



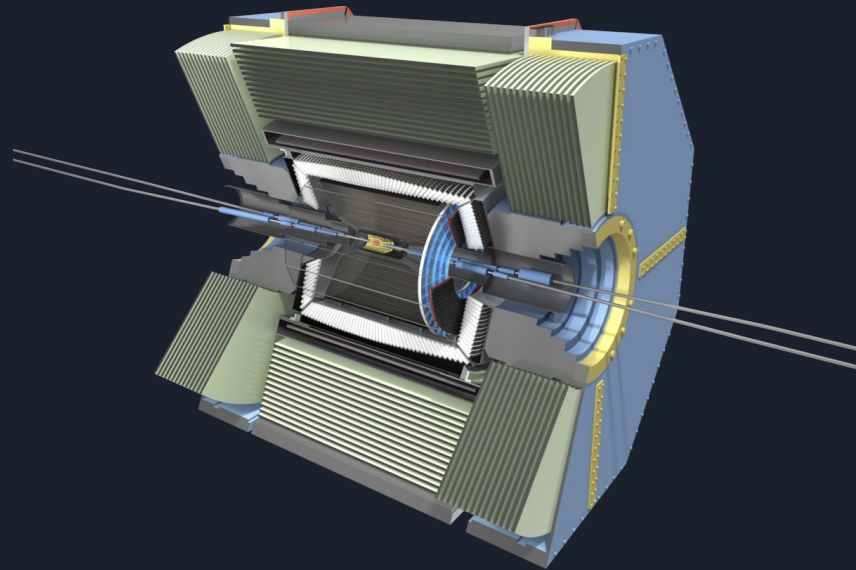
Time-dependent CP violation in $\mathbf{b \rightarrow s \gamma}$ transitions at Belle II



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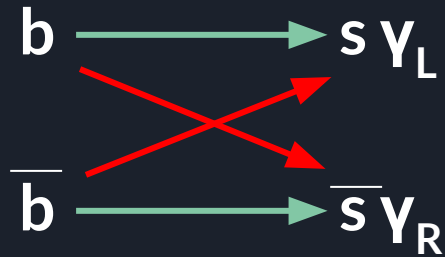
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3. Analysis strategies
4. Electroweak FCNCs
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6. K_S^0 reconstruction
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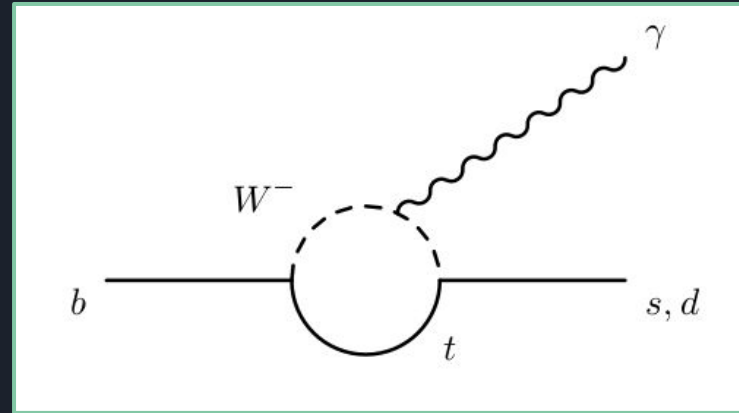


Motivation:

