple for small sensors, green for large sensors, and orange for trapezoidal sensors as described in Tab. 1.

Diamond sensors [3], used to monitor the radiation dose and for the beam abort system, are mounted on the IP beam pipe and the bellows pipes outside of the VXD. The pink circles in Fig. 1 indicate the locations of the diamond sensors on the IP beam pipe. The diamond's measured doses are used to estimate the dose in the SVD. The diamond system also sends beam abort requests to SuperKEKB if the radiation level gets too high to avoid severe damage to the detector.

82 2. Belle II Silicon Vertex Detector

The SVD is crucial for extrapolating the tracks to the PXD to measure the decay vertices with the PXD and point at a region-of-interest limiting the PXD readout data volume. Other roles of the SVD are the standalone track reconstruction of low-momentum charged particles and their particle identification using ionization energy deposits. The SVD also plays a critical role in the decay vertex measurement in the case of long-lived particles like K_S mesons, which decay inside the SVD volume.

The SVD [4] consists of four layers of double-sided silicon strip detectors 90 (DSSDs). The material budget of the SVD is about 0.7% of a radiation length 91 per layer. On each DSSD plane, a local coordinate is defined with u and v: 92 u-axis along n-side strips and v-axis perpendicular to u-axis. In other words, 93 p-side strips and n-side strips provide u and v information, respectively. In the 94 cylindrical coordinate, u corresponds to $r-\varphi$ information and v corresponds to 95 z information. The SVD consists of three types of sensors: "small" rectangular 96 sensors in layer 3, "large" rectangular sensors in the barrel region of layers 4, 5, 97 and 6, and "trapezoidal" sensors in the forward region of layers 4, 5, and 6, which 98 is slanted. They are indicated by purple, green, and orange boxes in Fig. 1. The 99 main characteristics of these three types of sensors are summarized in Tab. 1. 100 The sensors are manufactured by two companies: the small and large sensors 101 by Hamamatsu and trapezoidal sensors by Micron. The full depletion voltage is 102 60 V for Hamamatsu sensors and 20 V for Micron sensors; both types of sensors 103 are operated at 100 V. In total, 172 sensors are assembled, corresponding to a 104 sensor area of 1.2 m^2 and approximately 224,000 readout strips. 105

	Small	Large	Trapezoidal
No. of u/p-strips	768	768	768
u/p-strip pitch	$50 \ \mu m$	$75~\mu{ m m}$	50–75 μm
No. of v/n-strips	768	512	512
v/n-strip pitch	160 µm	$240~\mu\mathrm{m}$	$240~\mu\mathrm{m}$
Thickness	320 µm	$320~\mu\mathrm{m}$	$300~\mu{\rm m}$
Manufacturer	Hamamatsu		Micron

Table 1: Table of the main characteristics of the three types of sensors. Only readout strips are taken into account for number of strips and strip pitch. All sensors have one intermediate floating strip between two readout strips.

Sensor strips are AC coupled to the front-end ASIC, the APV25 [5], which was originally developed for the CMS Silicon Tracker. The APV25 tolerates more than 100 Mrad of radiation. It has 128 channels with a shaping time of about 50 ns. For the SVD, the APV25 is operated in "multi-peak" mode. The

mechanism of the data sampling in the multi-peak mode is explained in Fig. 2. 110 The chip samples the height of the signal waveform with the 32 MHz clock (31 ns 111 period) and stores each sample's information in an analog ring buffer. Since 112 the bunch-crossing frequency is eight times faster than the sampling clock, the 113 stored samples are not synchronous to the beam collision, in contrast to CMS, 114 which motivates operation in the multi-peak mode. In the present readout 115 configuration (the six-samples mode), at every reception of the Belle II global 116 Level-1 trigger, the chip reads out six successive samples of the signal waveform 117 stored in the buffers. The six-samples mode offers a wide enough time window 118 $(6 \times 31 \text{ ns} = 187 \text{ ns})$ to accommodate large timing shifts of the trigger. In 119 preparation for operation with higher luminosity, where background occupancy, 120 trigger dead-time, and the data size increase, we developed the three/six-mixed 121 acquisition mode (mixed-mode). The mixed-mode is a new method to read 122 out the signal samples from the APV25, in which the number of the samples 123 changes between three and six in each event, depending on the timing precision 124 of each Level-1 trigger signal in that event. For triggers with precise timing, 125 three-samples data are read out and the data have half time window and half 126 data size compared to ones of six-samples data, resulting in the reduction of the 127 effects due to higher luminosity. This functionality was already implemented 128 in the running system and confirmed by a few hours of smooth physics data 129 taking. Before we start to use the mixed-mode, the effect on the performance 130 due to the change of the acquisition mode is to be assessed. As the first step, 131 the effect in the hit efficiency was evaluated as described in Sec. 3. 132

