# Belle II input for y

+

[considerations on global y combinations]



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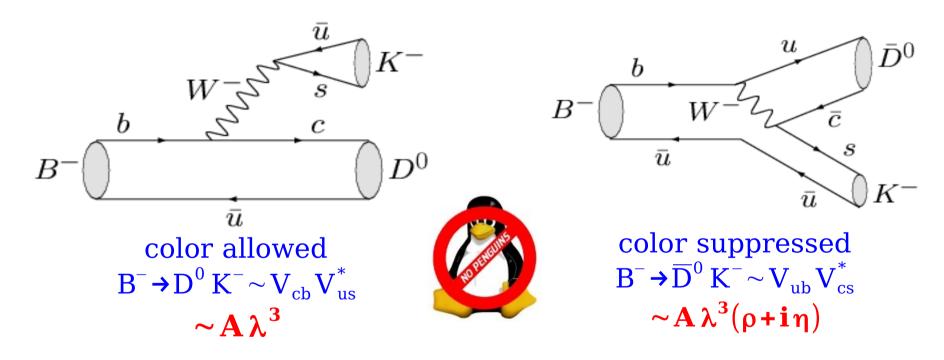
# **Outline**

- Quick estimation of Belle II sensitivity for  $\gamma$  with B→DK, D→K<sub>S</sub>π<sup>+</sup>π<sup>-</sup> as golden mode
- Potential improvements

Towards the Ultimate Precision in Flavour Physics

# $\gamma$ measurements from $B^{\pm} \rightarrow DK^{\pm}$

- ∘ Theoretically pristine B → DK approach
- ∘ Access  $\gamma$  via interference between  $B^- \rightarrow D^0 K^-$  and  $B^- \rightarrow \overline{D}^0 K^-$



relative magnitude of suppressed amplitude is r<sub>B</sub>

$$r_{\rm B} = \frac{|A_{\rm suppressed}|}{|A_{\rm favoured}|} \sim \frac{|V_{\rm ub}V_{\rm cs}^*|}{|V_{\rm cb}V_{\rm us}^*|} \times [{\rm color\ supp}] = 0.1 - 0.2$$

relative weak phase is  $\gamma$ , relative strong phase is  $\delta_B$ 

# $\gamma$ measurements from $B^{\pm} \rightarrow DK^{\pm}$

- Reconstruct D in final states accessible to both  $D^0$  and  $\overline{D}^0$ 
  - D = D<sub>CP</sub>, CP eigenstates as  $K^+K^-$ ,  $\pi^+\pi^-$ ,  $K_S\pi^0$ **GLW method (Gronau-London-Wyler)**
  - D = D<sub>sup</sub>, Doubly-Cabbibo suppressed decays as  $K\pi$ ADS method (Atwood-Dunietz-Soni)
  - − Three-body decays as D→ $K_S \pi^+ \pi^-$ ,  $K_S K^+ K^-$ GGSZ (Dalitz) method (Giri-Grossman-Soffer-Zupan)
  - Largest effects due to
    - charm mixing
    - charm CP violation

negligible
Y.Grossman, A.Soffer, J.Zupan
[PRD 72, 031501 (2005)]

- Different B decays (DK, D\*K, DK\*)
  - different hadronic factors  $(r_B, \delta_B)$  for each

# $\gamma$ measurements from $B^{\pm} \rightarrow DK^{\pm}$

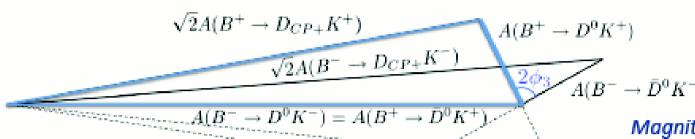
 $B^{\pm} \rightarrow DK^{\pm}$   $B^{\pm} \rightarrow D^{*}K^{\pm}, D^{*} \rightarrow D\pi^{0}$   $B^{\pm} \rightarrow D^{*}K^{\pm}, D^{*} \rightarrow D\gamma$   $B^{\pm} \rightarrow DK^{*\pm}$   $B^{0} \rightarrow DK^{*0}$   $B^{\pm} \rightarrow DK\pi\pi$   $B \rightarrow ...$ 



 $D \rightarrow K^{+}K^{-}, \pi^{+}\pi^{-}...$   $D \rightarrow K_{S}\pi^{0}, K_{S}\eta...$   $D \rightarrow KK\pi^{0}, \pi\pi\pi^{0}...$   $D \rightarrow K_{S}\pi\pi, K_{S}KK$   $D \rightarrow K_{S}\pi\pi\pi^{0}$  $D \rightarrow ...$ 

### D decays to CP eigenstates

#### Amplitude triangle:



#### measured observables:

$$R_{\mathrm{CP}^{\pm}} \equiv \frac{Br(B^{-} \rightarrow D_{\mathrm{CP}^{\pm}}K^{-}) + Br(B^{+} \rightarrow D_{\mathrm{CP}^{\pm}}K^{+})}{Br(B^{-} \rightarrow D^{0}K^{-}) + Br(B^{+} \rightarrow \overline{D}^{0}K^{+})}$$

Magnitude of one side is  $\sim$ 0.1 of the others while relative magnitude of the others help  $\phi_3$  constraint.

$$A_{\mathrm{CP}^{\pm}} \equiv \frac{Br(B^{-} \rightarrow D_{\mathrm{CP}^{\pm}}K^{-}) - Br(B^{+} \rightarrow D_{\mathrm{CP}^{\pm}}K^{+})}{Br(B^{-} \rightarrow D_{\mathrm{CP}^{\pm}}K^{-}) + Br(B^{+} \rightarrow D_{\mathrm{CP}^{\pm}}K^{+})}$$

Relation between 
$$(R_{CP+}, R_{CP-}, A_{CP+}, A_{CP-})$$
 and  $(\gamma, r_B, \delta_B)$ 

$$\mathbf{R}_{\mathrm{CP}+} = 1 + \mathbf{r}_{\mathrm{B}}^{2} + 2 \, \mathbf{r}_{\mathrm{B}} \cos \delta_{\mathrm{B}} \cos \gamma$$

$$\mathbf{A_{CP+}} = \frac{+2 \, \mathbf{r_B} \sin \delta_B \sin \gamma}{1 + \mathbf{r_B}^2 + 2 \, \mathbf{r_B} \cos \delta_B \cos \gamma}$$

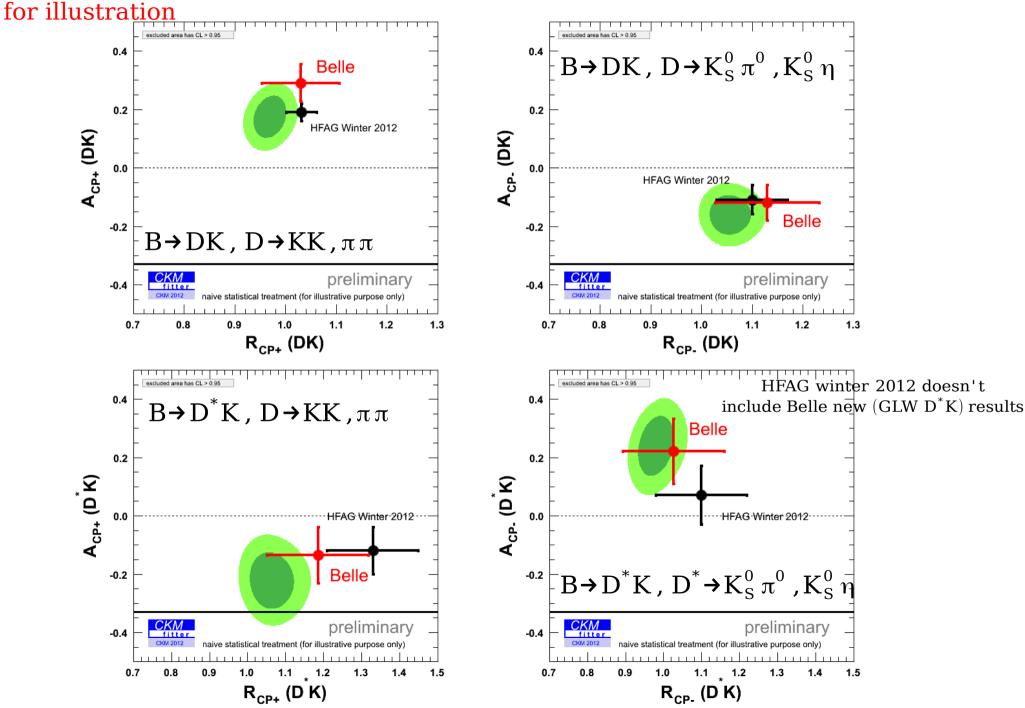
$$R_{CP} = 1 + r_B^2 - 2r_B \cos \delta_B \cos \gamma$$

$$\mathbf{A_{CP}} = \frac{-2 \, \mathbf{r_B} \sin \delta_B \sin \gamma}{1 + \mathbf{r_B}^2 - 2 \, \mathbf{r_B} \cos \delta_B \cos \gamma}$$

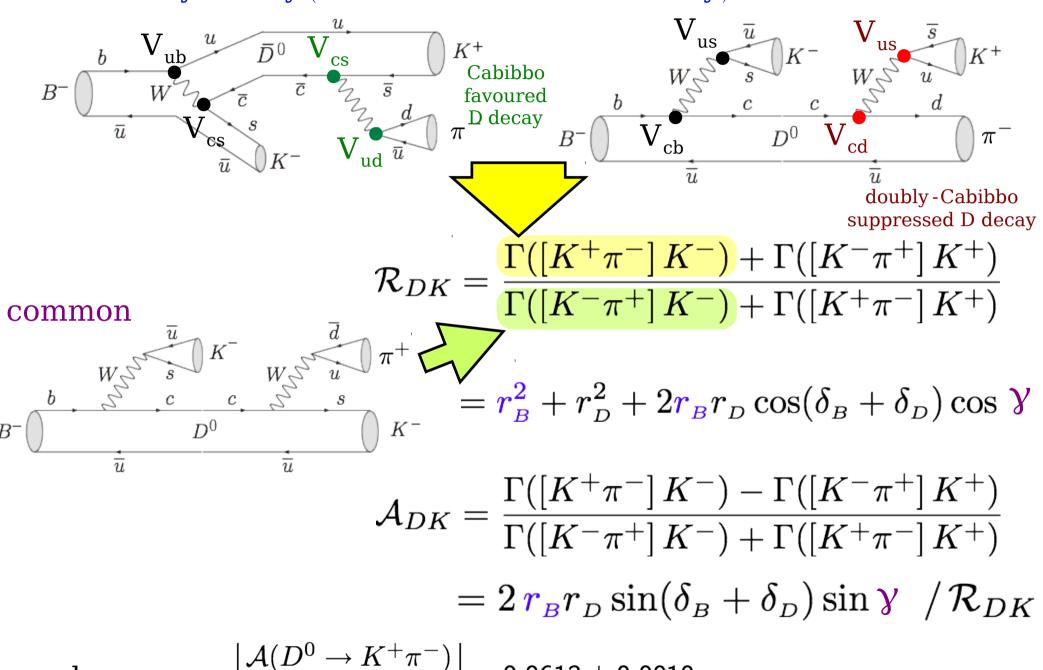
- $\Rightarrow$  look for  $R_{CP+} \neq 1$  and  $A_{CP+} \neq 0$
- $\Rightarrow \neq CP, \neq sign of asymmetry$

# Comparison of the results obtained for GLW D<sup>(\*)</sup>K with expectations

where ''expectations'' are derived from the GGSZ observables (W . A .),  $\delta_{D}$  and  $\gamma_{UT}$ 



<u>ADS method</u>:  $\gamma$  via the interference in rare  $B^- \rightarrow [K^+\pi^-]_D K^-$  decays rate and asymmetry (relative to the common decay):



where  $r_D = \left| \frac{\mathcal{A}(D^0 \to K^+\pi^-)}{\mathcal{A}(\bar{D}^0 \to K^+\pi^-)} \right| = 0.0613 \pm 0.0010$ 

#### How to get $\delta_{\rm D}$ and related (charm) hadronic parameters ?

- $\circ$  dedicated experiments (CLEO-c, BES III) using quantum correlations, running at  $\psi(3770)$ 
  - $\circ$  CLEO-c:  $R_D$ ,  $\cos \delta_D$ ,  $\sin \delta_D$  (but also BES III result...)
  - $\circ$  CLEO-c:  $R_{K\pi\pi^0}$ ,  $\delta_{K\pi\pi^0}$ ,  $R_{K3\pi}$ ,  $\delta_{K3\pi}$

 $R_f$ : coherence factor, can take any value from 0 to 1 indicates lack coherence between the intermediate states involved in the decay