In Standard Model $b \rightarrow s \gamma$ depend on b flavor



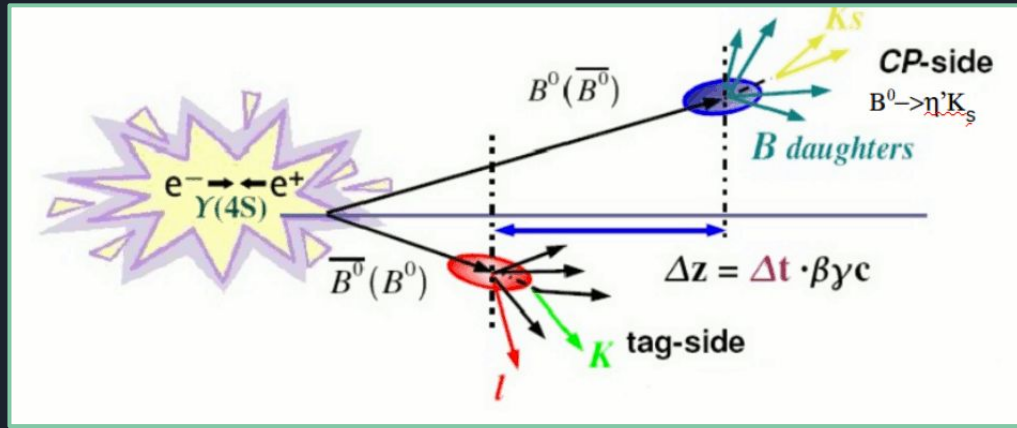
- ❖ allowed
- ❖ suppressed by m_s/m_b



Presence of significant mixing-induced CP violation would indicate the presence of right handed currents and clear hint of new physics.

- This type of new physics does not require a new phase.

Time-dependant evolution at Belle II



$(\Delta z)_{\text{Belle2}} \sim 130 \mu\text{m}$
 $(\Delta t)_{\text{Belle2}} \sim 0.94 \text{ ps}$

$$a_{f_{CP}}(\Delta t) \equiv \frac{\Gamma[B(\Delta t)] - \Gamma[\bar{B}(\Delta t)]}{\Gamma[B(\Delta t)] + \Gamma[\bar{B}(\Delta t)]} = C \cos(\Delta M \Delta t) - S \sin(\Delta M \Delta t)$$

Direct CPV (blue oval) and mixing induced CPV (green oval) are indicated by arrows pointing to the C and S terms in the equation above.

1. Fully reconstructed one of B mesons which decays to CP eigenstates
2. Tag-side determines its flavour (efficiency $\approx 30\%$)
3. Proper time (Δt) is measured from decay-vertex difference (Δz).

Inclusive analysis strategies for $b \rightarrow X_s \gamma$

Fully inclusive:

- Exploit clean decay environment on Belle II
- Can be fully hadronic tag (have full event information)
- ...or semi-leptonic tag (don't have full event)

&

Sum-of-exclusives:

- Reconstruct, the 'X' from many exclusive decays:
 $X_s \rightarrow K\eta\pi, 3K\eta\pi, K\eta\pi\pi$
 $(n>1, m\geq 1)$
- Know flavor of B
- Know isospin

reco. method	tagging	effi.	S/B	q	p_B	A_{CP}	Δ_{0+}	ΔA_{CP}
sum-of-exclusive	none	high	moderate	s or d	yes	yes	yes	yes
fully-inclusive	had. B	very low	very good	s and d	yes	yes	yes	yes
	SL B	very low	very good	s and d	no	yes	yes	yes
	L	moderate	good	s and d	no	yes	no	no
	none	very high	very bad	s and d	no	no	no	no

Electroweak FCNCs

$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{ts}^* V_{tb} \sum_{i=1}^{10} C_i(\mu) O_i(\mu)$$

Wilson coefficients

(calculated perturbatively; encode short-distance physics)

We can write the amplitude including RH contribution as:

$$\mathcal{M}(b \rightarrow s\gamma) \simeq -\frac{4G_F}{\sqrt{2}} V_{ts}^* V_{tb} \left[\underbrace{(C_{7\gamma}^{\text{SM}} + C_{7\gamma}^{\text{NP}})}_{\propto \mathcal{M}_L} \langle O_{7\gamma} \rangle + \underbrace{C'_{7\gamma}{}^{\text{NP}} \langle O'_{7\gamma} \rangle}_{\propto \mathcal{M}_R} \right]$$

Products of field operators

(non-perturbative hadronic matrix elements; Heavy quark expansion in inverse powers of m_b)