The APV25 chips are mounted on each middle sensor (chip-on-sensor con-133 cept) with thermal isolation foam in between. The merit of this concept is 134 shorter signal propagation length, leading to smaller capacitance of the signal 135 line and hence reduced noise level. To minimize the material budget the APV25 136 chips on the sensor are thinned down to $100 \ \mu m$. The APV25 chips are mounted 137 on a single side of the sensor and readout of the signals from the opposite side is 138 performed via wrapped flexible printed circuits. The power consumption of the 130 APV25 chip is 0.4 W/chip and 700 W in the entire SVD. The chips are cooled 140

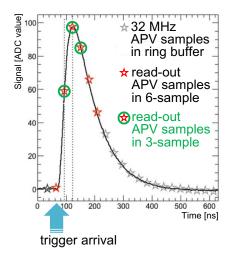


Figure 2: Example of sampling in "multi-peak" mode of the APV25. The black line shows the signal waveform after the CR-RC shaper circuit. The stars show the sampled signal height recorded in the analog ring buffer according to the 32 MHz sampling clock. The red stars indicate the six successive samples read out at the trigger reception in the six-samples mode. The red stars with a green circle indicate the samples read out in the three-samples acquisition.

¹⁴¹ by a bi-phase -20°C CO₂ evaporative cooling system.

¹⁴² 3. Performance

Since March 2019, the SVD has been operating reliably and smoothly for two and a half years. The total fraction of masked strips is about 1%. There was only one issue where one APV25 chip (out of 1,748 chips) was disabled during the spring of 2019, which was remediated by reconnecting a cable in the summer of 2019.

The SVD has also demonstrated stable and excellent performance [6]. The hit efficiency is continuously over 99% in most of the sensors. The cluster charge distributions are also reasonable. On the u/p-side, the most probable values agree with the calculated charge amount induced by MIPs within the uncertainty in calibration. On the v/n-side, 10–30% of the collected charge is lost compared to the signal collected on the u/p-side, due to the presence of the floating strip combined with the large pitch on the v/n-side. The most probable
values of the cluster signal-to-noise ratio distributions range from 13 to 30.

We measured the cluster position resolution by analyzing the $e^+e^- \rightarrow \mu^+\mu^-$ 156 data [7]. The resolution is estimated from the residual between the cluster po-157 sition and the track position, not biased by the target cluster, after subtracting 158 the effect of the track extrapolation error. The cluster position resolutions for 159 different incident angles are shown in Fig. 3. The observed resolution has the 160 expected shape, showing a minimum at the incident angle for which the projec-161 tion of the track along the direction perpendicular to the strips on the detector 162 plane corresponds to two strip pitches. Given the various sensor pitches with 163 one floating strip, the minimum is expected at 14 (21) degrees on the v/n-side 164 and at 4 (7) degrees on the u/p-side, respectively for layer 3 (4, 5, and 6). The 165 resolution for normal incident angle is also in good agreement with the expected 166 digital resolution, that is 23 (35) μ m on the v/n-side, 7 (11) μ m on the u/p-side, 167 respectively for layer 3 (4, 5, and 6). Still, some studies are ongoing to improve 168 the resolution especially for the layer-3 u/p-side, where at normal incidence a 169 slightly higher resolution is measured $(9 \ \mu m)$ compared to the expectations. 170

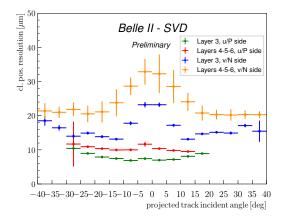


Figure 3: The SVD cluster position resolution depending on the projected track incident angle. The green (blue) plot shows the resolution in the u/p-side (n/v-side) of layer-3 sensors, and the red (yellow) one shows the u/p-side (n/v-side) of layers-4, 5, and 6 sensors.

¹⁷¹ The cluster hit-time resolution was also evaluated in candidate hadronic

events¹ using the reference event time estimated by the Central Drift Chamber 172 (CDC) outside of the SVD. The error on the event time, about 0.7 ns, was 173 subtracted to evaluate the intrinsic SVD hit-time resolution. The resulting 174 resolution is 2.9 ns on the u/p-side and 2.4 ns on the v/n-side. With such 175 precise hit-time information, it is possible to reject off-time background hits 176 efficiently. The hit-time distributions for signal² and background³ are shown 177 in Fig. 4. The signal distribution has a narrow peak, while the background 178 hit-time distribution is broad and almost flat in the signal peak region. The 179 separation power of the hit-time is high, as expected. For example, if we reject 180 hits with the hit-time less than -38 ns in this plot, we can reject 45% of the 181 background hits while keeping 99% of the signal hits. The background rejection 182 based on the hit-time is essential to sustain the good tracking performance in 183 the future high beam background condition. 184

