



Charm and time-dependent CPV in B decays at Belle II*

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Abstract

The Belle II measurements relying on the precise determination of the positions of displaced vertices are presented. It is the analysis of the $B^0 - \bar{B}^0$ mixing and CP violation based on the early Belle II data and the D^0 and D^+ lifetime measurement, where the world-leading accuracy was achieved already with the analysed data set. Lifetime measurements of other charm mesons are in the pipeline.

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9 1. INTRODUCTION

10 Excellent performance of the Belle II tracking system [1] allows for precise measurement
11 of the position of the displaced vertices, i.e. vertices not corresponding to the primary
12 e^+e^- interaction. Reconstruction of these vertices is essential for the lifetime measurements,
13 measurement of the $B^0-\bar{B}^0$ oscillations as well as for the CP violation parameters measured
14 in the time-dependent analysis of the B meson decays [2].

15 Compared to the LHCb, at Belle II the momenta of the produced particles are typically
16 smaller and so are the flight distances. On the other hand, at the e^+e^- B -factories the
17 event topology is much clearer and the event reconstruction profits from the knowledge of
18 the energies of the beams, in contrast to the pp collisions, where partons entering the hard
19 interaction carries only (unknown) fraction of the proton's momentum.

20 In this text a prompt measurement of the B^0 meson oscillation frequency Δm_d is pre-
21 sented, this analysis was recently superseded by [3] with 6times higher luminosity. Measuring
22 the oscillation frequency paves a highway toward high-precision measurements of the CKM
23 angle $\sin 2\phi_1$ related to the CP violation, where Belle has achieved the world-leading accu-
24 racy of 3% [4] and the full Belle II data set should improve it to 1% precision [5]. In this
25 text, the $\sin 2\phi_1$ measurement on the early Belle II data is shown as a demonstration of
26 Belle II capabilities. Measurement on data with higher statistics will follow.

27 The Belle II is not only a B -factory as the $e^+e^- \rightarrow c\bar{c}$ process has a similar cross section.
28 Especially, it allowed measuring the D^+ and D^0 lifetimes with world-leading accuracy already
29 with the current data set.

30 2. BELLE II EXPERIMENT

31 The Belle II experiment [6] is a B meson factory located in Tsukuba, Japan based on the
32 SuperKEKB accelerator complex. Electrons and positrons are collided at a centre-of-mass
33 (CM) energy at or near the mass of the $\Upsilon(4S)$ resonance, i.e. 10.58 GeV. The $\Upsilon(4S)$ decays
34 almost instantly into a $B\bar{B}$ pair. Belle II is an upgrade of its predecessor Belle [7], with a
35 target integrated luminosity of 50 ab^{-1} which is 50times the size of the Belle data set. Data
36 taking at Belle II has started in March 2019 and in April 2022 a total integrated luminosity
37 of 314 fb^{-1} has been reached. The data taking will continue until Summer 2022 when a
38 shut-down is planned.

39 3. $B^0-\bar{B}^0$ MIXING AND THE CP VIOLATION MEASUREMENT

40 Precise measurement of the $B^0-\bar{B}^0$ mixing frequency Δm_d and/or the B^0 -meson lifetime
41 is essential to validate the tools and techniques used for the time-dependent studies of the
42 $B^0\bar{B}^0$ system. In particular, proving that Belle II detector can measure these quantities
43 accurately ensures that it can also be used to perform a precise measurement of the CKM
44 angle ϕ_1 in the near future. The ϕ_1 measurement uses very similar experimental techniques
45 but with a data set approximately 10times smaller [8] than the one used for the Δm_d
46 determination. Therefore, the presented analysis allows probing the framework which a
47 much higher level of scrutiny.

48 In contrast to the ϕ_1 where Belle I and LHCb achieved similar accuracy, the world average
49 Δm_d value, $\Delta m_d = 0.5065 \pm 0.0019 \text{ ps}^{-1}$ is dominated by LHCb [9] measurement.

