

# Lepton Identification using Belle II Silicon-strip Vertex Detector

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**Abstract.** We improve the identification performance of low-momentum leptons, especially electrons, using the specific ionisation information from the silicon-strip vertex detector (SVD) of the Belle II experiment.

**Keywords:** Belle II, SVD, PID, Lepton identification

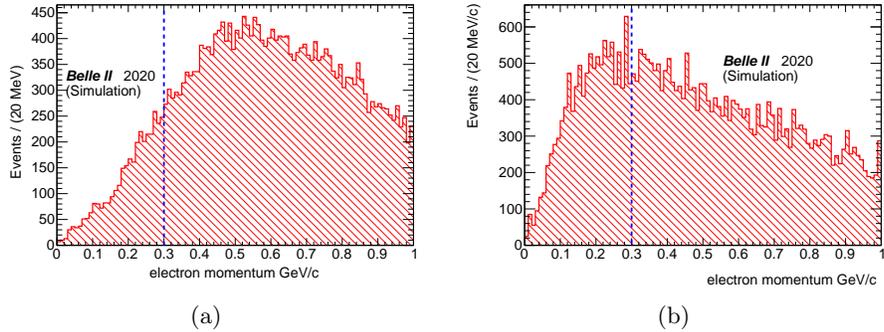
## 1 Introduction

Particle identification (PID) serves a crucial role in any flavor physics experiment. At Belle II [1], we use information from various subdetectors to identify a track as a lepton [2]. The PID algorithm heavily relies on the information provided by the electromagnetic calorimeter (ECL) to identify electrons and the  $K_L^0$  and muon detector (KLM) to identify muons. The performance deteriorates for low-momentum leptons that fail to reach the respective subdetectors and are instead reconstructed in the tracking system. We aim to improve the lepton identification performance for such tracks using the specific ionisation information from the SVD. Our study will aid the physics studies involving low-momentum leptons such as decays of  $B$  mesons to semileptonic final states:  $B \rightarrow K^{(*)}\ell\ell$ ,  $B \rightarrow K^{(*)}\tau\tau$  and lepton-flavor-violating decays of  $\tau \rightarrow e\ell\ell$  and  $\tau \rightarrow \mu\ell\ell$ , where  $\ell$  is  $e$  or  $\mu$ . Figure 1 shows the momentum distribution of electrons from two such decays, where a good number of electron candidates are found to have a momentum below the threshold to reach the ECL.

## 2 Electron identification using SVD

### 2.1 Event selection

The study is performed using Monte Carlo (MC) simulated data for  $e^+e^-$  collisions recorded at the  $\Upsilon(4S)$  peak with the Belle II detector. We use electrons coming from photon conversion ( $\gamma \rightarrow e^+e^-$ ) due to interaction with the two inner tracking systems including SVD. Electrons are reconstructed as primary particles using track-level information from the detector. A converted photon candidate is reconstructed by combining two oppositely charged tracks. We apply various kinematic and vertex requirements to suppress background events.



**Fig. 1.** Momentum distributions of slow electrons coming from  $B \rightarrow K^{(*)}\tau\tau$  (left) and  $\tau \rightarrow eee$  (right) decays. The dashed vertical lines indicate the momentum threshold required to reach the ECL.

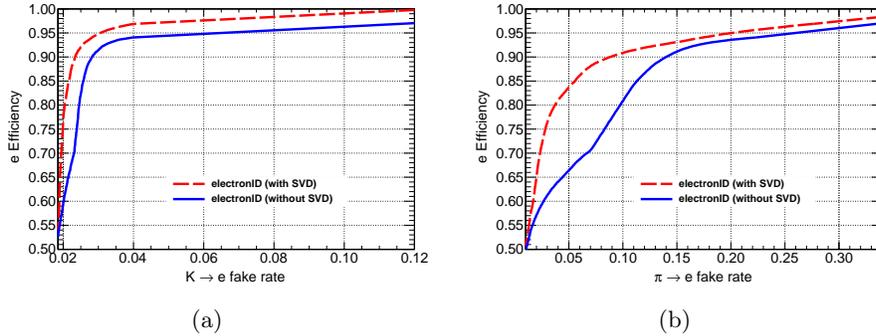
We use the  $s\mathcal{P}$ lot [3] technique to subtract the residual background. The  $s\mathcal{P}$ lot extracted two-dimensional distribution of energy loss ( $dE/dx$ ) vs. momentum is used as the contribution of SVD to the total PID for electron tracks.

