

¹³ Carlo sample. The introduction of additional information from the SVD is found to improve the

¹⁴ overall PID performance in the low momentum region.

15 1. INTRODUCTION

Particle identification (PID) plays a central role in the physics program of Belle II. Low-16 momentum charged particles having a transverse momentum $p_{\rm T} \lesssim 65 \,\mathrm{MeV}/c$ are unable to 17 reach the central drift chamber (CDC), the main tracking system of the experiment, owing 18 to their highly curved trajectories. Our goal is to exploit specific ionization (dE/dx) by 19 these low-momentum particles in the silicon-strip vertex detector (SVD) towards identifying 20 them. Even if the particles have a $p_{\rm T}$ greater than $65 \,{\rm MeV}/c$ and are thus able to reach the 21 CDC, the dE/dx values measured in the SVD can provide complementary information to 22 that obtained from the main PID detectors of Belle II: CDC, TOP, and ARICH. 23

We use clean samples of $D^{*+} \to D^0(K^-\pi^+)\pi^+$ and $\Lambda \to p\pi$ decays to first obtain the SVD dE/dx calibration for charged pions, kaons and protons. Later we check the impact of dE/dx information on overall PID performance using the same decay channels. We have also verified the dE/dx values to be independent of the mass of traversing particles depending only on their $\beta\gamma$ values. The study is based on e^+e^- collision data recorded at the $\Upsilon(4S)$ peak by Belle II and the results are compared with that of a Monte Carlo (MC) sample.

To assess the impact of SVD dE/dx information on the overall PID performance, we plot the identification efficiency and fake rate as a function of momentum applying a requirement on the binary PID likelihood $\mathcal{L}(i/j) > 0.5$. The efficiency is defined as:

$$\epsilon_i = \frac{\text{\# charged particle tracks identified kinematically as well as with PID requirement under the hypothesis i}{\text{\# charged particle tracks identified kinematically under the hypothesis i}}$$

³⁴ and the fake rate is given by:

 $f_{j \to i} = \frac{\# \text{ charged particle tracks identified kinematically as well as with PID requirement under the hypothesis }{\# \text{ charged particle tracks identified kinematically under the hypothesis }_{j}}$

36 2. RESULTS

³⁷ In the following, we will have a series of plots describing various aspects of the study.

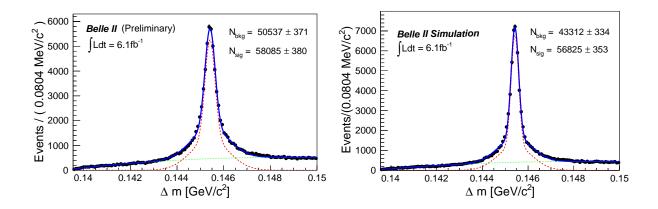


FIG. 1. Fitted distributions of D^*-D^0 mass difference (Δm) from the D^* sample in data (left) and MC (right) events. We require the charged particle tracks to have a transverse (longitudinal) impact parameter less than 0.5 cm (2.0 cm). These tracks must have at least one SVD hit and a track-fit χ^2 value greater than 10^{-5} . To further purify the sample, we require the reconstructed D^0 mass to lie between 1.85 and 1.88 GeV/ c^2 , corresponding to a $\pm 3\sigma$ window around the nominal D^0 mass. The reconstructed D^* mass must be within 1.95 and 2.05 GeV/ c^2 . We apply a loose criterion on kaon and pion PID likelihoods, calculated without SVD information, to remove low-momentum secondary pions and kaons produced due to hadronic interaction in the detector material. We model the signal and background Δm shape by a sum of two Gaussian functions with a common mean and a threshold function, respectively. The ${}_s\mathcal{P}lot$ [1] technique is used to subtract the residual background contributions.