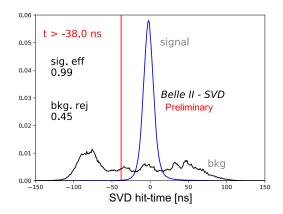


Figure 4: Example of the background hit rejection using hit-time. The blue distribution shows the signal, and the black distribution shows the background. Assuming the hit-time cut at -38 ns, the signal hit efficiency of 99% and the background hit rejection of 45% are achieved.

185 186 The performance in three-samples data was compared with that in sixsamples data to evaluate the performance in the mixed-mode. If the trigger

¹The events with more than three good tracks and not like Bhabha scattering.

 $^{^{2}}$ The clusters found to be used in the tracks in the hadronic events.

³The clusters in events triggered by delayed-Bhabha pseudo-random trigger.

timing has no deviation, the three-samples data will show comparable perfor-187 mance to the six-samples data because the relevant part of the signal waveform 188 to evaluate the necessary signal properties, i.e., the signal height and the sig-189 nal timing can be accommodated in the three-sample's time window. However, 190 when the trigger has a jitter and the timing shift happens, some part of the 191 signal waveform can be out of the three-sample's time window, and the recon-192 struction performance deteriorates. We examined the effect on the hit efficiency 193 as a function of the trigger timing shift. The effect is evaluated by the rel-194 ative hit efficiency, which is defined as the ratio of the hit efficiency in the 195 three-samples data to the one in the six-samples data. For this study, the three-196 samples data are emulated in the offline analysis from the six-samples data by 197 selecting consecutive three samples at a fixed latency with respect to the Level-198 1 trigger signal. The trigger timing shift is evaluated by the CDC event time. 199 The resulting relative efficiencies as a function of the trigger timing shift in the 200 hadronic events are shown in Fig. 5. The decreasing trend is observed for the 201 shift of the trigger timing, as expected. As a result, the relative efficiency is 202 over 99.9% for the trigger timing shift within ± 30 ns, which is almost all the 203 events. 204

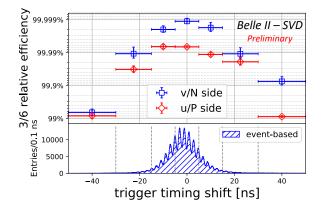


Figure 5: The relative hit efficiencies (the ratios of the hit efficiency in the three-samples data to the one in the six-samples data) as a function of the trigger timing shift for v/n-side (blue square) and u/p-side (red diamond). The positive (negative) trigger timing shift corresponds to early (late) trigger timing.

²⁰⁵ 4. Beam-related background effects on SVD

The beam-related background (BG) increases the hit occupancy of the SVD, 206 which in turn degrades the tracking performance. Considering this performance 207 degradation, we set the occupancy limit in layer-3 sensors to be about 3%, which 208 will be loosened roughly by a factor of two after we apply the hit-time rejection 209 described in Sec. 3. With the current luminosity, the average hit occupancy 210 in layer-3 sensors is below 0.5%. However, the projection of the hit occupancy 211 at the luminosity of 8×10^{35} cm⁻²s⁻¹ is about 3% in layer-3 sensors. The 212 projected occupancy comes from the Monte Carlo (MC) simulation scaled by 213 the data/MC ratio determined from the BG data of the current beam optics. 214 The corresponding integrated dose, using the data/MC-rescaled BG extrapo-215 lation, is about 0.2 Mrad/smy, and the equivalent 1-MeV neutron fluence is 216 about $5 \times 10^{11} n_{eq}/cm^2/smy$ (smy: Snowmass Year = 10^7 sec). Considering the 217 radiation hardness of the SVD sensors, about 10 Mrad and about $10^{13} n_{eq}/cm^2$, 218 based on the experience of similar DSSD sensors used in the BaBar Silicon Ver-219 tex Tracker [8], we expect to be able to safely operate the SVD even for ten 220 years at high luminosity, with safety margin of factor two to three against BG 221 extrapolation. The long-term BG extrapolation is affected by large uncertain-222 ties from the optimization of collimator settings in MC and the future evolution 223 of the beam injection background, which is not simulated. This uncertainty, 224 together with the relatively small safety factor of two to three between the BG 225 extrapolation and the detector limits, motivates the VXD upgrade to improve 226 the tolerance of the hit rates and the radiation damage, and the technology 227 assessment is ongoing for multiple sensor options. 228