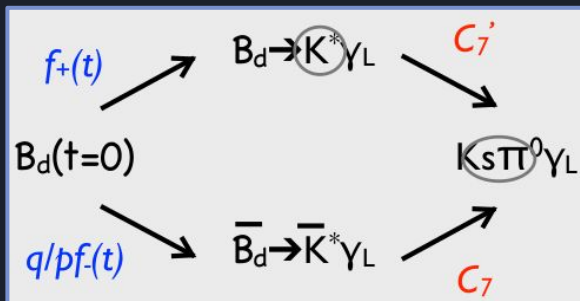
We have a constraint from inclusive branching ratio measurement:

$$Br(B \rightarrow X_S \gamma) \propto |C_{7\gamma}^{\text{SM}} + C_{7\gamma}^{\text{NP}}|^2 + |C'_{7\gamma}{}^{\text{NP}}|^2$$

The polarization measurement carries information on

$$\frac{\mathcal{M}_R}{\mathcal{M}_L} \simeq \frac{C'_{7\gamma}{}^{\text{NP}}}{C_{7\gamma}^{\text{SM}} + C_{7\gamma}^{\text{NP}}}$$

Time dependent CP violation analysis of $B^0 \rightarrow K_S^0 \pi^0 \gamma$



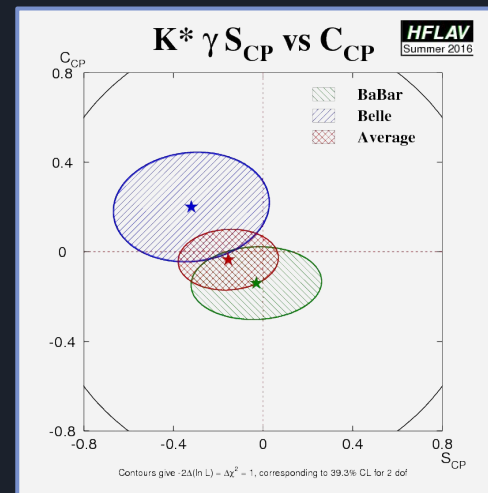
Prediction:

$$S_{K_S^0 \pi^0 \gamma}^{\text{SM}} \sim -2 \frac{m_s}{m_b} \sin 2\phi_1 = -(2.3 \pm 1.6)\%$$

Current world average:

$$S_{K_S^0 \pi^0 \gamma}^{\text{exp}} = -0.16 \pm 0.22$$

- $b \rightarrow s \gamma_R$ is helicity suppressed (m_s/m_b) wrt $b \rightarrow s \gamma_L$
- $B^0 \rightarrow f_{CP} \gamma_R$ interferes with $B^0 \rightarrow B^0\text{-bar} \rightarrow f_{CP} \gamma_R$ only for helicity suppressed $b \rightarrow s \gamma_R$ decay
- TDCPV analysis is sensitive to the decay rate of b into “wrongly” polarized γ .
- New physics can enhance the $b \rightarrow s \gamma_R$ decay rate

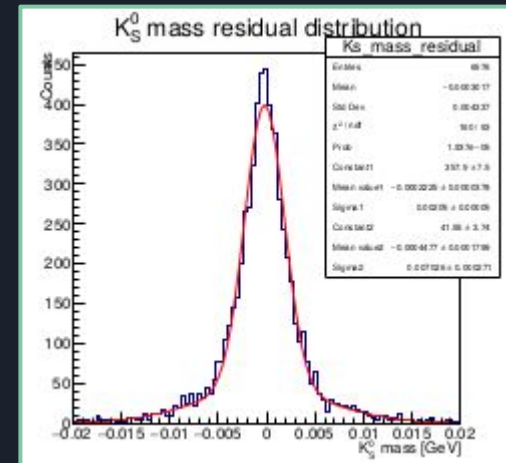
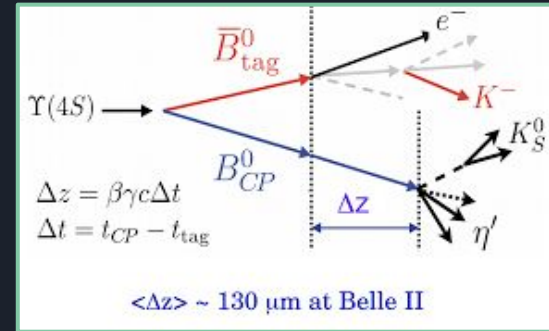