- mixing/CPV results from BaBar, Belle, CDF, LHCb...
  - ∘ D→KK,  $\pi\pi$ :  $y_{CP}$ ,  $A_{\Gamma}$  (BaBar, Belle, LHCb)
  - ∘ D $\rightarrow$ K<sub>S</sub><sup>0</sup> $\pi\pi$ : x, y, |q/p|,  $\phi$  (BaBar, Belle)
  - ∘ D→Klv: R<sub>M</sub> (BaBar, Belle...)
  - ∘ D →  $K \pi \pi^0$ : x'', y'' (BaBar)
  - ∘ D→Kπ: x', y' (BaBar, Belle, CDF, LHCb)
  - 0
- CLEO-c/BES III, use external inputs to access the relevant physics parameters
- strong phases information in B-factories/LHCb
- x, y are also needed for D-mixing corrections in ADS observables

$$R^{\mp} = r_{B}^{2} + r_{D}^{2} + 2r_{B}r_{D}\cos(\delta_{B}\mp\gamma + \delta_{D})$$

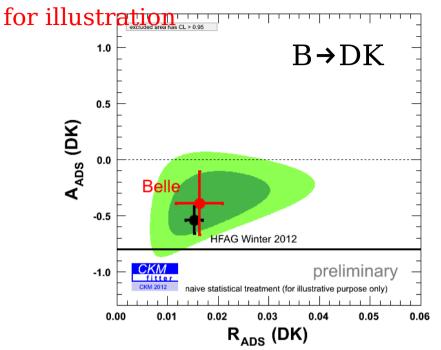
$$\Rightarrow R^{\mp} = r_{B}^{2} + r_{D}^{2} + 2r_{B}r_{D}\cos(\delta_{B}\mp\gamma + \delta_{D}) - yr_{D}\cos\delta_{D} - yr_{B}\cos(\delta_{B}\mp\gamma) + xr_{D}\sin\delta_{D} - xr_{B}\sin(\delta_{B}\mp\gamma)$$

 $\Rightarrow$  combine charm observables to obtain  $\gamma$  and mixing/CPV charm parameters

 $\delta_{\rm D}$  grand combination à la HFAG  $\sim$  35 observables  $\overline{n}$ pt a great fit  $(\sim 3\sigma)$ Kπ (LHCb) (1.64) 8 parameters: ( 0.57 Kπ (CDF) x, y,  $\delta_D^{K\pi}$ ,  $r_D$ ,  $A_D$ , |q|/|p|,  $\phi$ ,  $\delta_D(K\rho)$ Kπ (Belle) ₹ 0.01 ) Kπ (BaBar) (3.10)  $K\pi\pi^{0}$  (BaBar) ( 3.37 ) CLEO-c (1.88) **BES III** ( 0.37 ) ( 0.36 ) K<sub>s</sub>ππ (Belle) (1.08) $K_s\pi\pi$  (BaBar) (0.90)( 0.68 ) (1.82) ---- ADS+GLW Combined GGSZ+ADS+GLW ---- All charm 0 Pull (σ) (include K3 $\pi$ , K $\pi\pi^0$  info, see next slides) 8.0 p-value 0.6 All charm:  $\delta_{\rm D}^{\rm K\pi} = (191.4^{+8.2}_{-11.4})^{\circ} (^{+16}_{-30})$ 0.4 GGSZ+GLW+ADS:  $\delta_{D}^{K\pi} = (193^{+18}_{-23})^{\circ} (^{+34}_{-77})$ 0.2 0.0 300 50 100 150 200 250 350  $\delta_{\rm D}(K\pi)$ 

# Comparison of the results obtained for D<sup>(\*)</sup>K with expectations

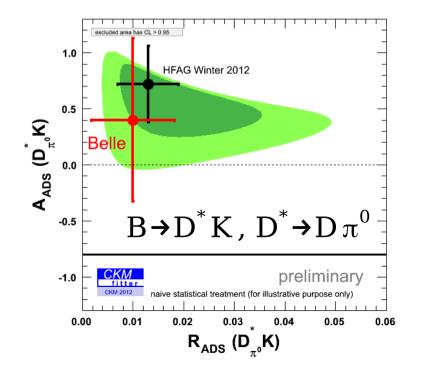
where ''expectations'' are derived from the GGSZ observables,  $\delta_D$  and  $\gamma_{UT}$ 

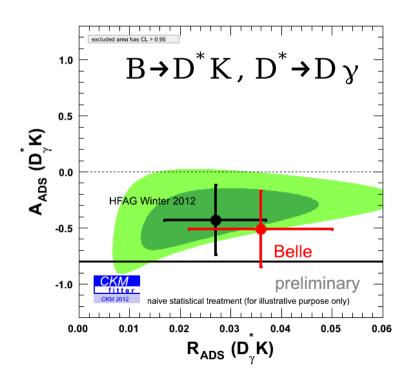


$$\mathbf{R}_{ADS}(\mathbf{DK}) = \mathbf{r}_{B}^{2} + \mathbf{r}_{D}^{2} + 2\mathbf{r}_{B}\mathbf{r}_{D}\mathbf{cos}(\delta_{B} + \delta_{D})\mathbf{cos}\gamma$$
$$\mathbf{A}_{ADS}(\mathbf{DK}) = 2\mathbf{r}_{B}\mathbf{r}_{D}\mathbf{sin}(\delta_{B} + \delta_{D})\mathbf{sin}\gamma/\mathbf{R}_{ADS}(\mathbf{DK})$$

$$\begin{aligned} \mathbf{R}_{\mathrm{ADS}}(\mathbf{D}_{\pi^{0}}^{*}\mathbf{K}) &= \mathbf{r}_{\mathrm{B}}^{*2} + \mathbf{r}_{\mathrm{D}}^{2} + \mathbf{2}\mathbf{r}_{\mathrm{B}}^{*}\mathbf{r}_{\mathrm{D}}\mathbf{\cos}(\delta_{\mathrm{B}}^{*} + \delta_{\mathrm{D}})\mathbf{\cos}\gamma \\ \mathbf{A}_{\mathrm{ADS}}(\mathbf{D}_{\pi^{0}}^{*}\mathbf{K}) &= \mathbf{2}\mathbf{r}_{\mathrm{B}}^{*}\mathbf{r}_{\mathrm{D}}\mathbf{\sin}(\delta_{\mathrm{B}}^{*} + \delta_{\mathrm{D}})\mathbf{\sin}\gamma / \mathbf{R}_{\mathrm{ADS}}(\mathbf{D}_{\pi^{0}}^{*}\mathbf{K}) \end{aligned}$$

$$\begin{aligned} \mathbf{R}_{ADS}(\mathbf{D}_{\gamma}^{*}\mathbf{K}) &= \mathbf{r}_{B}^{*2} + \mathbf{r}_{D}^{2} - 2\mathbf{r}_{B}^{*}\mathbf{r}_{D}\mathbf{cos}(\delta_{B}^{*} + \delta_{D})\mathbf{cos}\,\gamma \\ \mathbf{A}_{ADS}(\mathbf{D}_{\gamma}^{*}\mathbf{K}) &= -2\mathbf{r}_{B}^{*}\mathbf{r}_{D}\mathbf{sin}(\delta_{B}^{*} + \delta_{D})\mathbf{sin}\,\gamma/\mathbf{R}_{ADS}(\mathbf{D}_{\gamma}^{*}\mathbf{K}) \end{aligned}$$





# Sensitivity to $\gamma$ in $B \rightarrow D(K_S \pi \pi)K$ mode

sensitivity to  $\gamma/\phi_3$  varies across the Dalitz plot

 $\gamma = 75^{\circ}$ ,  $\delta = 180^{\circ}$ ,  $r_{B} = 0.125$  $w=1/(d^2L/d\gamma^2)$ **GLW** like Interference of BABAR $B^- \rightarrow D^0 K^-$ ,  $D^0 \rightarrow K_S^0 \rho^0$ pre linanur y with  $B^- \rightarrow \overline{D}^0 K^-$ ,  $\overline{D}^0 \rightarrow K_s^0 \rho^0$ DCS  $K^*(1^2430)$ 1.5 **ADS** like Interference of  $B^- \rightarrow D^0 K^-$ ,  $D^0 \rightarrow K^{*+} \pi^-$ DCS K\*(892 10 with 0.5  $B^- \rightarrow \overline{D}^0 K^-$ ,  $\overline{D}^0 \rightarrow K^{*+} \pi^$  $m_{\perp}^{2} (GeV^{2}/c^{4})$ 

- golden mode!! even more for Belle II than for LHCb
- focusing our efforts/resources on this mode

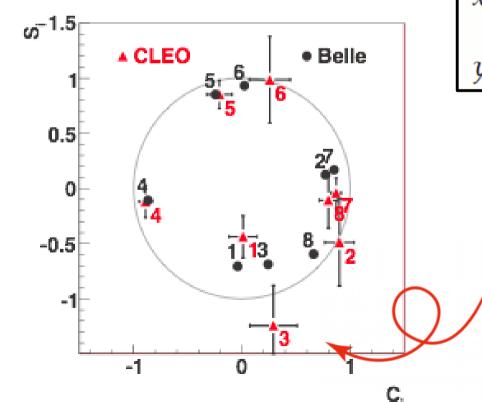
# **Binned Dalitz method:** avoid the modeling error by ''optimal'' binning of the Dalitz plot

[choice of bins guided by model, but extraction of  $\gamma$  is not biased by this choice]

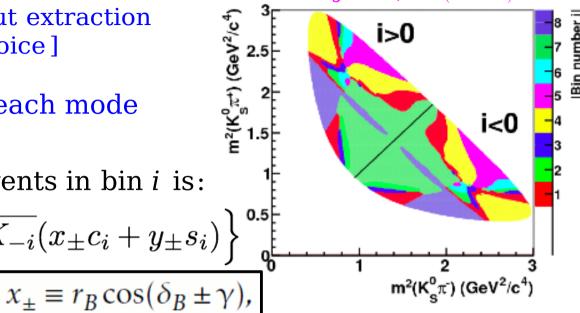
minimize  $\chi^2$  in fit to all bins for each mode

Expected number of  $B^{\pm} \rightarrow DK^{\pm}$  events in bin *i* is:

$$N_{i}^{\pm} = h \left\{ K_{i} + r_{B}^{2} K_{-i} + 2 \sqrt{K_{i} K_{-i}} (x_{\pm} c_{i} + y_{\pm} s_{i}) \right\}_{\mathbf{q}}^{\mathbf{0.5}}$$



Bondar and Poluektov EPJ C55, 51 (2008)



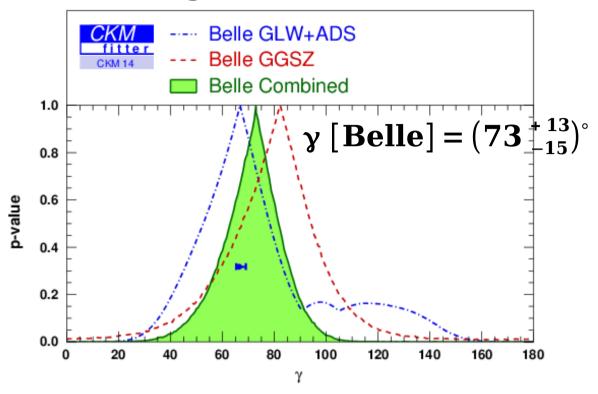
 $y_{\pm} \equiv r_B \sin(\delta_B \pm \gamma).$ 

 $K_i$  is the # of events in bin *i* from a flavour-tagged sample  $(D^{*\pm} \rightarrow D\pi^{\pm})$ 

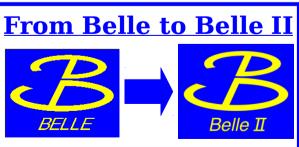
 $c_i$  and  $s_i$  contain information about the strong-phase difference in bin i

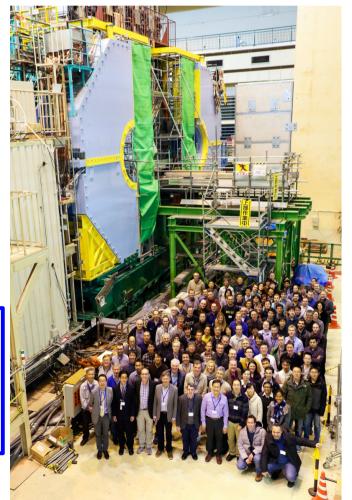
(use CLEO data for  $\psi(3770) \rightarrow D^0 \overline{D}^0$  here; measured by BES-III too)