In the first two and a half years of operation, the integrated dose in the layer-3 mid-plane sensors, which are the most exposed in the SVD, is estimated to be 70 krad. The estimation is based on the measured dose by the diamonds on the beam pipe exploiting the measured correlation between the SVD occupancy and the diamond dose [9]. Thanks to a new random trigger line recently introduced, we improved the dose analysis, removing an overestimation of about factor three in the previous study. The new estimate still has an uncertainty of about 50% mainly due to the unavailability of this newly introduced trigger before December 2020. Assuming the dose/n_{eq} fluence ratio of $2.3 \times 10^9 \text{ n}_{eq}/\text{cm}^2/\text{krad}$ from MC, 1-MeV equivalent neutron fluence is evaluated to be about $1.6 \times 10^{11} \text{ n}_{eq}/\text{cm}^2$.

The effect of the integrated dose on the sensor leakage current is measured, 240 and the results show a clear linear correlation as in the upper plot of Fig. 6. The 241 slopes for all the sensors are $2-5 \ \mu A/cm^2/Mrad$, as summarized in the lower 242 plot of Fig. 6. The large variations can be explained by temperature effects and 243 the deviation of sensor-by-sensor dose from the average in each layer used in 244 the estimation. The slopes are in the same order of magnitude as previously 245 measured in the BaBar experiment [8], 1 μ A/cm²/Mrad at 20°C. The precise 246 temperature in layer 3 of the SVD is unknown but expected to be in a similar 247 regime. While the leakage current is increasing, the impact on the strip noise 248 is suppressed by the short shaping time (50 ns) in APV25. It is expected to be 249 comparable to the strip-capacitive noise only after 10 Mrad irradiation and not 250 problematic for ten years where the integrated dose is estimated to be 2 Mrad. 251 The evolution of the noise with the integrated dose is shown in Fig. 7. The 252 noise increase of 20–25% is observed in layer 3, but this does not affect the 253 SVD performance. This noise increase is likely due to the radiation effects on 254 the sensor surface. Fixed oxide charges on sensor surface increase with dose, 255 with some saturation expected at around 100 krad, enlarging also non-linearly 256 the inter-strip capacitance, also expected to saturate with dose. The noise 257 saturation is already observed on the v/n-side and also starts to be seen on the 258 u/p-side. 259

The full depletion voltage of the sensor is also a key property that can be affected by the radiation damage. It can be measured from the v/n-side strip noise, which suddenly decreases at the full depletion voltage because the sensor substrate is n-type and thus the v/n-side strips are only fully isolated at full depletion. From this measurement full depletion voltages consistent with measurements performed on the bare sensors before the installation were obtained, ranging from 20 to 60 V, and so far no change in full depletion voltage is observed in the first two and a half years of operation, which is consistent with the expectation from low integrated neutron fluence of $1.6 \times 10^{11} \text{ n}_{eq}/\text{cm}^2$.

269 5. Conclusions

The SVD has been taking data in Belle II since March 2019 smoothly and 270 reliably. The detector performance is excellent and agrees with expectations. 271 We are ready to cope with the increased background during higher luminosity 272 running by rejecting the off-time background hits using hit-time and operating 273 in the three/six-mixed acquisition mode. In the recent study, the efficiency 274 loss in the three-samples data is confirmed to be less than 0.1% for the trigger 275 timing shift within ± 30 ns. The observed first effects of radiation damage are 276 also within expectation and do not affect the detector performance. 277

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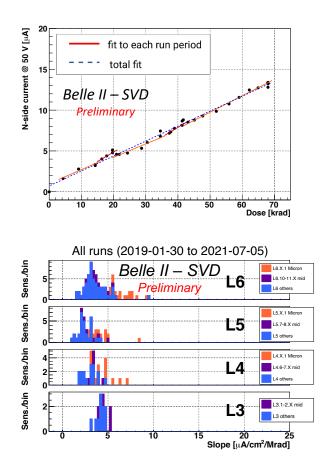


Figure 6: (upper) Effect of the integrated dose on the leakage current in the n/v-side of one layer-3 sensor. The slope is fitted for each run period (solid red line) and all the runs (dashed blue line). Both fit results agree with each other and are consistent with the linear increase. (lower) The fit results of all the sensors for all runs. The sensors are classified as trapezoidal sensors in the forward region (Micron), sensors around the midplane, and the others.

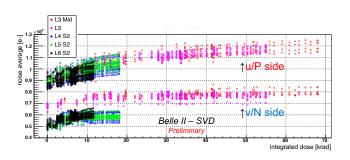


Figure 7: Effect of the integrated dose on the noise average in electron. The upper (lower) series shows the u/p-side (v/n-side) results, respectively.