BaBar ($N_{BB} = 467 \cdot 10^6$) [PRD 78 (2008) 071102]

Belle ($N_{BB} = 535 \cdot 10^6$) [PRD 74 (2006) 111104(R)]

K_S^0 reconstruction

- The reconstruction of K_S^0 is the only source of information to determine the vertex position of B_{sig}^0 .
- The channel of K_S^0 reconstruction is $K_S^0 \rightarrow \pi^+ \pi^-$
- The K_S^0 flight direction is extrapolated backwards and matched to the estimated region in which the $e^+ e^-$ collisions take place.

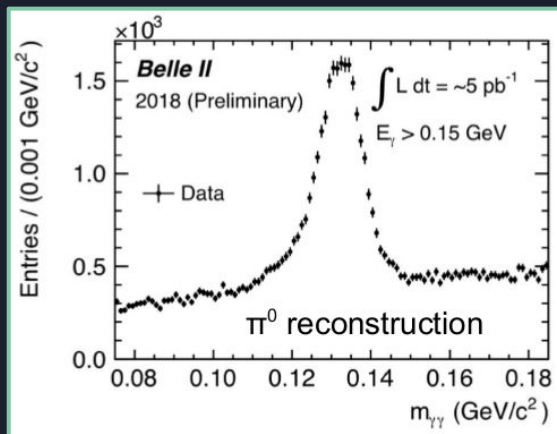
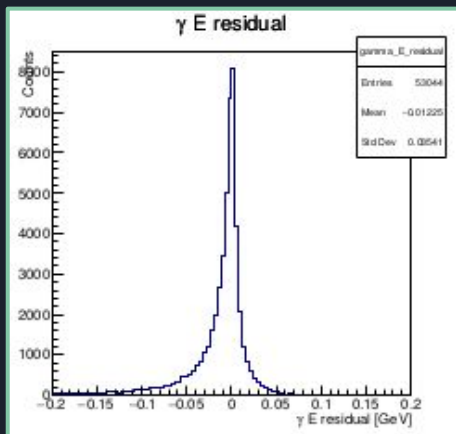


	K_S^0	π^0	γ	B^0
ϵ^{reco}	58.6 %	53.7 %	83.4 %	26.2 %

Photon polarization

Indirect probe of photon polarization by measuring:

$$A_{CP}(t) \equiv \frac{\Gamma(\bar{B}(t) \rightarrow f_{CP}\gamma) - \Gamma(B(t) \rightarrow f_{CP}\gamma)}{\Gamma(\bar{B}(t) \rightarrow f_{CP}\gamma) + \Gamma(B(t) \rightarrow f_{CP}\gamma)} \simeq -2 \left(\frac{m_s}{m_b} \right) \sin(2\beta) \sin(\Delta m \cdot t)$$

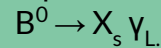


Photon polarization

Standard Model makes definite prediction of photon helicity (D. Atwood et al., Phys. Rev. Lett. 79, 185 (1997)):

- $B^0 \rightarrow X_s \gamma_R$
- $B^0 \rightarrow X_s \gamma_L$

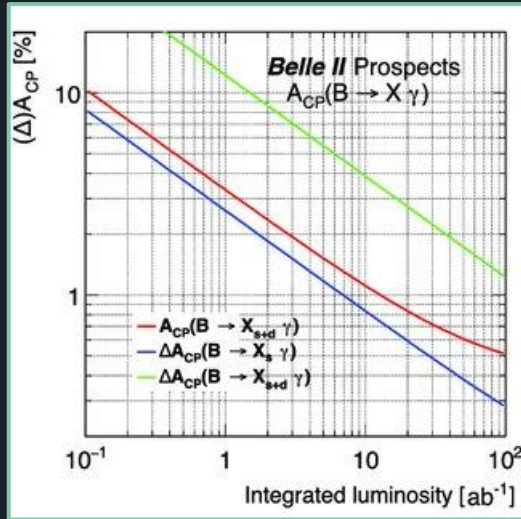
If a helicity flip occurs, the photon will also flip its helicity, producing



No common final state for B^0 and B^0 -bar

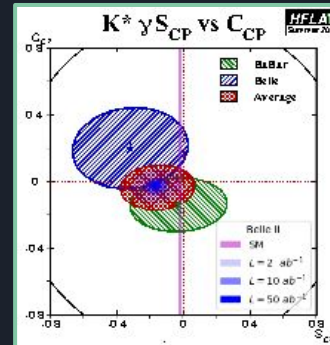
- Suppression of asymmetry S due to interference between B^0 mixing and decay diagrams
- TD CP asymmetry measurements give an indirect measurement of photon polarization.