# **Combining measurements for y from all methods**

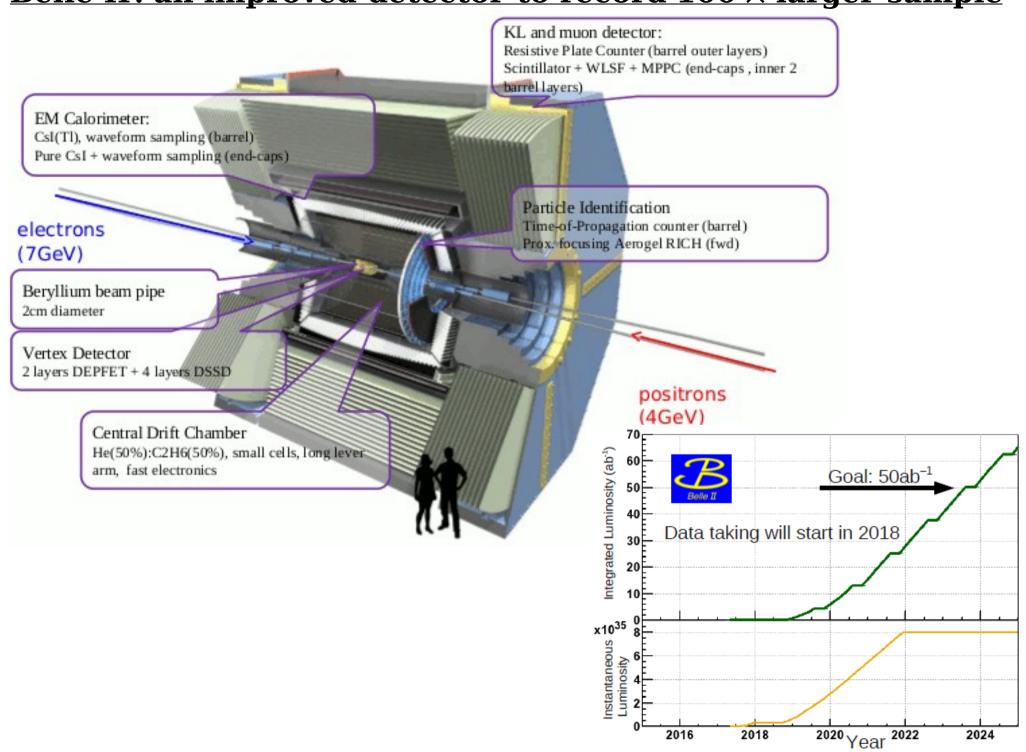








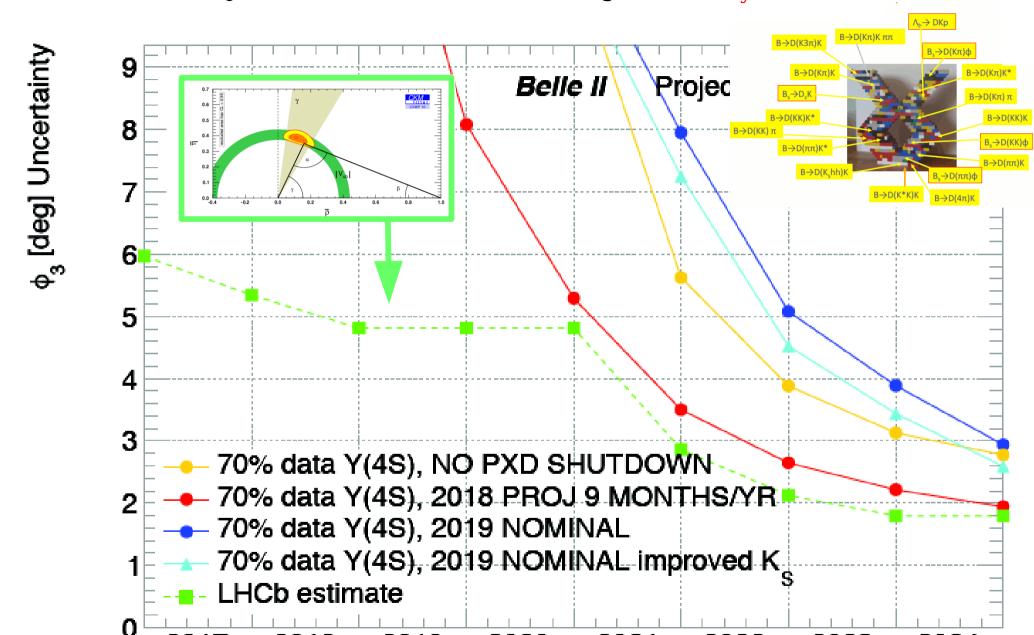
#### Belle II: an improved detector to record 100 × larger sample



# Ultimate y-from-tree decays

precision will be reached through many individual measurements

 $(\sigma_{theory} \ negligible)$ 



Year

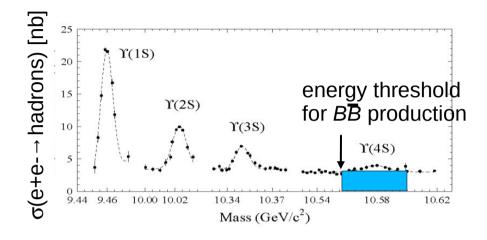
# Potential improvements Belle II vs Belle

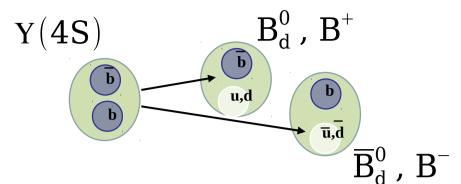
(beyond only statistics)

- continuum suppression
- PID performances
- o new possible avenues...

# Y(4S) B-factory

#### but also continuum factory....

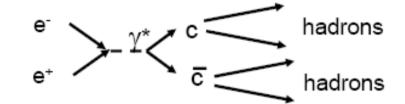




- 2 B's and nothing else!
- 2 B mesons are created simultaneously in a L=1 coherent state
  - $\Rightarrow$  before first decay, the final states contains a B and a  $\overline{B}$

#### "on resonance" production

$$e^+e^- \rightarrow Y(4\,S) \rightarrow B_d^0\overline{B}_d^0$$
,  $B^+B^-$   
 $\sigma(e^+e^- \rightarrow B\,\overline{B}) \simeq 1.1 \ nb \ (\sim 10^9 \ B\,\overline{B} \ pairs)$ 



• ''continuum'' production  $(q\overline{q} = u\overline{u}, d\overline{d}, s\overline{s}, c\overline{c})$ 

$$\sigma(e^+e^- \rightarrow c\overline{c}) = 1.3 \text{ nb}$$

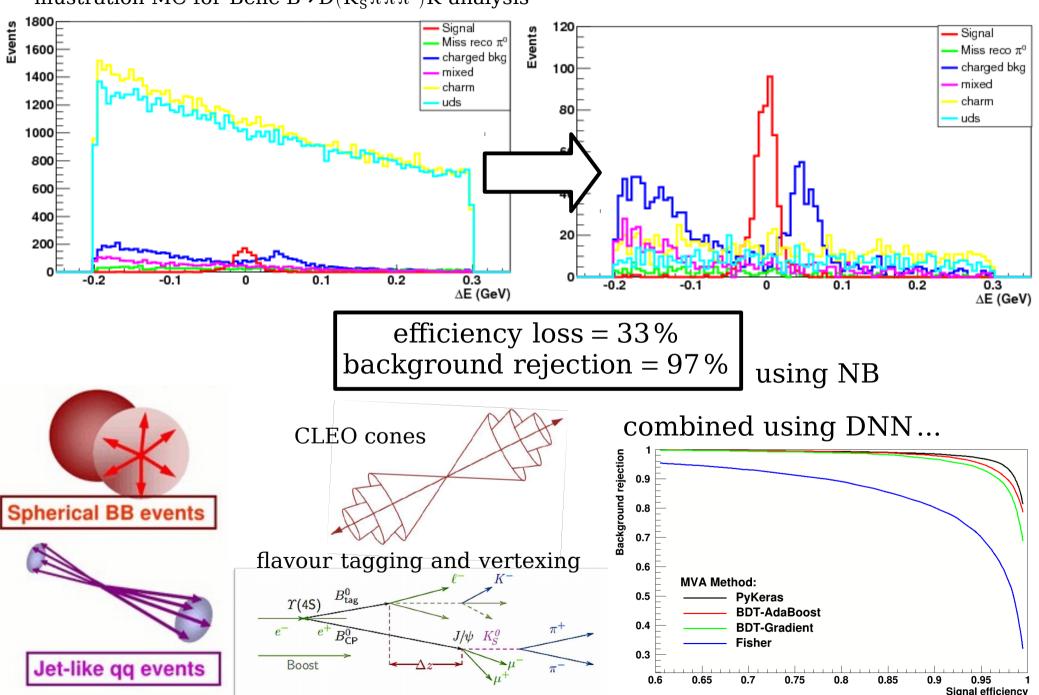
$$\sigma(e^+e^- \rightarrow s\overline{s}) = 0.4 \text{ nb}$$

$$\sigma(e^+e^- \rightarrow u\overline{u}) = 1.6 \text{ nb}$$

$$\sigma(e^+e^- \rightarrow d\overline{d}) = 0.4 \text{ nb}$$

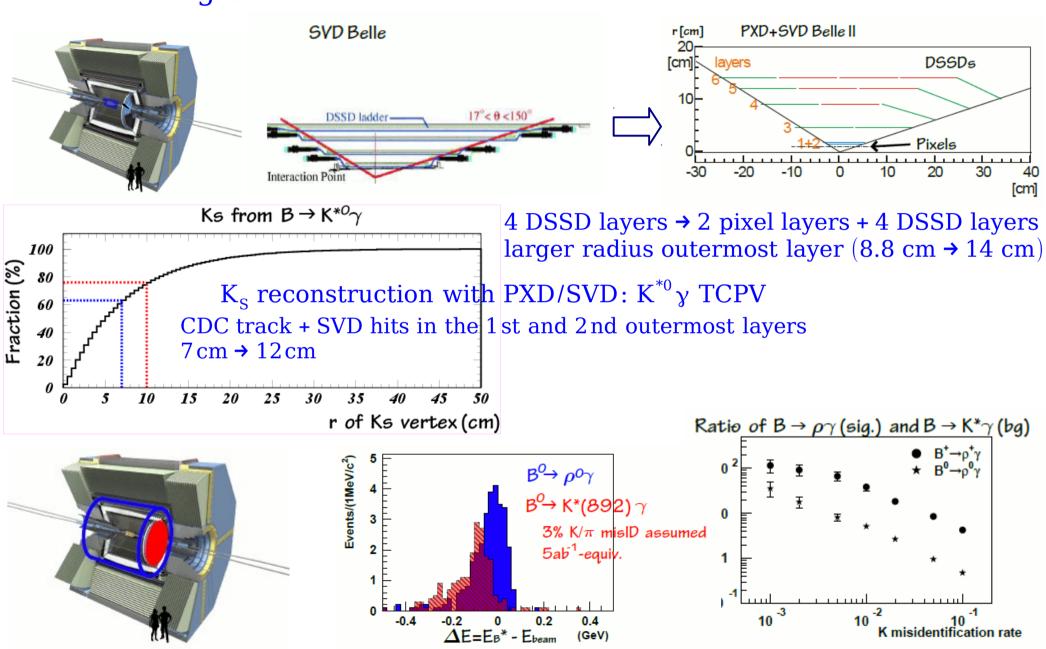
# $B \rightarrow [K_S \pi^+ \pi^-]_D K^{\pm}$ Dalitz Analysis with Belle II

illustration MC for Belle  $B \rightarrow D(K_S \pi \pi \pi^0)K$  analysis



# Belle II in few words

 $\circ$  collecting  $50ab^{-1}$  from 2019 to 2027



⇒ new detectors (CDC, TOP, ARICH) in place (see P. Urquijo's talk)

### B→DK<sup>±</sup> at Belle II

 $B \rightarrow D\pi$  $B \rightarrow DK$ 

illustration with Belle  $B \rightarrow D(K\pi)K$  analysis

$$\begin{split} N_{\eta,\,KID > 0.6}^{DK} &= \frac{1}{2} \left( 1 - \eta A^{DK} \right) \, N_{tot}^{D\pi} \, R_{K/\pi} \, \epsilon \\ N_{\eta,\,KID < 0.6}^{DK} &= \frac{1}{2} \left( 1 - \eta A^{DK} \right) \, N_{tot}^{D\pi} \, R_{K/\pi} \, \left( 1 - \epsilon \right) \\ N_{\eta,KID > 0.6}^{D\pi} &= \frac{1}{2} \left( 1 - \eta A^{D\pi} \right) \, N_{tot}^{D\pi} \, \kappa \\ N_{\eta,KID < 0.6}^{D\pi} &= \frac{1}{2} \left( 1 - \eta A^{D\pi} \right) \, N_{tot}^{D\pi} \, \left( 1 - \kappa \right) \end{split}$$

