



Belle(II) XYZ results in charm sector

Chengping Shen

shencp@fudan.edu.cn



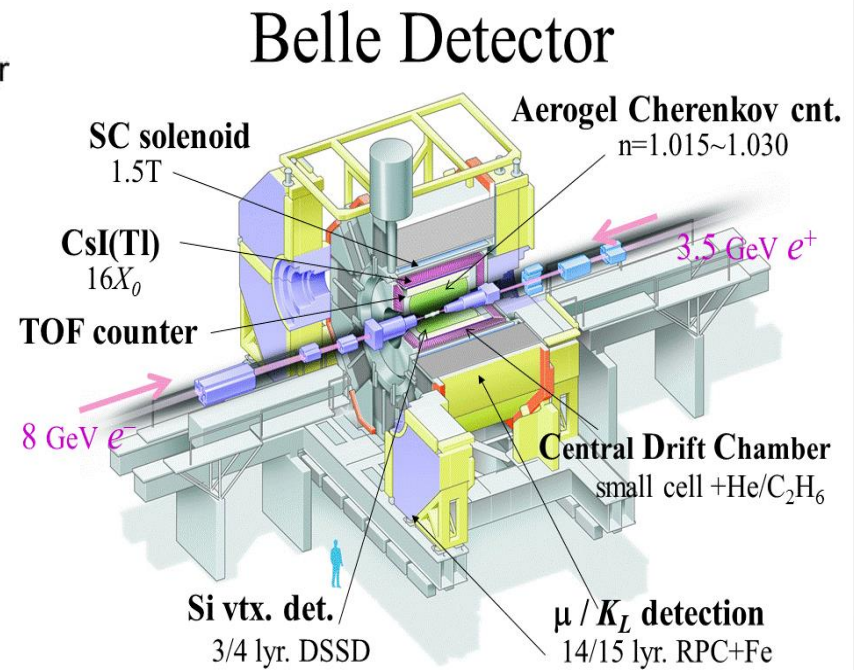
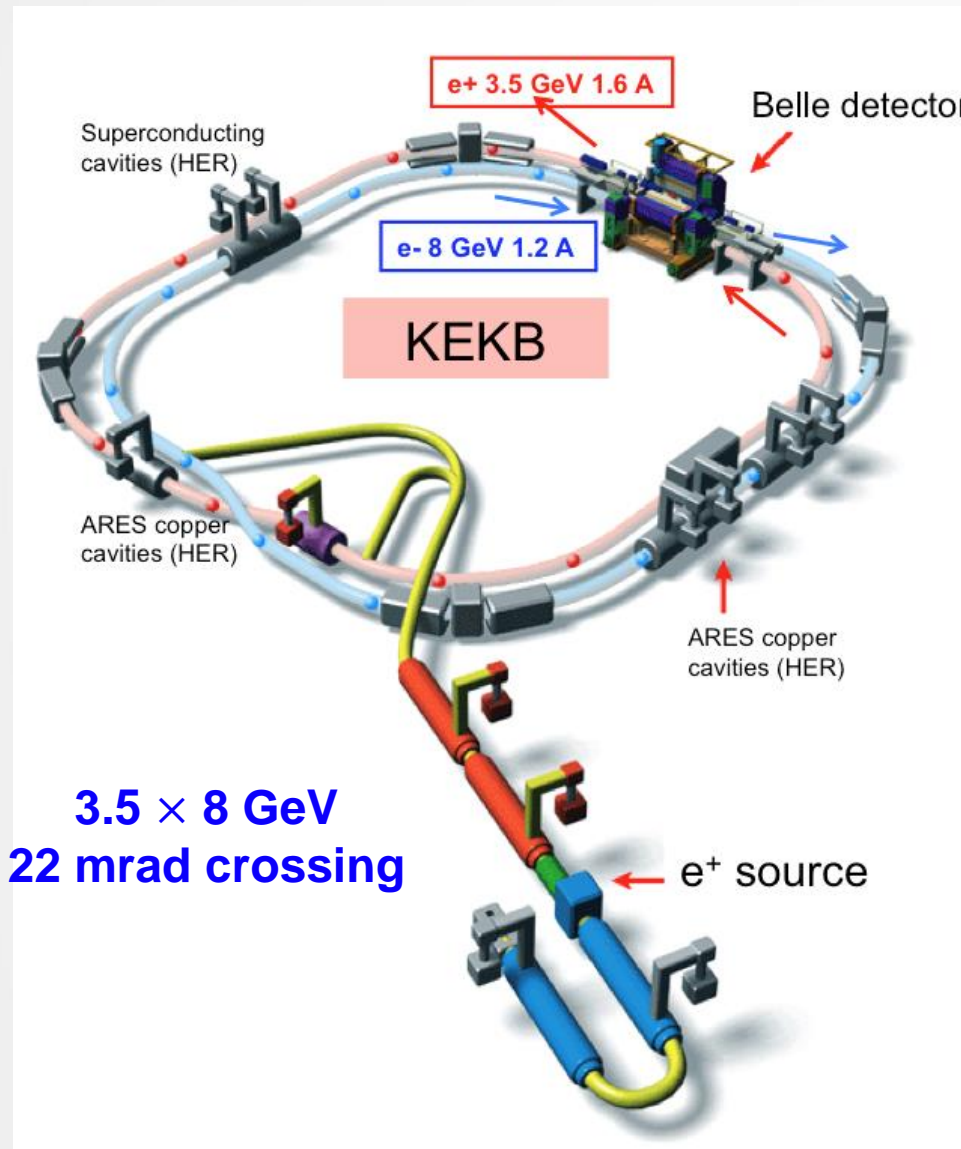
Outline