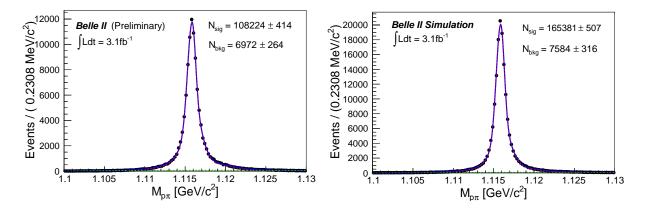


FIG. 2. Fitted $M_{p\pi}$ distributions from the Λ sample in data (left) and MC (right) events. We require the reconstructed $p\pi$ invariant mass of Λ candidates to be in the range [1.10, 1.13] GeV/ c^2 , and they are further subjected to a vertex fit. The distance between the interaction point and the vertex of the Λ candidates is required to be greater than 1 cm and the vertex fit χ^2 must be greater than 10^{-3} to remove the random combination of two tracks. We also require at least one SVD hit for both daughter tracks of Λ candidates. We suppress the contamination of charged pions coming from the K_S^0 decay by rejecting events that have the $M_{\pi^+\pi^-}$ value in the range [488, 508] MeV/ c^2 , corresponding to a $\pm 3\sigma$ window around the nominal K_S^0 mass. Similarly, events with electrons from converted photons are suppressed by excluding $M_{e^+e^-} < 50 \,\text{MeV}/c^2$, a veto decided by the figure-of-merit optimization. We impose an additional requirement of at least one CDC hit and a loose criterion on the proton PID calculated without SVD information to remove low-momentum secondary pions produced due to hadronic interaction with the detector material. We model the signal shape with the sum of a Gaussian and two asymmetric Gaussian functions of a common mean and the background shape with a second-order Chebyshev polynomial. The $_s \mathcal{P}lot$ [1] technique is used to subtract the residual background contributions.

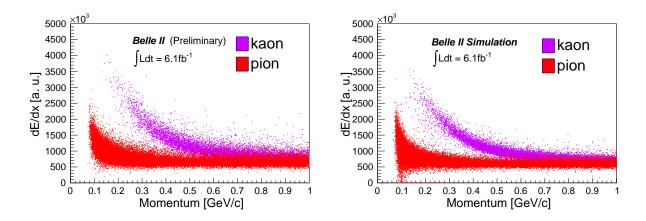


FIG. 3. Scatter plot of dE/dx values of charged pions and kaons as a function of their momentum for data (left) and MC (right) events from the D^* sample. The two-dimensional distributions of dE/dx vs. momentum show a clear separation between different particles in the low momentum region, and are uploaded to the calibration database.

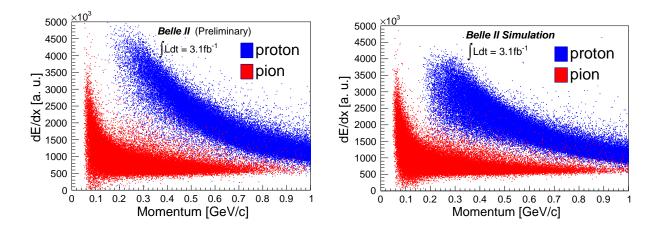


FIG. 4. Scatter plot of dE/dx of proton and pion as a function of their momentum in data (left) and MC (right) events from the Λ sample. The two-dimensional distributions of dE/dx vs. momentum show a clear separation between different particles in the low momentum region, and are uploaded to the calibration database.

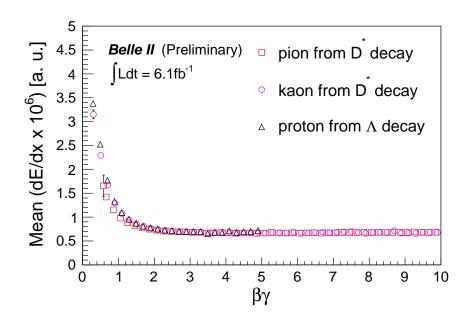


FIG. 5. The $\beta\gamma$ universality curve obtained in data with D^* and Λ samples. As the specific ionisation depends only on the velocity (β) of traversing particles, we check the $\beta\gamma$ universality of dE/dx values obtained from D^* and Λ samples. The minimum energy loss occurs at $\beta\gamma \approx 3$ regardless of the particle type. We get a flat curve beyond that threshold, as the relativistic rise of energy loss is suppressed by the density effect in silicon.