#### KID>0.6 (kaon-like)

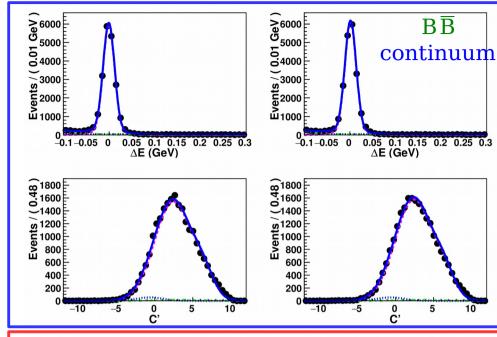
140	kaon fake $(1-\epsilon)$	kaon eff ε	pion eff $(1-\kappa)$	pion fake κ	⇐
MC data	$14.70 \pm 0.06$ $15.86 \pm 0.40$	$85.41 \pm 0.06$ $84.32 \pm 0.39$	$95.42 \pm 0.03$ $92.13 \pm 0.46$		

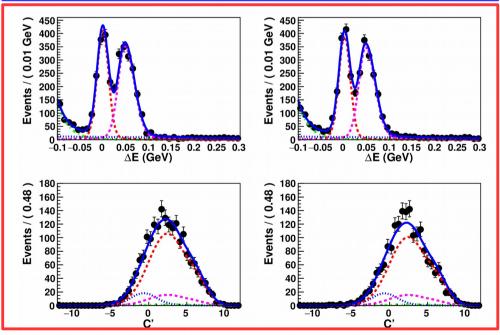
for Belle

for Belle II: performances expected

to be as good (better?) as for Belle MC...

one of the important outputs of current data taking (jury is still out)





# Lot of interesting modes...

not used until now

D mode	$2F_{+}-1$	branching ratio
		$( imes 10^{-3})$
$K^+K^-$	+1	$3.96 \pm 0.08$
$\pi^{+}\pi^{-}$	+1	$1.40\!\pm\!0.03$
$\pi^0 \ \pi^0$	+1	$0.82 \!\pm\! 0.04$
$K_{ m L}^0\pi^0$	+1	$10.0 \pm 0.7$
$ extstyle K_{ extstyle S}^0 \pi^0 \ \pi^0$	+1	$9.1 \pm 1.1$
$K^0_S\eta\pi^0$	+1	$5.5 \!\pm\! 1.1$
$\mathbf{K}_{\mathrm{S}}^{0}\mathbf{K}_{\mathrm{S}}^{0}\mathbf{K}_{\mathrm{S}}^{0}$	+1	$0.91 \pm 0.13$
$\pi\pi\pi^0$		$14.3 \pm 0.6$
${\rm K}{\rm K}\pi^0$		$3.3 \pm 0.1$
ππππ		$7.4 \pm 0.2$

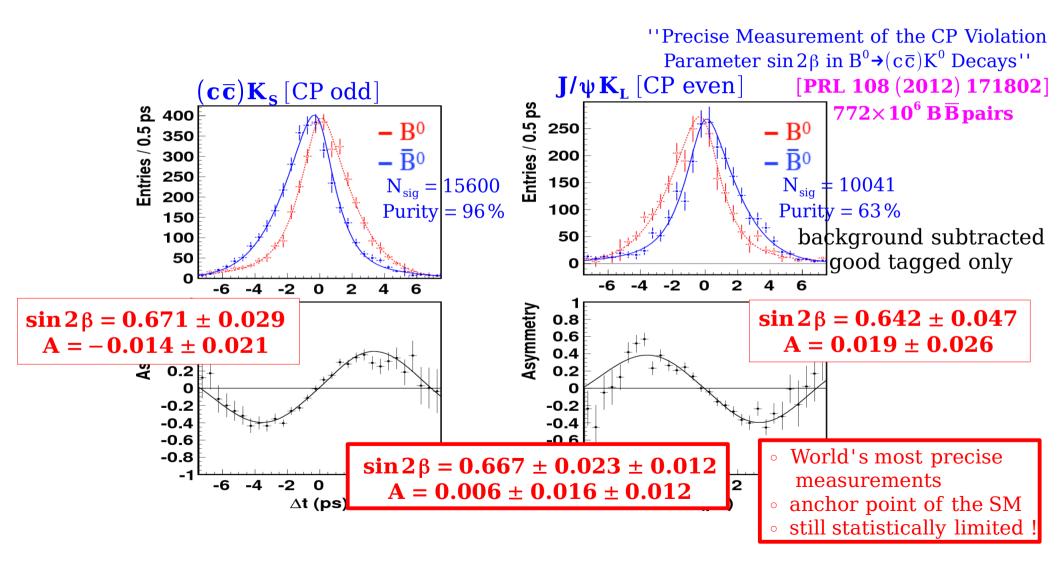
D mode	$2F_{+}-1$	branching ratio
		$( imes 10^{-3})$
${\rm K}_{\rm S}^0\pi^0$	-1	$11.9\!\pm\!0.4$
${ m K}_{ m S}^0\eta$	-1	$4.8 \!\pm\! 0.3$
extstyle  ext	-1	$9.4\!\pm\!0.5$
$\mathbf{K}_{\mathrm{S}}^{0}\mathbf{K}_{\mathrm{S}}^{0}\mathbf{K}_{\mathrm{L}}^{0}$	-1	1.0
$\eta\pi^0\pi^0$	-1	unknown
$\eta$ ' $\pi^0$ $\pi^0$	-1	unknown
$\mathrm{K}^0_\mathrm{S}\mathrm{K}^0_\mathrm{S}\pi^0$	-1	< 0.6
$K^0_S K^0_S \eta$	-1	unknown

#### **D** mode branching ratio $(\times 10^{-3})$

challenging modes with  $K_L$ , two  $\pi^0$ 's...

# $\mathbf{B} \to \mathbf{D}(\mathbf{K}_{\mathbf{L}} \pi \pi) \mathbf{K}$

- ∘ D→ $K_L\pi\pi$  has never been explored in B-factories
- $\circ$  However,  $J/\psi K_L$  has been used for  $\sin 2\beta$  extraction
- with a reasonable efficiency/purity (and a significant impact)
- potential is even more promising in Belle II (upgraded KLM with scintillators)



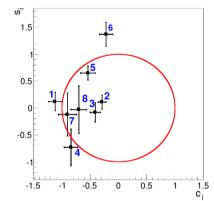
# Estimates of $\gamma$ sensitivity with $B^{\pm} \rightarrow D(K_S \pi \pi \pi \pi^0) K^{\pm}$

- ∘ The decay  $D^0 \rightarrow K_S^0 \pi^+ \pi^- \pi^0$  has a relatively large branching fraction of 5.2%, almost twice that of  $K_S^0 \pi^+ \pi^-$
- Interesting resonance substructure
  - $-K_S^0 \omega CP$  eigenstate GLW like
  - $-K^{*+}\pi^{-}\pi^{0}$  Cabibbo favored state (CF) ADS like
  - CLEO-c obtained  $F_+ = 0.240 \pm 0.021$  (significantly CP-odd)

Bin number	Specification
1	$m(\pi^+\pi^-\pi^0)\approx m(\omega)$
2	$m(K_S^0\pi^-) \approx m(K^{*-}) \& m(\pi^+\pi^0) \approx m(\rho^+)$
3	$m(K_S^0\pi^+) \approx m(K^{*+}) \& m(\pi^-\pi^0) \approx m(\rho^-)$
4	$m(K_S^0\pi^-)\approx m(K^{*-})$
5	$m(K_S^0\pi^+) \approx m(K^{*+})$
6	$m(K_S^0\pi^0)\approx m(K^{*0})$
7	$m(\pi^+\pi^0) pprox m( ho^+)$
8	Remainder

∘  $c_i < 0 \Rightarrow CP$  oddness of  $K_S^0 \pi^+ \pi^- \pi^0$ 

•			
Bin	$c_i$	Si	ώ¯
1	$-1.12 \pm 0.12$	0.12 ± 0.17	1.5
2	$-0.29 \pm 0.07$	$0.11 \pm 0.13$	1 <u>-</u>
3	$-0.41 \pm 0.09$	$-0.08 \pm 0.18$	0.5
4	$-0.84 \pm 0.12$	$-0.73 \pm 0.34$	1 82,4
5	$-0.54 \pm 0.13$	$0.65 \pm 0.13$	0 7 7
6	$-0.22 \pm 0.12$	$1.37 \pm 0.22$	-0.5
7	$-0.90 \pm 0.16$	$-0.12 \pm 0.40$	1 1
8	$-0.70 \pm 0.14$	$-0.03 \pm 0.44$	-1.5 -1 -0.5



- $\circ$  Project to a 50 ab $^{-1}$  sample  $\sigma_{\gamma} \sim 3.5^{\circ}$
- ∘ compare to  $B^{\pm} \rightarrow D(K_S^0 \pi^+ \pi^-) K^{\pm} \sigma_{\gamma} \sim 2^{\circ}$
- on-going Belle analysis should give us a more precise estimation soon

# c<sub>i</sub> and s<sub>i</sub> at charm factory

at  $\psi(3770)$ ,  $J^{PC}=1^{-1}$ , decays to a  $D\overline{D}$  pair (decay are quantum related) D mesons decay to final states  $f_a$  and  $f_b$  with CP eigenvalues  $\eta_a$  and  $\eta_b$  CP conservation requires that  $\eta_a \eta_b (-1)^L = 1$ , hence  $\eta_a / \eta_b = -1$   $\Rightarrow$  if one D meson is reconstructed in a CP even (odd) eigenstate, other D meson must be CP odd (even) eigenstate

measurements of  $c_i$  and  $s_i$  require that one of the D mesons decays to  $K_S^0\pi^+\pi^-$  final state and the other decays to final state  $X_D$  if  $X_D$  is CP even (odd) eigenstate, D meson decaying to  $K_S^0\pi^+\pi^-$  must be CP-odd (even) amplitude and partial width of  $D_+$  at Dalitz plot coordinate  $(m_-^2, m_+^2)$ :

$$\begin{split} A(D_{\pm} \to K_{\rm S}^0 h^+ h^-) &= \frac{1}{\sqrt{2}} \left( A_D \pm \overline{A}_D \right), \\ \frac{d\Gamma(D_{\pm} \to K_{\rm S}^0 h^+ h^-)}{dm_-^2 dm_+^2} &= \frac{1}{2} \left( A_D^2 + \overline{A}_D^2 \right) \pm A_D \overline{A}_D \cos \delta_D. \end{split}$$

decay rate to bin i of the D<sub>+</sub> Dalitz plot:

$$\Gamma_i(D_{\pm} \to K_{\rm S}^0 h^+ h^-) \propto \frac{1}{2} (T_i + T_{-i}) \pm \sqrt{T_i T_{-i}} c_i.$$

$$\frac{\text{if } X_{\rm D} \text{ is } K_{\rm S}^0 \pi^+ \pi^-:}{\Gamma_{ij} \propto T_i T_{-j} + T_{-i} T_j - 2 \sqrt{T_i T_{-i} T_j T_{-j}} (c_i c_j + s_i s_j).}$$

arXiv:1010.2817, arXiv:0903.1681

$$c_{i}$$
,  $s_{i}$  for  $D \rightarrow K_{S}^{0} \pi^{+} \pi^{-}$  
$$M_{i}^{\pm} = h_{CP\pm}(K_{i} \pm 2c_{i}\sqrt{K_{i}K_{-i}} + K_{-i}),$$

$$M_{ij} = h_{corr}(K_i K_{-j} + K_{-i} K_j - 2\sqrt{K_i K_{-j} K_{-i} K_j} (c_i c_j + s_i s_j)).$$

$$C_i'$$
,  $S_i'$  for  $D \rightarrow K_L^0 \pi^+ \pi^ M_i'^{\pm} = h_{CP\pm}(K_i' \mp 2c_i' \sqrt{K_i' K_{-i}'} + K_{-i}')$ ,

$$M'_{ij} = h_{corr}(K_i K'_{-j} + K_{-i} K'_j + 2\sqrt{K_i K'_{-j} K_{-i} K'_j} (c_i c'_j + s_i s'_j))$$