- Recent results on XYZ results at Belle
- Belle II status
- Potential XYZ results in charm sector at Belle II
- Summary



Success= $X+Y+Z=XYZ$ states

Belle experiment and data samples



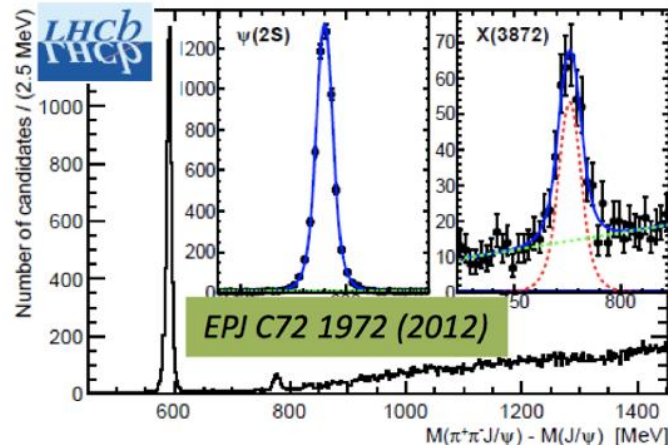
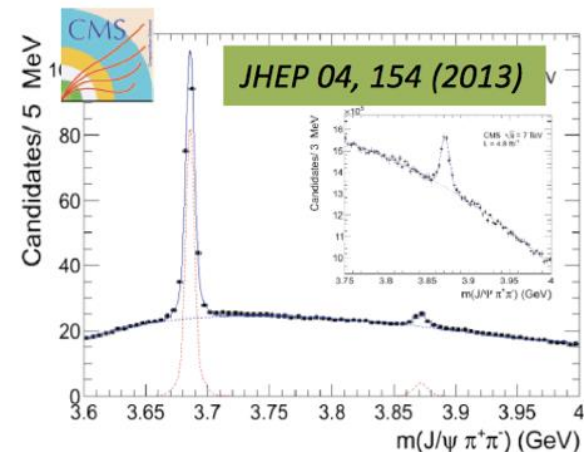
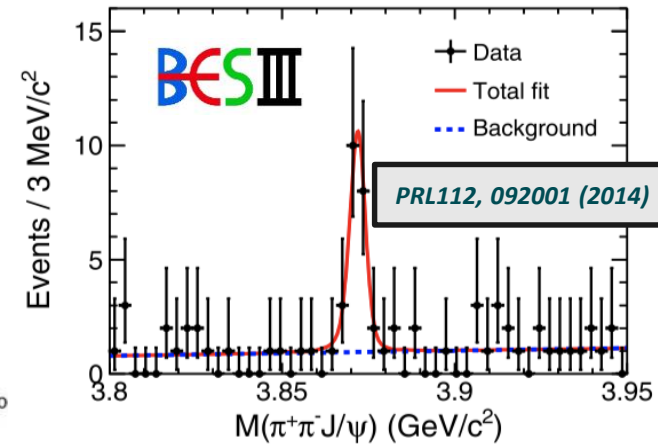
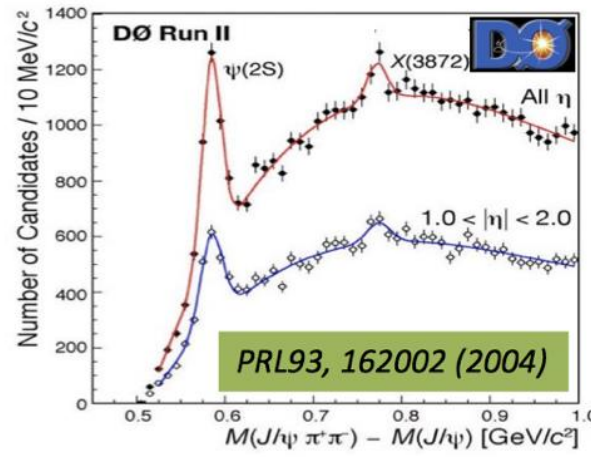
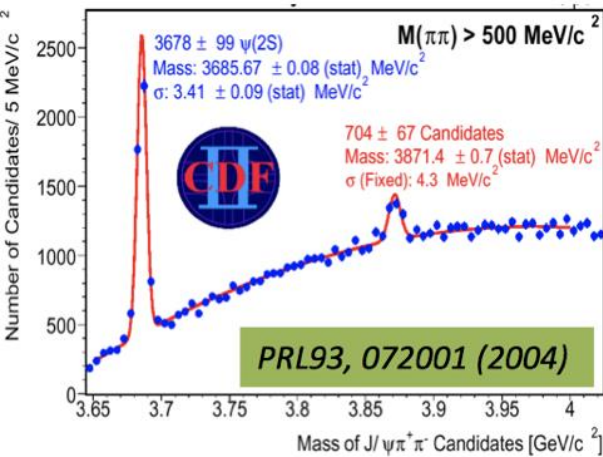
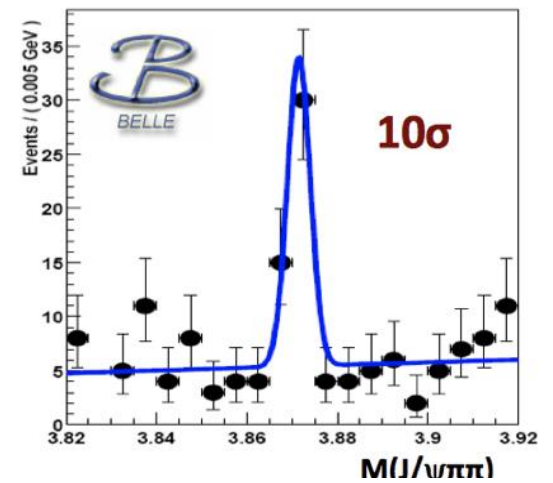
Data taking: 1999 – 2010
On/off/Scan Υ (nS) peaks
Total luminosity: 980 fb⁻¹
772M $B\bar{B}$ events @ Υ (4S)