## 2.2 PID performance

We compare the PID performance evaluated with and without the SVD information. To calculate the electron identification efficiency, we use the  $\gamma \rightarrow e^+e^-$  sample, while for calculating  $\pi \rightarrow e$  and  $K \rightarrow e$  fake rates we use the  $D^{*+} \rightarrow D^0[\rightarrow K^-\pi^+]\pi^+$  sample. The latter was used earlier for hadron identification studies [4]. We define:

$$\text{efficiency } (\epsilon_i) = \frac{\text{No. of tracks identified with PID under the hypothesis } i}{\text{No. of tracks kinematically identified under the hypothesis } i}$$

$$\text{fake rate } (f_{j \rightarrow i}) = \frac{\text{No. of tracks identified with PID under the hypothesis } i}{\text{No. of tracks kinematically identified under the hypothesis } j}$$



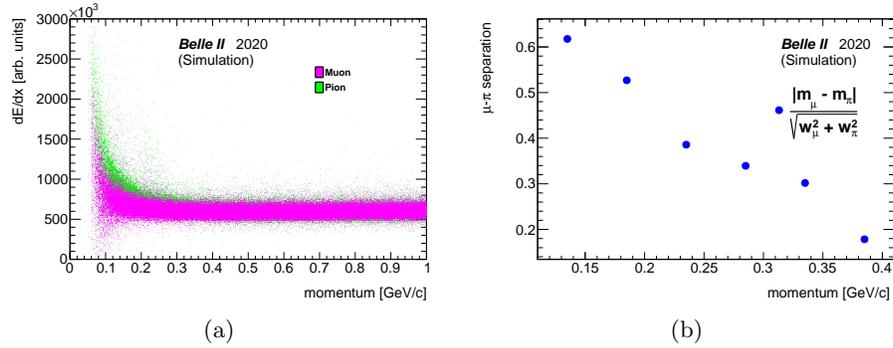
**Fig. 2.** Efficiency vs fake rate plots for different criteria on PID variable.

Figure 2 shows that the performance improves significantly with the introduction of SVD information in the PID. For a fixed 10%  $\pi \rightarrow e$  fake rate in MC, the electron efficiency increases from around 80% to 90%. Similarly, for a fixed 4%  $K \rightarrow e$  fake rate in MC, the electron efficiency increases from around 94% to 97%.

### 3 Muon identification using SVD

The performance of SVD to separate  $dE/dx$  distributions of muons and pions is studied using simulated charged tracks. We generate two million muon and pion tracks and let them pass through the detector simulation. Figure 3(a) shows the distribution of  $dE/dx$  vs. momentum for these tracks. We obtain the  $dE/dx$  distribution for muons and pions in momentum bins of 50 MeV/c. The distributions are fitted with a Gaussian function to obtain their mean (m) and width (w) values. We then look at the distribution of  $\mu$ - $\pi$  separation vs. momentum as shown in Fig. 3(b). The  $\mu$ - $\pi$  separation is defined as

$$\frac{|m_{\mu} - m_{\pi}|}{\sqrt{w_{\mu}^2 + w_{\pi}^2}}$$



**Fig. 3.**  $dE/dx$  vs. momentum distributions for muons and pions (left), and the  $\mu$ - $\pi$  separation vs. momentum (right).

We observe that muon and pion tracks are not so well separated in the SVD owing to their small mass difference (around  $30 \text{ MeV}/c^2$ ). There is a maximum separation of around 0.6 standard deviations between the two hypotheses. This means we can use the same  $dE/dx$  vs. momentum distribution as the SVD information to the PID for both pions and muons, and we have decided to use the one obtained for pions.

## References

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