At Belle II the Δm_d or $\sin 2\phi_1$ are determined using the $e^+e^- \rightarrow \Upsilon(4S) \rightarrow B^0\bar{B}^0$ events. As the e^+ and e^- beam energies are different, the $\Upsilon(4S)$ is produced with a Lorentz boost $\beta\gamma \approx 0.288$. The B^0 mesons are produced almost at the rest in the $\Upsilon(4S)$ reference frame. One B^0 meson, B_{sig}^0 is fully reconstructed either as the flavor specific final state (mixing analysis) or as the CP symmetric final state (ϕ_1 analysis). The other B^0 meson, B_{tag}^0 , is reconstructed by combining all other tracks in the event not used for the B_{sig}^0 reconstruction. The flavor of the B_{tag}^0 meson is determined using the category-based flavor tagging algorithm [10]. Neglecting the detector effects on the decay time measurement, the time evolution of the $B_{\text{sig}}^0 B_{\text{tag}}^0$ system is for the mixing measurement described as:

$$f_{\text{phys}}^{\text{mixing}}(\Delta t) = \frac{e^{-|\Delta t|\tau_{B^0}}}{4\tau_{B^0}}(1 + q \cos \Delta m_d \Delta t), \quad (1)$$

where $q = 1$ (-1) for events where B_{sig}^0 and B_{tag}^0 have opposite (same) flavor at their respective decay time. For B_{sig}^0 corresponding to the CP eigen-state, the expression is similar:

$$f_{\text{phys}}^{\text{CP}}(\Delta t) = \frac{e^{-|\Delta t|\tau_{B^0}}}{4\tau_{B^0}}(1 + qS \sin \Delta m_d \Delta t), \quad (2)$$

50 where $S = \sin 2\phi_1$ is related to the time-dependent CP asymmetry, whereas the direct
 51 CP asymmetry is assumed to be zero. Flavor of B_{tag}^0 is described by q , $q = 1$ ($q = -1$)
 52 corresponds to particle (anti-particle).

53 The proper-time difference between the B_{sig}^0 and B_{tag}^0 decays, Δt , is reconstructed from
 54 the difference between the two B meson vertex position vectors projected into the direction
 55 of $\Upsilon(4S)$ momentum Δl , i.e. $\Delta t' = \Delta l / (\beta\gamma c)$.

56 The reconstructed $\Delta t'$ differs from the truth Δt due to the detector resolution effects,
 57 the efficiency of the flavor tagger, as well as the residual movement of the B -mesons in the
 58 $\Upsilon(4S)$ reference frame. They are described by the detector resolution function which is the
 59 probability density of the difference between the truth and the reconstructed Δt . A perfect
 60 understanding of the Δt resolution is essential for precise measurement of the CKM angle
 61 ϕ_1 . The Δt resolution is dominated by the reconstruction of the B_{tag}^0 vertex, while the fully
 62 reconstructed signal B^0 vertex is determined with higher precision. The resolution function
 63 obtained from the mixing analysis is also of use in the CPV measurement.

64 In the analysis of the prompt 2020 Belle II data [11] with integrated luminosity 34.6 fb^{-1} ,
 65 the parameters of the resolution function are kept identical for both mixing and CPV mea-
 66 surement. The flavor specific B -meson decay is reconstructed in $B^0 \rightarrow D^-(K^+\pi^-\pi^-)\pi^+$ and
 67 the B -meson in CP eigen-state is reconstructed via $B_{\text{CP}}^0 \rightarrow J/\psi(\ell\ell)K_s(\pi^+\pi^-)$ decay. The
 68 measured spectra of the Δt together with the fitted curve are shown in Figure 1. Several se-
 69 lection criteria are applied to suppress the background, for example on the beam-constrained
 70 invariant mass of the B_{sig}^0 meson or on the reduced Fox-Wolfram moment R_2 to suppress
 71 the continuum background which typically has jet-like topology.

72 The obtained values $\Delta m_d = 0.531 \pm 0.046 \pm 0.013 \text{ ps}^{-1}$ and $\sin 2\phi_1 = 0.55 \pm 0.21 \pm 0.04$ are
 73 in agreement with the world average values and their precision is driven by the statistical
 74 uncertainty. It is also consistent with recent Belle II measurement [3], $\Delta m_d = 0.516 \pm$
 75 $0.008 \pm 0.005 \text{ ps}^{-1}$, based on data with $L_{\text{int}} = 190 \text{ fb}^{-1}$. It demonstrates that the employed
 76 analysis workflow can be used for the state-of-the-art ϕ_1 measurement with emerging full
 77 Belle II data set.