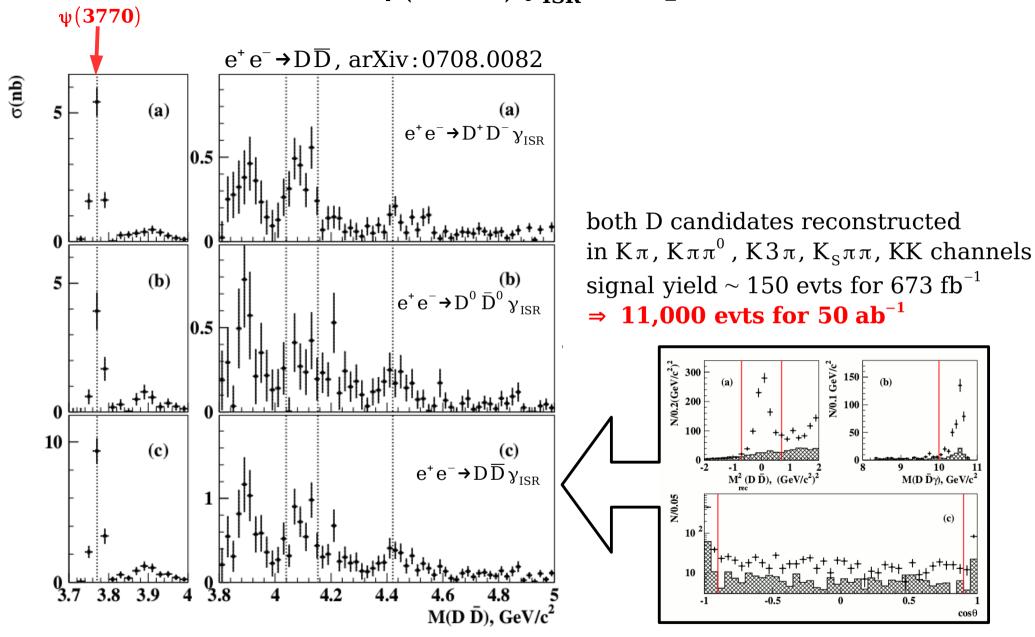
#### $\Delta c_i$ , $\Delta s_i$ are model-dependent

with assumption made to deduce  $\Delta c_i$ , as DCS decays contribute with opposite sign, CP-eigenstate amplitudes related by factor  $(1-2\,r\,e^{i\delta})$ ,  $r=\tan^2\theta_C$ ,  $\delta$  any value use BaBar model

i	$\Delta c_i$	$\Delta s_i$
1	$0.39 \pm 0.17$	$0.07 \pm 0.06$
2	$0.18 \pm 0.05$	$0.01 \pm 0.10$
3	$0.61 \pm 0.15$	$0.30 \pm 0.12$
4	$0.09 \pm 0.08$	$0.00 \pm 0.08$
5	$0.16 \pm 0.17$	$0.06 \pm 0.06$
6	$0.57 \pm 0.21$	$-0.15 \pm 0.24$
7	$0.03 \pm 0.01$	$-0.04 \pm 0.06$
8	$-0.10 \pm 0.15$	$-0.15 \pm 0.21$

By the way,  $c_i$  and  $s_i$  come from charm factories (CLEO-c, BESIII) but could we use  $e^+e^- \rightarrow \psi(3770)\gamma_{ISR}$  sample?

# Could we use $e^+e^- \rightarrow \psi(3770)\gamma_{ISR}$ sample?

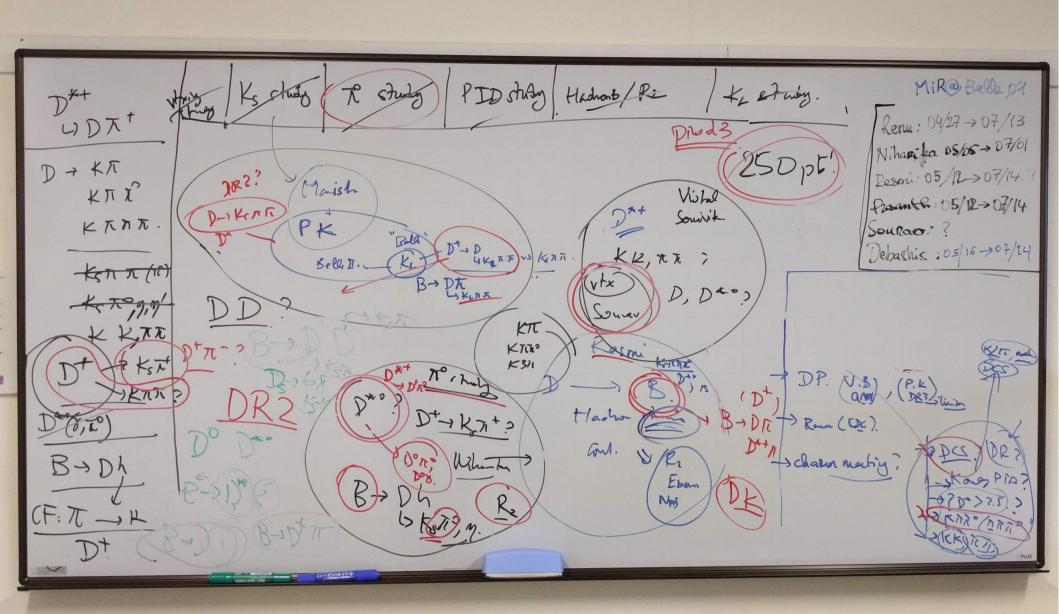


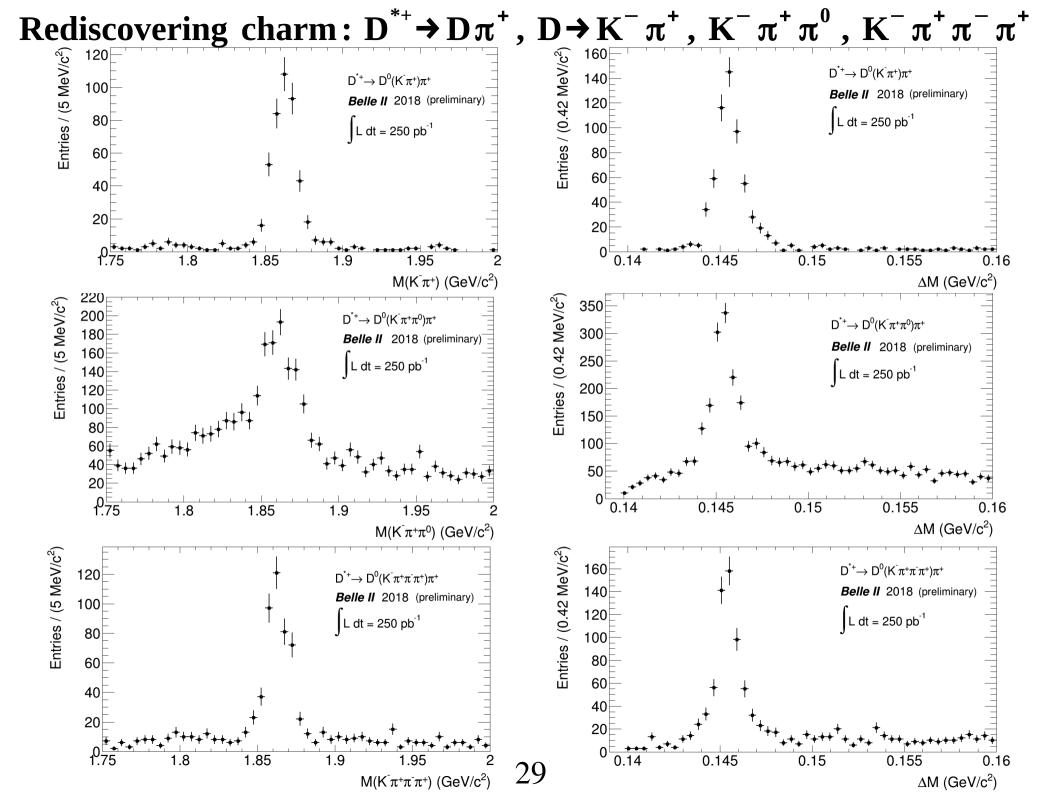
to be compared to CLEO-c, DT yield  $(K_S\pi\pi, K\pi + K\pi\pi^0 + K3\pi + K_S\pi\pi + KK) = 7,000$  evts @ 0.8 fb<sup>-1</sup> only for  $K_S\pi\pi$  mode, only for 0.8 fb<sup>-1</sup> so doesn't seem to be competitive

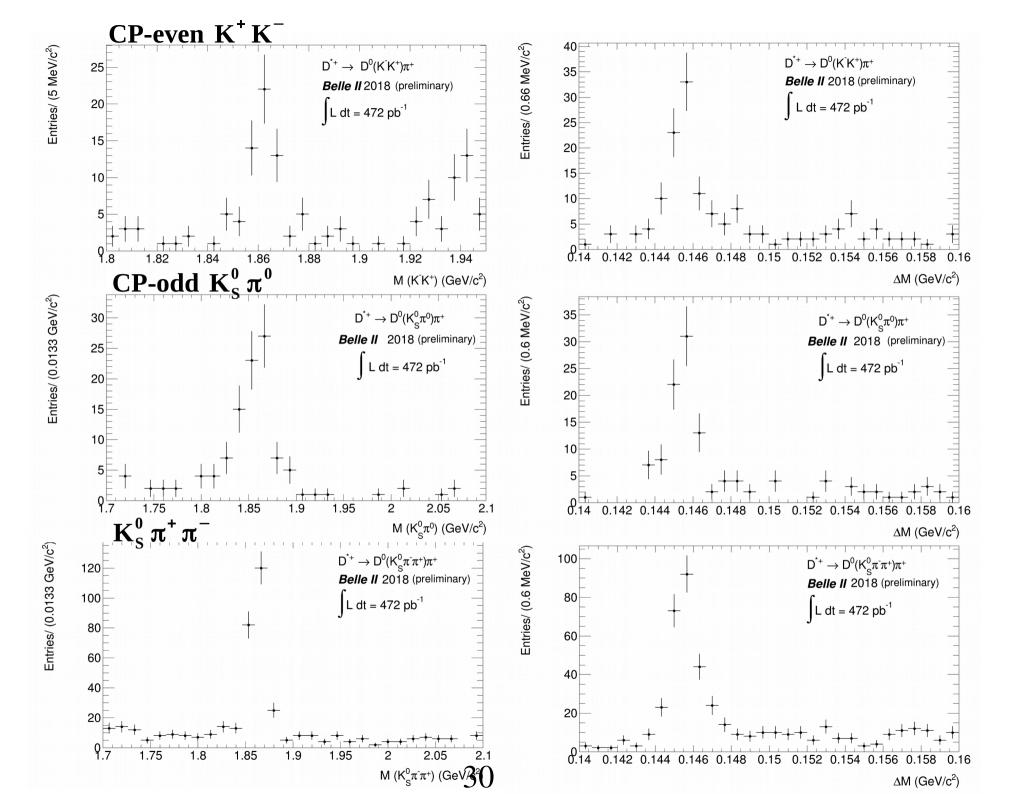
# Time - dependent measurements

- All of the measurements presented so far were time-independent
- Time-dependent measurements (mixing induced CPV) also possible:
  - B<sup>0</sup> → D<sup>(\*)</sup>π, B<sup>0</sup> → D<sup>(\*)</sup>ρ
- ∘ In order to extract  $\gamma$  from B → SS/SV decays, must supply  $r = |A_{DCS}/A_{CF}|$  externally (expected to be ~ 1-2%), usually assuming SU(3) symmetry
  - $\Rightarrow$  not good idea to include those measurements in  $\gamma$  average
- $\circ$  In B  $\rightarrow$  VV decays, one can extract all physics parameters from data
- Belle study:  $\sim 100 \, \text{k}$  evts per  $ab^{-1}$ , 3 helicity configurations:  $A = \sum_{\lambda} A_{\lambda}$  we use Cartesian coordinates  $\{r_{\lambda}, \delta_{\lambda}, \phi_{w}\} \rightarrow \{x_{\lambda}, y_{\lambda}, \overline{x}_{\lambda}, \overline{y}_{\lambda}\}$   $\sigma(2\beta + \gamma) \approx 11^{\circ}$  for Belle II with 50  $ab^{-1}$
- on-going Belle analysis should give us a more precise estimation soon

# A look at first data... (phase 2)

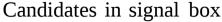


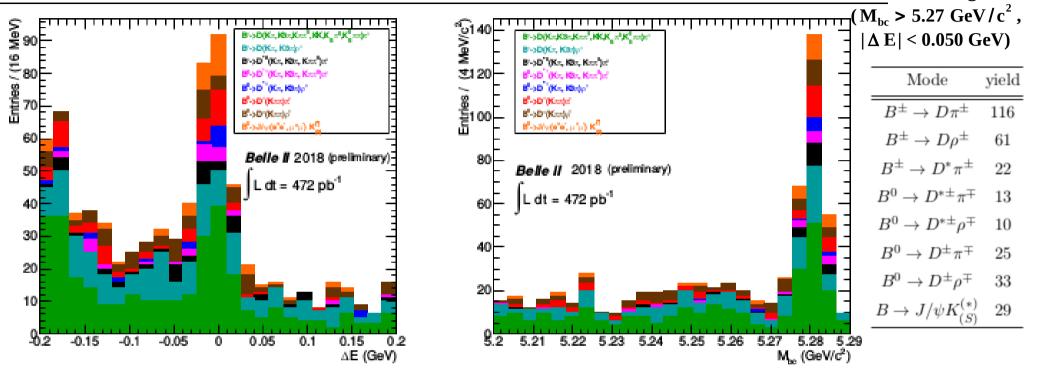




# Rediscovering beauty: $B \rightarrow D^{(*)}h + B \rightarrow J/\psi K^{(*)}$

Results for 0.5 fb<sup>-1</sup>





#### Show capacity for charm physics in $e^+e^- \rightarrow c \bar{c}$

- $\circ$  D<sup>0</sup>, D<sup>+</sup>, D<sup>\*</sup>
- Cabibbo favoured and suppressed modes