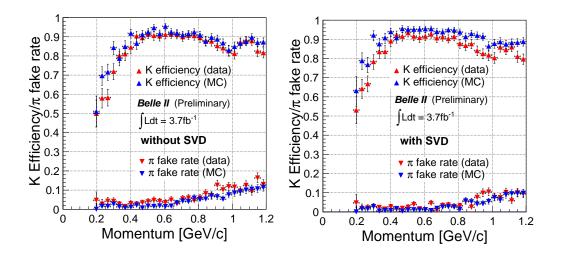


FIG. 6. K efficiency and π fake rate vs. momentum without (left) and with (right) SVD dE/dx information. To assess the impact of this information to the overall PID, we plot the efficiency and fake rate as a function of momentum applying a requirement on the PID likelihood $\mathcal{L}(K/\pi) > 0.5$. The study shows an improvement in kaon efficiency after adding SVD dE/dx information.

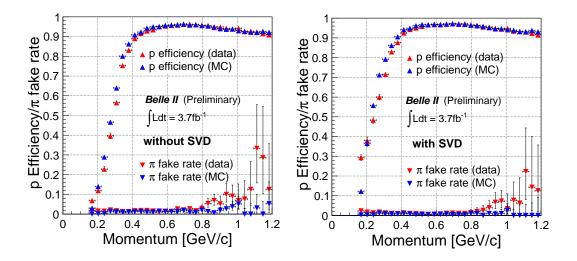


FIG. 7. p efficiency and π fake rate vs. momentum without (left) and with (right) SVD dE/dx information. To assess the impact of this information to the overall PID, we plot the efficiency and fake rate as a function of momentum applying a requirement on the PID likelihood $\mathcal{L}(p/\pi) > 0.6$. The study shows an improvement in proton efficiency after adding SVD dE/dx information.

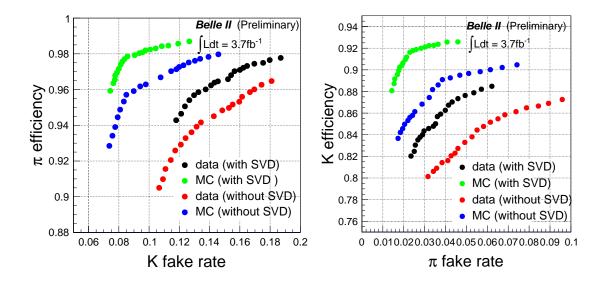


FIG. 8. π efficiency vs. K fake rate (left) and K efficiency vs. π fake rate (right) with and without SVD for p < 1 GeV/c. We plot the efficiency vs. fake rate to better appreciate the improvement in PID performance by adding the SVD dE/dx information. The PID likelihood criterion is varied from 0 to 1 in order to produce these plots. Our study confirms that addition of dE/dx information improves the pion (kaon) efficiency for a given kaon (pion) fake rate in the low momentum region.

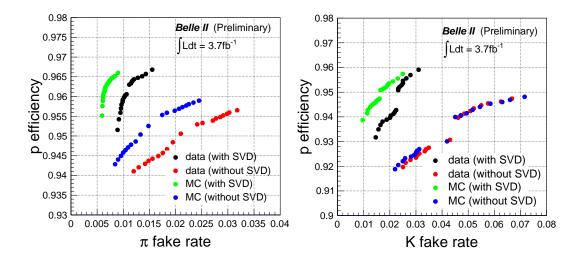


FIG. 9. p efficiency vs. π fake rate (left) and p efficiency vs. K fake rate (right) with and without SVD for p < 1 GeV/c. We plot the efficiency vs. fake rate to better appreciate the improvement in PID performance by adding the SVD dE/dx information. The PID likelihood criterion is varied from 0 to 1 in order to produce these plots. Our study confirms that addition of dE/dx information improves the proton efficiency for a given paon or kaon fake rate in the low momentum region.

 $_{\tt 38}$ [1] M. Pivk and F. R. Le Diberder, Nucl. Instrum. Meth. A555, 356 (2005).