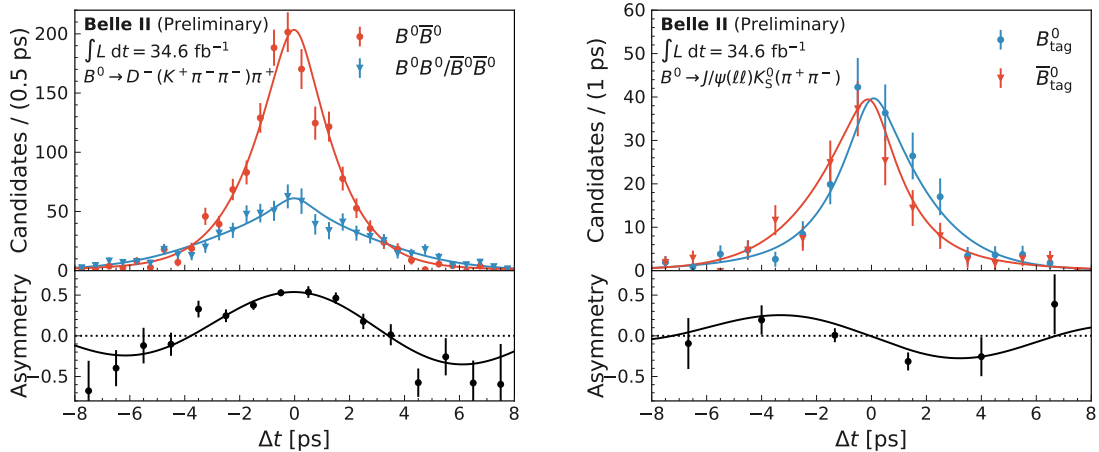


FIG. 1: The Δt distribution for the $B^0 \rightarrow D^-(K^+\pi^-\pi^-)\pi^+$ flavor-specific state (left) and for $B^0 \rightarrow J/\psi(\ell\ell)K_s^0(\pi^+\pi^-)$ CP eigen-state (right) [11]. The data points and the fitted curves are plotted for both flavor states of the B_{tag}^0 .

78 4. D LIFETIME MEASUREMENT

79 Accurate predictions of lifetimes of weakly decaying charmed and bottom hadrons are
80 challenging because they involve strong-interaction theory at low energy. Predictions must
81 resort to effective models, such as the heavy-quark expansion [12]. Precise lifetime measure-
82 ments provide excellent tests of such effective models whose precision is still far behind the
83 precision of the experiment.

84 Up to now, the most precise estimates of the D^+ and D^0 lifetimes are from the FOCUS
85 experiment [13], where D mesons were produced in $\gamma(180 \text{ GeV}) + \text{BeO}$ interactions. There is
86 no estimate of these lifetimes from Belle, BaBar or LHCb. However, the LHCb collaboration
87 precisely measured the lifetimes of the D_s^+ [14] mesons and several charmed baryons, e.g.
88 [15], relative to that of the D^+ meson. Such relative measurements minimise systematic
89 uncertainties due to decay-time-biasing event-selection criteria that are particularly severe
90 at hadron colliders. For these LHCb lifetime estimates the external uncertainty coming from
91 the D^+ lifetime is one of the dominant systematic sources.

92 By contrast, experiments at e^+e^- colliders, owing to the reconstruction of large charmed
93 hadron yields without decay-time-biasing selections, have a great potential for absolute
94 lifetime measurements. The high precision Belle II measurement of the D^0 and D^+ lifetimes
95 is based on $D^{*+} \rightarrow D^0(\rightarrow K^-\pi^+)\pi^+$ and $D^{*+} \rightarrow D^+(\rightarrow K^-\pi^+\pi^+)\pi^0$ decay modes [16]. The
96 analysed data correspond to the integrated luminosity of 72 fb^{-1} . At Belle II the D^* meson
97 is produced in the $e^+e^- \rightarrow c\bar{c}$ interactions and the decay vertex of D^* is identical to the
98 primary vertex due to the short D^* lifetime. The flight lengths of D^0 and D^+ are $\approx 200 \mu\text{m}$
99 and $\approx 500 \mu\text{m}$, respectively, which is higher than for B^0 -mesons ($\approx 130 \mu\text{m}$) due to the higher
100 boost of produced D mesons. The decay time t is reconstructed as $t = m_D \vec{L} \cdot \vec{p} / |\vec{p}|^2$, where
101 \vec{p} is the momentum of the D meson and \vec{L} is the vector pointing from the reconstructed
102 D meson production vertex to its decay vertex. The production vertex is estimated with a
103 help of the known beam spot, i.e. the probability distribution of the primary vertexes. At