#### ... for B-physics

- hadronic modes from b→c
- ∘ semileptonic decay modes from b→c

# Conclusion

"Data! data!" he cried impatiently
"I can't make bricks without clay." (Arthur Conan Doyle)

- $\circ$  Promising perspectives at Belle II for  $\gamma$  measurement
- To stay competitive, we need to stay on schedule...
- With first data, more realistic estimation on going
- But also plenty of room for improvements
  - improved methods
  - new modes (some pioneered on Belle data sample)



equations for the rate of events in bins i and -i of the Dalitz plots

 $n_+^2 \left[ \text{GeV}^2 / c^4 \right]$ 

for B<sup>-</sup> and B<sup>+</sup> decays:

$$x_{\pm} \equiv r_B \cos(\delta_B \pm \gamma),$$
  
 $y_{\pm} \equiv r_B \sin(\delta_B \pm \gamma).$ 

$$\Gamma_{+i}(B^{-} \to D(\to K_{\rm S}^{0}h^{+}h^{-})K^{-}) \propto \left[T_{+i} + (x_{-}^{2} + y_{-}^{2})T_{-i} + 2\sqrt{T_{+i}T_{-i}}(x_{-}c_{+i} + y_{-}s_{+i})\right],$$

$$\Gamma_{-i}(B^{-} \to D(\to K_{\rm S}^{0}h^{+}h^{-})K^{-}) \propto \left[T_{-i} + (x_{-}^{2} + y_{-}^{2})T_{+i} + 2\sqrt{T_{+i}T_{-i}}(x_{-}c_{-i} + y_{-}s_{-i})\right],$$

$$\Gamma_{+i}(B^{+} \to D(\to K_{\rm S}^{0}h^{+}h^{-})K^{+}) \propto \left[T_{-i} + (x_{+}^{2} + y_{+}^{2})T_{+i} + 2\sqrt{T_{+i}T_{-i}}(x_{+}c_{+i} - y_{+}s_{+i})\right],$$

$$\Gamma_{-i}(B^{+} \to D(\to K_{\rm S}^{0}h^{+}h^{-})K^{+}) \propto \left[T_{+i} + (x_{+}^{2} + y_{+}^{2})T_{-i} + 2\sqrt{T_{+i}T_{-i}}(x_{+}c_{-i} - y_{+}s_{-i})\right].$$

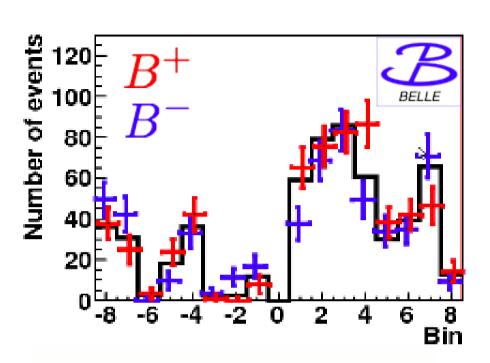
parameters  $T_{\pm i}$  can be determined by measuring decay rates of flavour-tagged  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$  decays, i.e. where D meson can be identified as  $D^0$  or  $\overline{D}^0$ 

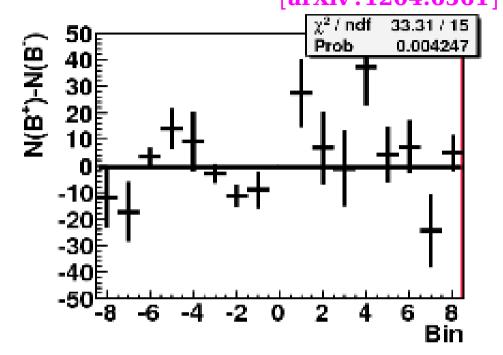
measuring B $\rightarrow$ DK decay rates in each bin, 2k+3 unknowns =  $\mathbf{c_i}$ ,  $\mathbf{s_i}$ ,  $\mathbf{r_B}$ ,  $\delta_{\mathbf{B}}$  and  $\gamma$  k  $\geq$  2: greater number of equations than unknowns and  $\gamma$  can be determined preferable to perform dedicated measurements of  $\mathbf{c_i}$  and  $\mathbf{s_i}$ , use them as inputs

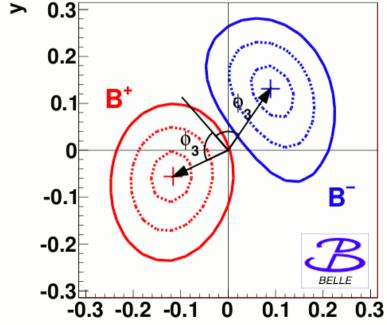
#### Binned Dalitz method result in B → DK

X

 $772 \, M \, B \overline{B}$  PRD 85, 112014 (2012) [arXiv:1204.6561]







$$\gamma = (77.3^{+15.1}_{-14.9} \pm 4.1 \pm 4.3)^{\circ}$$

$$r_B = 0.145 \pm 0.030 \pm 0.010 \pm 0.011$$

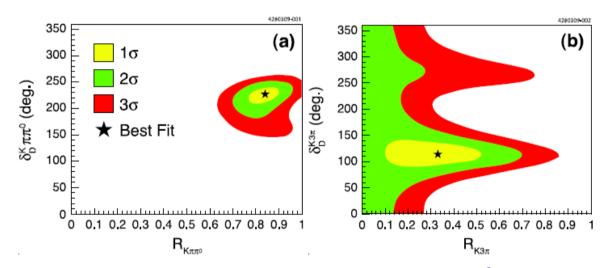
$$\delta_{\rm B} = (129.9 \pm 15.0 \pm 3.8 \pm 4.7)^{\circ}$$

uncertainty in c<sub>i</sub>, s<sub>i</sub> from CLEO data size (can be reduced using future BES-III data)

# quasi-GLW, quasi-ADS...

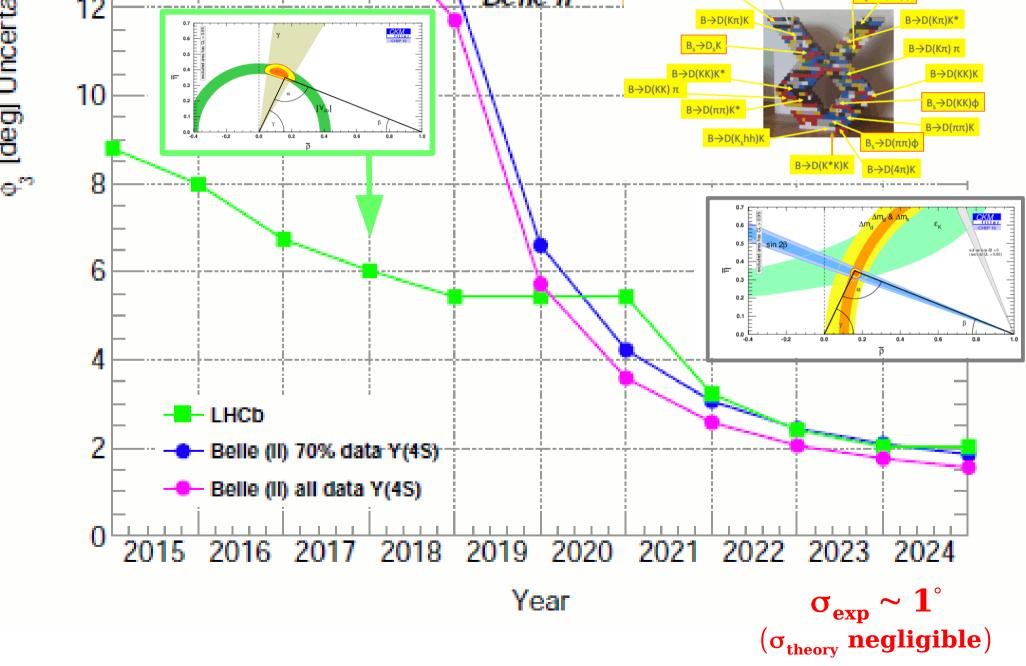
certain multi-body decays are almost pure CP-eigenstates:  $\Rightarrow$  quasi-GLW, for example for D $\rightarrow$ 4 $\pi$ , 2F $_{+}$ -1 = 0.737 ± 0.028

other like ADS modes: for example  $D \rightarrow K \pi \pi^0$ , coherence factor  $\sim 1$ 



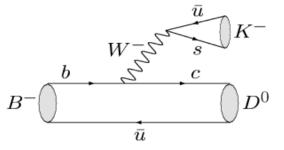
yields of double-tagged events where one meson decays into  $K^-\pi^+\pi^0$  (or  $K3\pi$ ), and the other meson decays into CP-odd, CP-even and  $K\pi$ 

Ultimate y-from-tree decays precision will be reached through many individual measurements  $\Lambda_b \rightarrow DKp$ ္နဲ့ [deg] Uncertainty  $B \rightarrow D(K\pi)K \pi\pi$  $B \rightarrow D(K3\pi)K$ Belle II 12  $B \rightarrow D(K\pi)K$  $B \rightarrow D(K\pi)K^*$  $B_s \rightarrow D_s K$  $B \rightarrow D(K\pi) \pi$  $B \rightarrow D(KK)K^*$  $B \rightarrow D(KK)K$ 10  $B_s \rightarrow D(KK) \varphi$  $B \rightarrow D(\pi\pi)K$  $B \rightarrow D(\pi\pi)K$  $B \rightarrow D(K_c hh)K$  $B_c \rightarrow D(\pi\pi) \Phi$ 8 6 4 - LHCb 2 Belle (II) 70% data Y(4S): Belle (II) all data Y(4S)



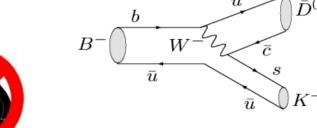
## $\gamma$ measurements from $B^{\pm} \rightarrow DK^{\pm}$

- Theoretically pristine B→DK approach
- ∘ Access  $\gamma$  via interference between  $B^- \rightarrow D^0 K^-$  and  $B^- \rightarrow \overline{D}^0 K^-$





color allowed  $B^- \rightarrow D^0 K^- \sim V_{cb} V_{us}^*$  $\sim A \lambda^3$ 

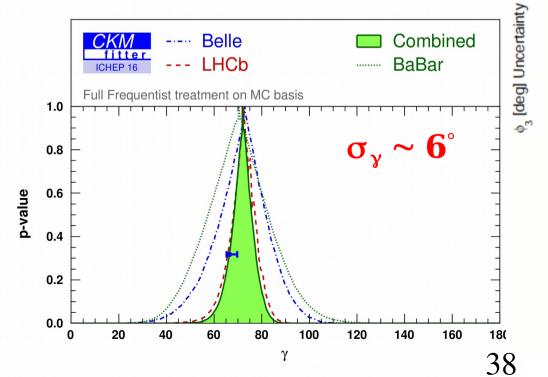


relative weak phase is  $\gamma$  relative strong phase is  $\delta_B$ 

 $r_{\rm B} \simeq 0.1$ 

color suppressed  $B^- \rightarrow \overline{D}^0 K^- \sim V_{ub} V_{cs}^*$ 

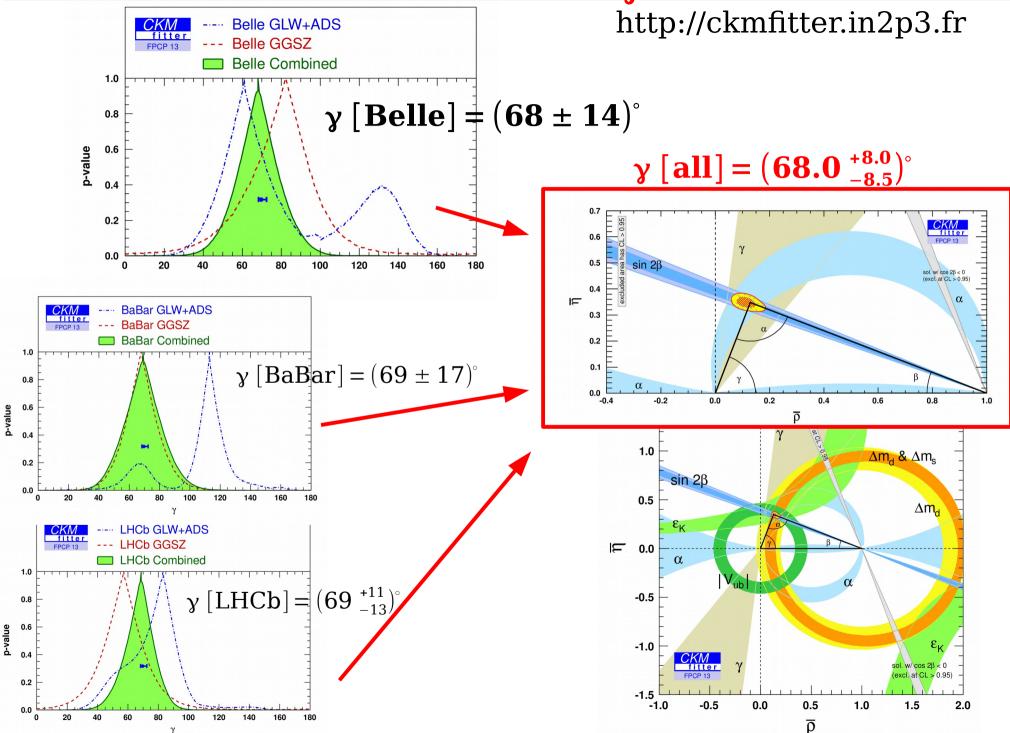
 $\sim A \lambda^3 (\rho + i \eta)$ 





long way to go ... ( $\stackrel{\mathsf{Year}}{\to} \sigma_{\mathsf{y}} = 1^{\circ} \text{ or less ?}$ )