Predictions for Belle II



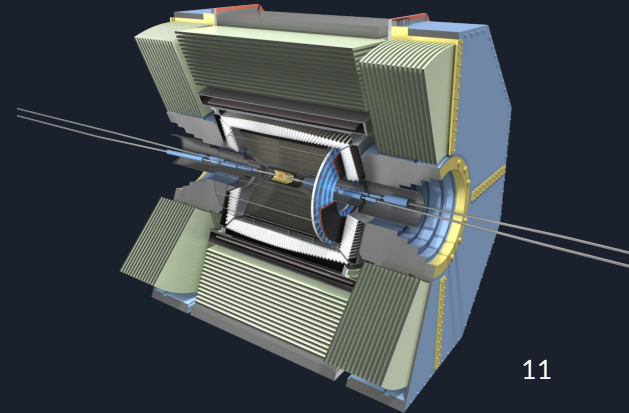
Observables	Belle 0.71 ab ⁻¹	Belle II 5 ab ⁻¹	Belle II 50 ab ⁻¹
$\text{Br}(B \rightarrow X_s \gamma)_{\text{inc}}^{\text{lep-tag}}$	5.3%	3.9%	3.2%
$\text{Br}(B \rightarrow X_s \gamma)_{\text{inc}}^{\text{had-tag}}$	13%	7.0%	4.2%
$\text{Br}(B \rightarrow X_s \gamma)_{\text{sum-of-ex}}$	10.5%	7.3%	5.7%
$\Delta_{0+}(B \rightarrow X_s \gamma)_{\text{sum-of-ex}}$	2.1%	0.81%	0.63%
$\Delta_{0+}(B \rightarrow X_{s+d} \gamma)_{\text{inc}}^{\text{had-tag}}$	9.0%	2.6%	0.85%
$A_{CP}(B \rightarrow X_s \gamma)_{\text{sum-of-ex}}$	1.3%	0.52%	0.19%
$A_{CP}(B^0 \rightarrow X_s^0 \gamma)_{\text{sum-of-ex}}$	1.8%	0.72%	0.26%
$A_{CP}(B^+ \rightarrow X_s^+ \gamma)_{\text{sum-of-ex}}$	1.8%	0.69%	0.25%
$A_{CP}(B \rightarrow X_{s+d} \gamma)_{\text{inc}}^{\text{lep-tag}}$	4.0%	1.5%	0.48%
$A_{CP}(B \rightarrow X_{s+d} \gamma)_{\text{inc}}^{\text{had-tag}}$	8.0%	2.2%	0.70%
$\Delta A_{CP}(B \rightarrow X_s \gamma)_{\text{sum-of-ex}}$	2.5%	0.98%	0.30%
$\Delta A_{CP}(B \rightarrow X_{s+d} \gamma)_{\text{inc}}^{\text{had-tag}}$	16%	4.3%	1.3%

Channel	WA (2017)		5 ab ⁻¹		50 ab ⁻¹	
	$\sigma(S)$	$\sigma(A)$	$\sigma(S)$	$\sigma(A)$	$\sigma(S)$	$\sigma(A)$
$J/\psi K^0$	0.022	0.021	0.012	0.011	0.0052	0.0090
ϕK^0	0.12	0.14	0.048	0.035	0.020	0.011
$\eta' K^0$	0.06	0.04	0.032	0.020	0.015	0.008
ωK_S^0	0.21	0.14	0.08	0.06	0.024	0.020
$K_S^0 \pi^0 \gamma$	0.20	0.12	0.10	0.07	0.031	0.021
$K_S^0 \pi^0$	0.17	0.10	0.09	0.06	0.028	0.018