104 Belle II the beam spot is described by 3D Gaussian distribution and its parameters, which
 105 are heavily time-dependent, are continuously monitored using $e^+e^- \rightarrow \mu^+\mu^-$ events. The
 106 measured spectra of the reconstructed D meson mass and the decay times t are shown in
 Figure 2. The tail towards negative t values is related to the resolution which is about 70 fs

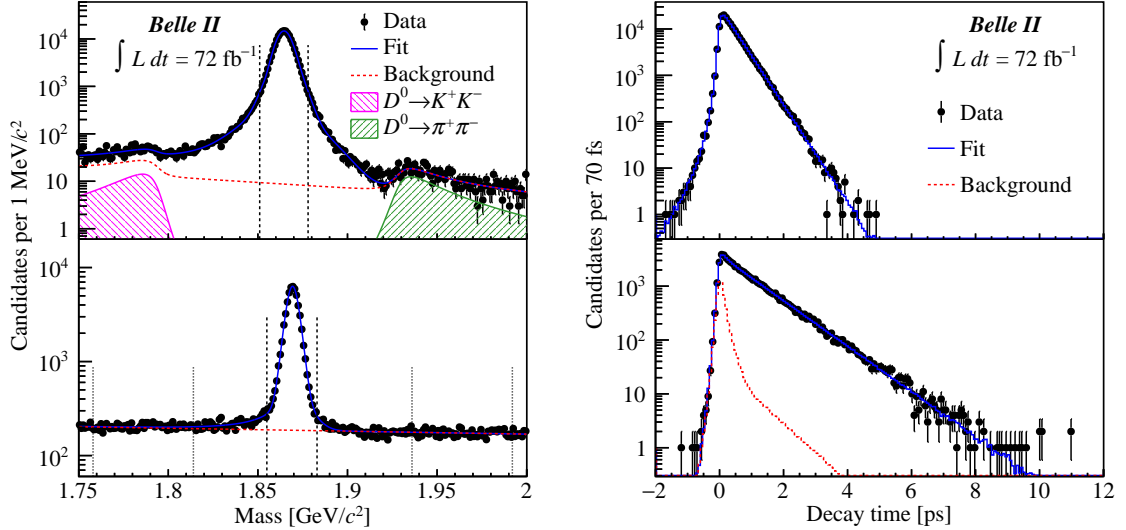


FIG. 2: Mass and decay time distribution of the D^0 (top) and D^+ (bottom) candidates [16]. Vertical lines depict borders of signal and side-band regions.

107
 108 for D^0 and 60 fs for D^+ , respectively.

109 The unbinned maximum likelihood fits of t performed simultaneously for the signal and
 110 side-band D mass regions result in $\tau(D^0) = 410.5 \pm 1.1 \pm 0.8$ fs and $\tau(D^+) = 1030.4 \pm 4.7 \pm$
 111 3.1 fs. For both D meson charges the obtained values are consistent with the world average
 112 but have higher precision. The analysis is statistically limited, the dominant systematic
 113 uncertainty originates from the background modelling for D^+ and from detector alignment
 114 for D^0 .

115 5. CONCLUSIONS

116 The Belle II detector is an excellent machine to study the time-dependent CP violation
 117 in the $B^0 - \bar{B}^0$ and the lifetimes of the heavy-flavor hadrons. For the D^+ and D^0 mesons,
 118 the measured lifetimes have world-leading accuracy. In the case of the parameters related
 119 to the CP violation, like $\sin 2\phi_1$, higher integrated luminosity is needed for the competitive
 120 results.

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