### Combined measurements for y from all methods



# Charm mixing in $D^0 \rightarrow K^+\pi^-$

The ratio R(t) of WS  $D^{*+} \rightarrow D^0 \pi^+ \rightarrow K^+ \pi^- \pi^+$  to RS  $D^{*+} \rightarrow D^0 \pi^+ \rightarrow K^- \pi^+ \pi^+$ decay rates can be approximated (assuming |x|,  $|y| \le 1$  and no CPV) by:

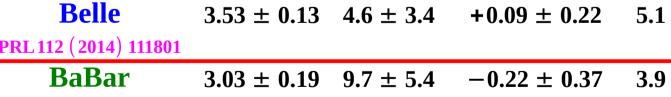
$$R(t) = R_{D} + \sqrt{R_{D}} y't + \frac{x'^{2} + y'^{2}}{4} t^{2}$$

$$= x \cos \delta + y \sin \delta \qquad \text{So a streng where differences}$$

$$x' = x \cos \delta_{K\pi} + y \sin \delta_{K\pi}$$
$$y' = y \cos \delta_{K\pi} - x \sin \delta_{K\pi}$$

 $\delta_{K_{\pi}}$ : strong phase difference btw DCS and CF amplitudes

Exp	$\mathbf{R}_{\mathbf{D}}$	$\mathbf{y}$	$\mathbf{x}^{'2}$	Σ
	$(10^{-3})$	$\left(10^{-3}\right)$	$\left(10^{-3}\right)$	



PRL 98 (2007) 211802

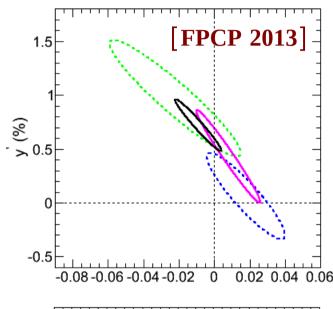
LHCb 
$$3.57 \pm 0.07 + 4.8 \pm 1.0 + 0.055 \pm 0.049$$
 ? PRL 111 (2013) 251801

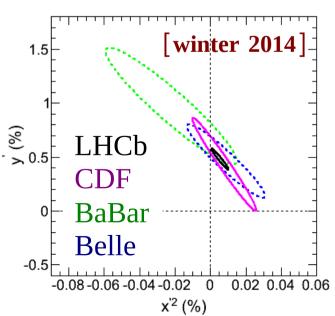
preliminary (2013)

 $3.51 \pm 0.35 \quad 4.3 \pm 4.3$ 

 $+0.08 \pm 0.18$ 

6.1





## **ADS observables**

•  $(R_+, R_-)$  instead of  $(R_{ADS}, A_{ADS})$  whenever available Effect of D- $\overline{D}$  mixing on  $\gamma$ 

- M.Rama, arXiv:1307.4384
- $\circ \ R^{\mp} = r_{B}^{2} + r_{D}^{2} + 2r_{B}r_{D}\cos(\delta_{B}\mp\gamma + \delta_{D})$   $\rightarrow \ R^{\mp} = r_{B}^{2} + r_{D}^{2} + 2r_{B}r_{D}\cos(\delta_{B}\mp\gamma + \delta_{D}) yr_{D}\cos\delta_{D} yr_{B}\cos(\delta_{B}\mp\gamma) + xr_{D}\sin\delta_{D} xr_{B}\sin(\delta_{B}\mp\gamma)$
- tried on the current LHCb average (DK): ~ 1 degree difference

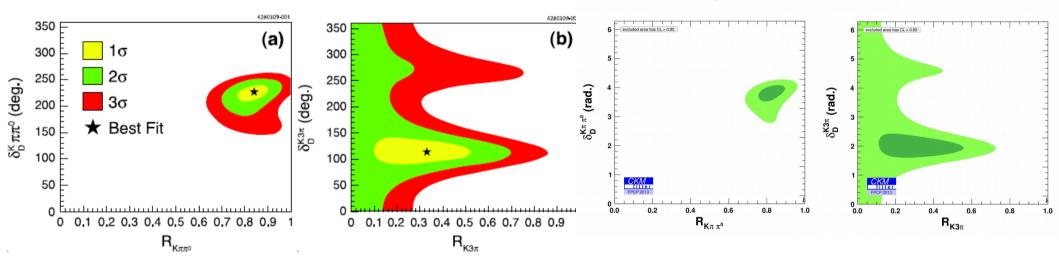
### $K\pi\pi^0$ , $K3\pi$ from CLEO-c

yields of double-tagged events where one meson decays into  $K^-\pi^+\pi^0$  (or  $K3\pi$ ), and the other meson decays into CP-odd, CP-even and  $K\pi$ 

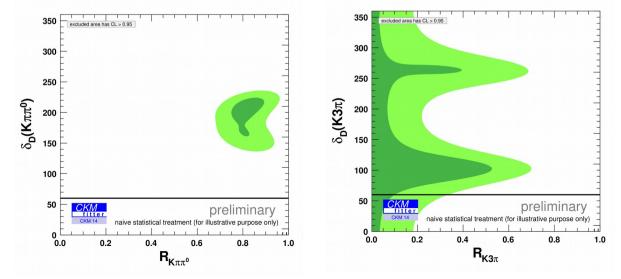
[arXiv:0903.4853, N.Lowrey et al]

(combined with external inputs: x, y,  $\delta_{D}$ ...)

that we could reproduce earlier extending the charm fitter (+ Br's)



### 2014 version (currently used in our γ combination):

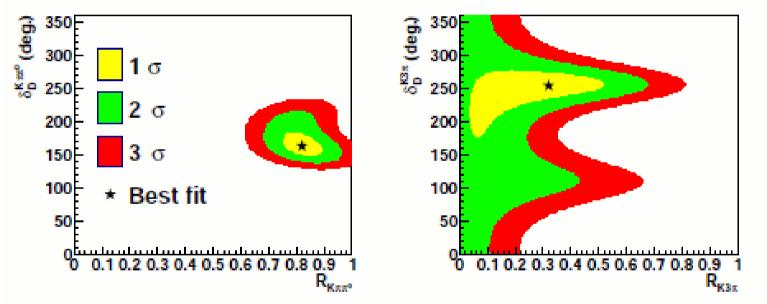


### $\mathbf{K}$ $\mathbf{\pi}$ $\mathbf{\pi}^{\mathbf{0}}$ , $\mathbf{K}$ $\mathbf{3}$ $\mathbf{\pi}$ from CLEO-c [J.Libby et al, arXiv:1401.1904]

yields of double-tagged events where one meson decays into  $K^-\pi^+\pi^0$  (or  $K3\pi$ ), and the other meson decays into  $K_S^0\pi^+\pi^-$ 

$$Y_i = H_{K\pi\pi^0} \Big( K_i + (r_D^{K\pi\pi^0})^2 K_{-i} - \frac{2r_D^{K\pi\pi^0} \sqrt{K_i K_{-i}} R_{K\pi\pi^0} [c_i \cos \delta_D^{K\pi\pi^0} + s_i \sin \delta_D^{K\pi\pi^0}] \Big),$$
 measure by CLEO-c

K<sub>i</sub>: fractional yield of D<sup>0</sup> decays that fall into bin i



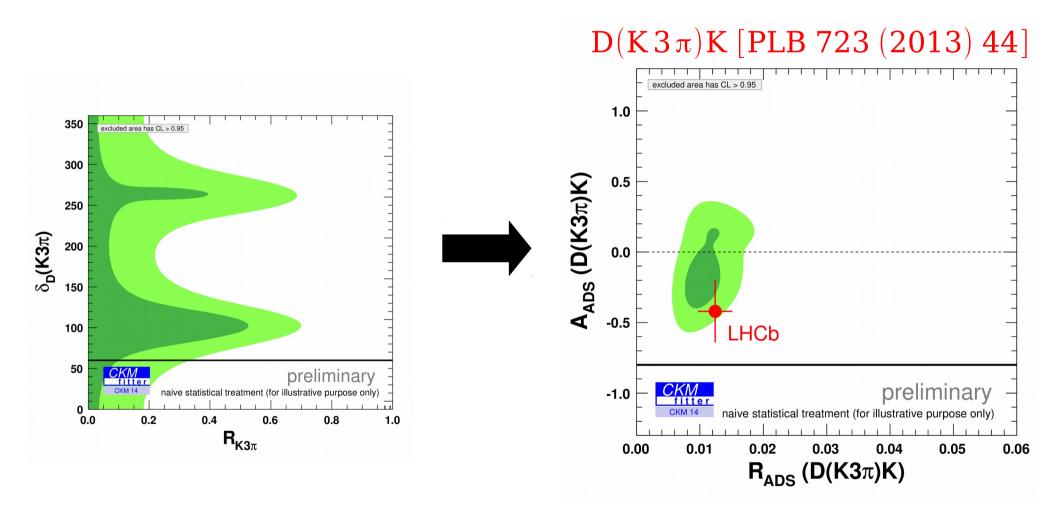
⇒ will soon include this information

#### $K3\pi$ charm information is limited:

- possible additional inputs from BES III
- B factories/LHCb[S.Harnew and J.Rademacker, arXiv:1309.0134]

## ADS $B \rightarrow D(K3\pi)K$

where ''expectations'' derived from the GGSZ observables,  $\delta_D$ ,  $r_D$  and R (for K3 $\pi$ )



 $\Rightarrow$  D(K3 $\pi$ )K LHCb result included in the  $\gamma$  combination

## ADS $B \rightarrow D(K\pi\pi^0)K$

P

R. M

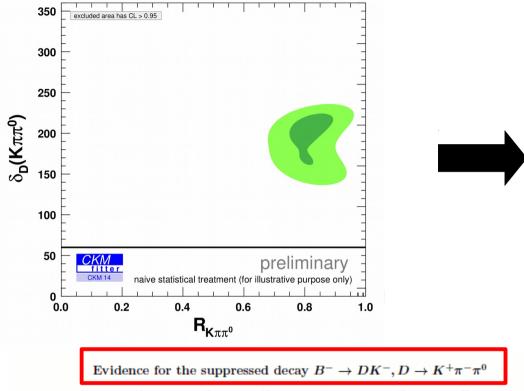
P. Pa

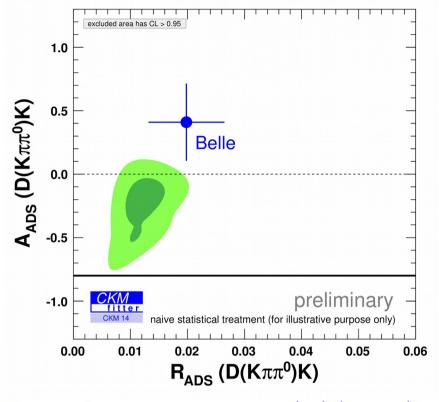
V. Sa

M. S

A. S

where ''expectations'' derived from the GGSZ observables,  $\delta_D$ ,  $r_D$  and R (for  $K\pi\pi^0$ )





M. Nayak, <sup>16</sup> J. Libby, <sup>16</sup> K. Trabelsi, <sup>12</sup> I. Adachi, <sup>12</sup> H. Aihara, <sup>55</sup> D. M. Asner, <sup>42</sup> T. Aushev, <sup>20</sup> A. M. Bakich, <sup>49</sup> A. Bala, <sup>43</sup> P. Behera, <sup>16</sup> K. Belous, <sup>18</sup> V. Bhardwaj, <sup>34</sup> G. Bonvicini, <sup>60</sup> A. Bozek, <sup>38</sup> M. Bračko, <sup>27</sup>, <sup>21</sup> T. E. Browder, <sup>11</sup> D. Červenkov, <sup>5</sup> M.-C. Chang, <sup>8</sup> P. Chang, <sup>37</sup> V. Chekelian, <sup>28</sup> A. Chen, <sup>35</sup> B. G. Cheon, <sup>10</sup> R. Chistov, <sup>20</sup> I.-S. Cho, <sup>62</sup> K. Cho, <sup>24</sup> V. Chobanova, <sup>28</sup> Y. Choi, <sup>48</sup> D. Cinabro, <sup>60</sup> J. Dalseno, <sup>28</sup>, <sup>51</sup> M. Danilov, <sup>20</sup>, <sup>30</sup> Z. Doležal, <sup>5</sup> Z. Drásal, <sup>5</sup>