Summary

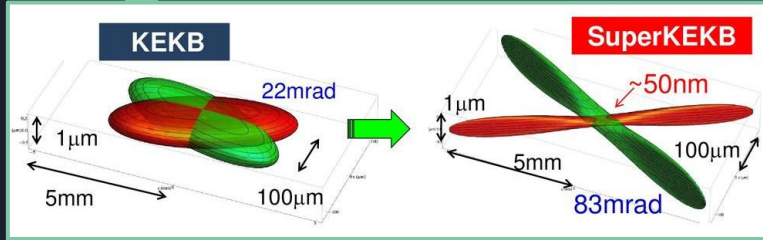
1. Belle II provides a large dataset + improved detector and physics software (Flavor tagging and Vertex reconstruction).
2. Unique possibilities for modes with final state with neutral particles (π^0 , γ)
3. For $b \rightarrow s \gamma$ transitions we can safely assume that all the channels will be dominated by the statistical uncertainties.
4. For most of the penguin dominated modes Belle II is projected to reduce the WA errors by a factor of:
 - 2 to 3 already with 5 ab^{-1}
 - 8 with 50 ab^{-1}





Backup

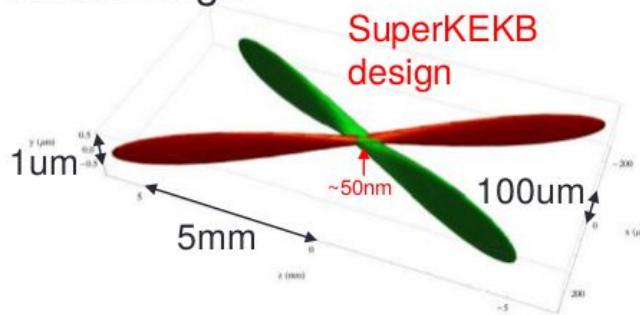
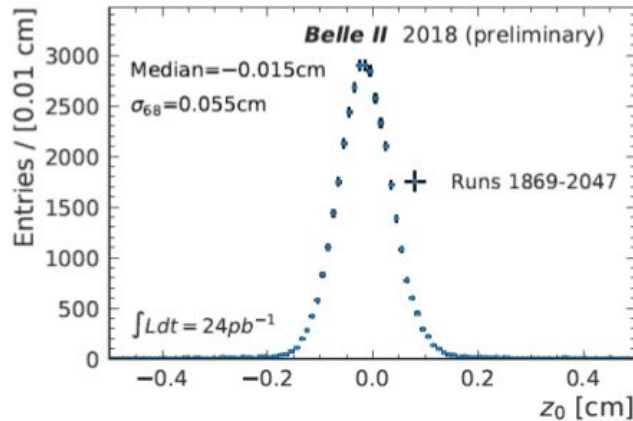
Belle II Nano-Beam



	KEKB Achieved		SuperKEKB	
	LER	HER	LER	HER
RF frequency f [MHz]		508.9		
# of Bunches N		1584	2500	
Horizontal emittance ϵ_x [nm]	18	24	3.2	4.6
Beta at IP β_x^*/β_y^* [mm]		1200/5.9	32/0.27	25/0.30
beam-beam param. ξ_y	0.129	0.090	0.088	0.081
Bunch Length S_z [mm]	6.0	6.0	6.0	5.0
Horizontal Beam Size s_x^* [μm]	150	150	10	11
Vertical Beam Size s_y^* [nm]		0.94	48	62
Half crossing angle ϕ [mrad]		11	41.5	
Beam energy E_b [GeV]	3.5	8	4	7.007
Beam currents I_b [A]	1.64	1.19	3.6	2.6
Lifetime t [min]	133	200	6	6
Luminosity L [$\text{cm}^{-2}\text{s}^{-1}$]		2.1×10^{34}	8×10^{35}	

Interaction vertex

- Distribution of the longitudinal component of the interaction vertex is much smaller than the bunch length



- **The nano beam scheme is working!**