DK [PRD 88, 091104(R) (2013)]

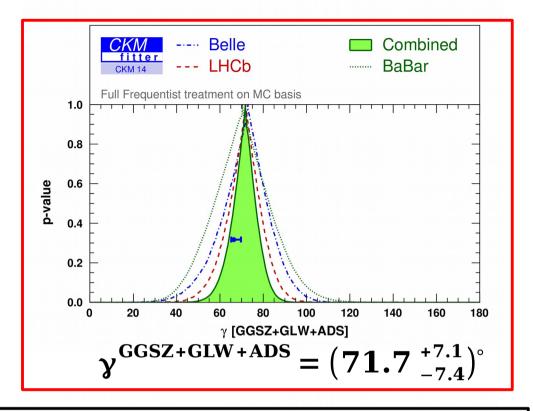
We report a study of the suppressed decay  $B^- \to DK^-, D \to K^+\pi^-\pi^0$ , where D denotes either a  $D^0$  or a  $\overline{D}^0$  meson. The decay is sensitive to the CP-violating parameter  $\phi_3$ . Using a data sample of  $772 \times 10^6$   $B\overline{B}$  pairs collected at the  $\Upsilon(4S)$  resonance with the Belle detector, we measure the ratio of branching fractions of the above suppressed decay to the favored decay  $B^- \to DK^-, D \to K^-\pi^+\pi^0$ . Our result is  $R_{DK} = [1.98 \pm 0.62 ({\rm stat.}) \pm 0.24 ({\rm syst.})] \times 10^{-2}$ , which indicates the first evidence of the signal for this suppressed decay with a significance of 3.2 standard deviations. We measure the direct CP asymmetry between the suppressed  $B^-$  and  $B^+$  decays to be  $A_{DK} = 0.41 \pm 0.30 ({\rm stat.}) \pm 0.05 ({\rm syst.})$ . We also report measurements for the analogous quantities  $R_{D\pi}$  and  $A_{D\pi}$  for the decay  $B^- \to D\pi^-$ ,  $D \to K^+\pi^-\pi^0$ .

G. Varner, <sup>11</sup> K. E. Varvell, <sup>49</sup> M. N. Wagner, <sup>9</sup> C. H. Wang, <sup>36</sup> M.-Z. Wang, <sup>37</sup> Y. Watanabe, <sup>22</sup> K. M. Williams, <sup>59</sup> F. Wong, <sup>25</sup> V. Varnashita, <sup>39</sup> S. Vashchonko, <sup>7</sup> V. Varna, <sup>40</sup> V. Zhilloh, <sup>4</sup> W. Zhulanov, <sup>4</sup> Ond. A. Zunane, <sup>23</sup>

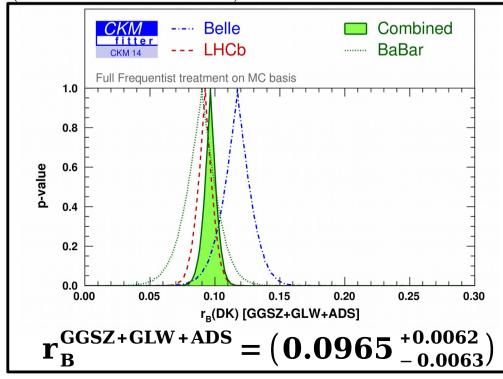
 $\overset{\text{E. Won,}^{25}}{\Rightarrow}\overset{\text{Y. Yamashita,}^{39}}{\text{Belle}}(\overset{\text{S. Yashchenko}}{\text{and}}\overset{\text{Y. Yusa,}^{40}}{\text{Ba}}\overset{\text{V. Zhillanov,}^{40}}{\text{nd}}\overset{\text{N. Zhullanov,}^{40}}{\text{nd}}\overset{\text{N. Zhullanov,}^{40}}{\text{N. Tessults included in the }}\chi \text{ combination}$ 

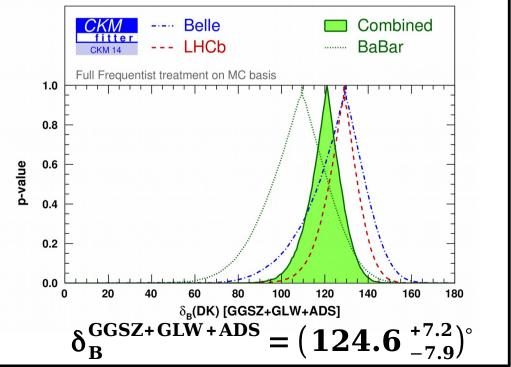
## **GGSZ+GLW+ADS**

+20 obs.



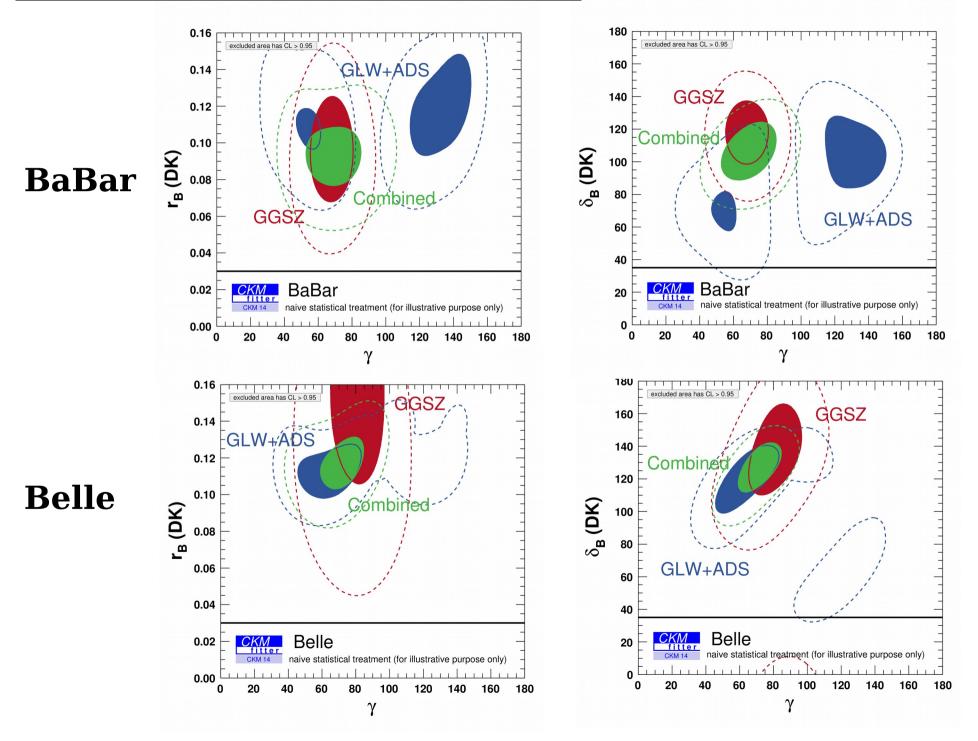
(results for DK)



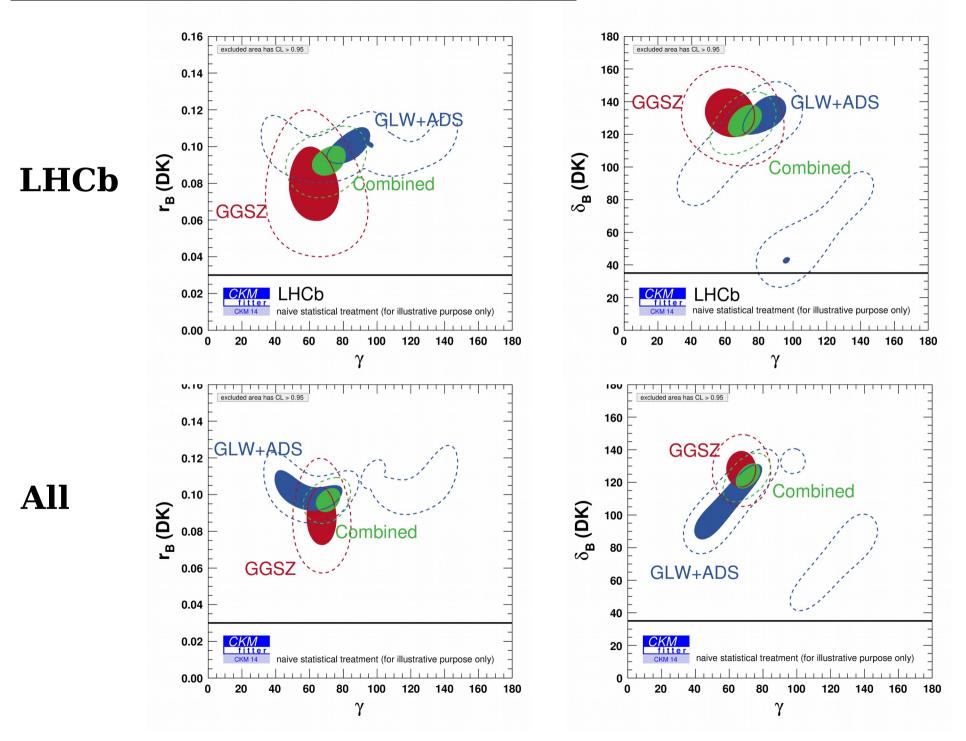


## **GGSZ versus GLW+ADS**

 $(\mathbf{r}_{\mathbf{B}}(\mathbf{D}\mathbf{K}) \mathbf{vs} \boldsymbol{\gamma}, \delta_{\mathbf{B}}(\mathbf{D}\mathbf{K}) \mathbf{vs} \boldsymbol{\gamma})$ 



## GGSZ versus GLW+ADS $(r_B(DK) \text{ vs } \gamma, \delta_B(DK) \text{ vs } \gamma)$



## The small $r_B$ issue

clearly in the  $r_B \rightarrow 0$  limit the interference disappears and there is no sensitivity to the phase  $\gamma$ 

when the true value of  $r_B$  is small, then the distribution of  $\hat{r_B}$  best fit values for randomly generated data is biased towards larger values, until the experimental errors are sufficiently small to exclude the  $r_B \sim 0$  region

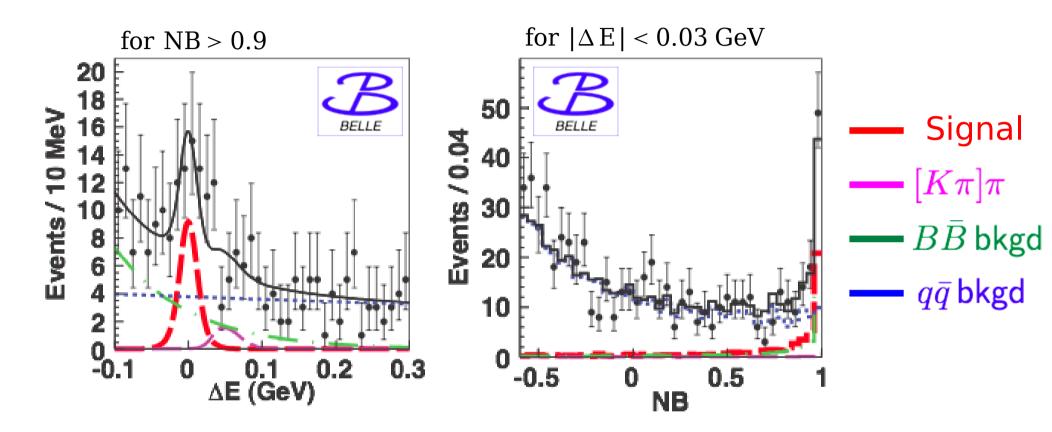
on the other hand the error on  $\gamma$  is roughly proportional to  $1/r_B$ , hence for small r<sub>B</sub> it is biased towards smaller values

in the language of frequentist statistics it means that the usual  $\Delta \ln \mathcal{L} = 1/2$  rule does not work here, the 68%CL interval extracted from it does not cover the true value of  $\gamma$  at 68% frequency (undercoverage)

to correct for this effect one has to compute the actual distribution of the profile log-likelihood, and from that distribution deduce a p-value or a CL interval

problem: as soon as the log-likelihood is not distributed as a  $\chi^2$ , its distribution a priori depends on the nuisance parameters, namely  $r_B$ ,  $\delta_B$  etc.

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 $56.0^{+15.1}_{-14.2}$  events

$$\mathbf{R}_{\mathrm{DK}} = (\mathbf{1.63}^{+0.44}_{-0.41}^{+0.07}) \times \mathbf{10}^{-2}$$

$$A_{\rm DK} = -0.39^{+0.26}_{-0.28}^{+0.04}_{-0.03}$$

First evidence obtained with a significance of  $4.1\sigma$  (including syst.)