$X(3872) \rightarrow J/\psi \pi^+ \pi^-$

The most-cited article at Belle: 1900+

First observed by Belle in $B \rightarrow K J/\psi \pi^+ \pi^-$ *PRL91, 262001 (2003)*

- M_X close to $D^0 \bar{D}^{*0}$ threshold $M = (3871.68 \pm 0.17)$ MeV
- Surprisingly narrow: $\Gamma_{\text{tot}} < 1.2$ MeV at 90% C.L.



$X(3872) \rightarrow J/\psi \gamma$: C-even

Angular analysis:

Belle 2006: $J^{PC} = 1^{++}$ or ≥ 2

CDF 2008: $J^{PC} = 1^{++}$ or 2^{-+}

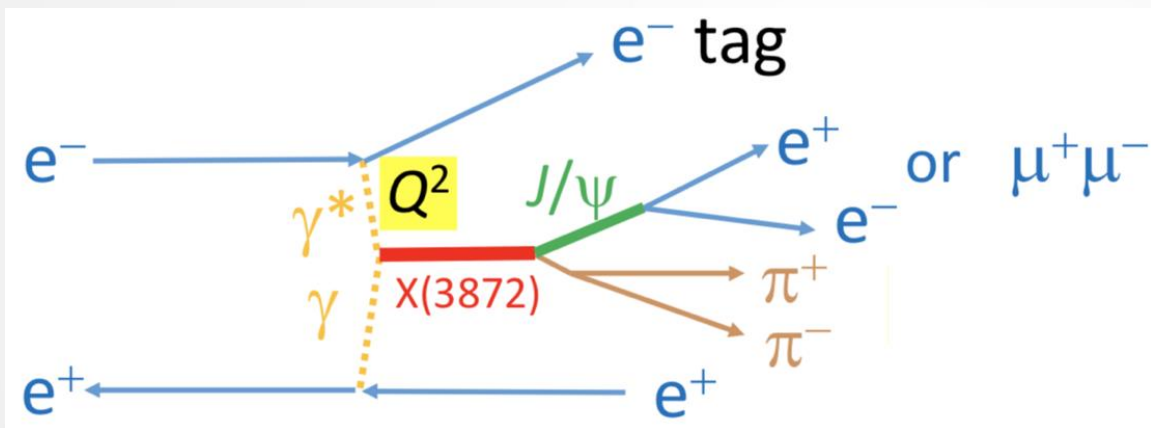
Belle 2011: $J^{PC} = 1^{++}$ or 2^{-+}

LHCb 2013: $J^{PC} = 1^{++}$

Evidence for $X(3872) \rightarrow \pi^+ \pi^- J/\psi$ produced in single-tag two-photon interactions

Phys. Rev. Lett. **126**, 122001 (2021)

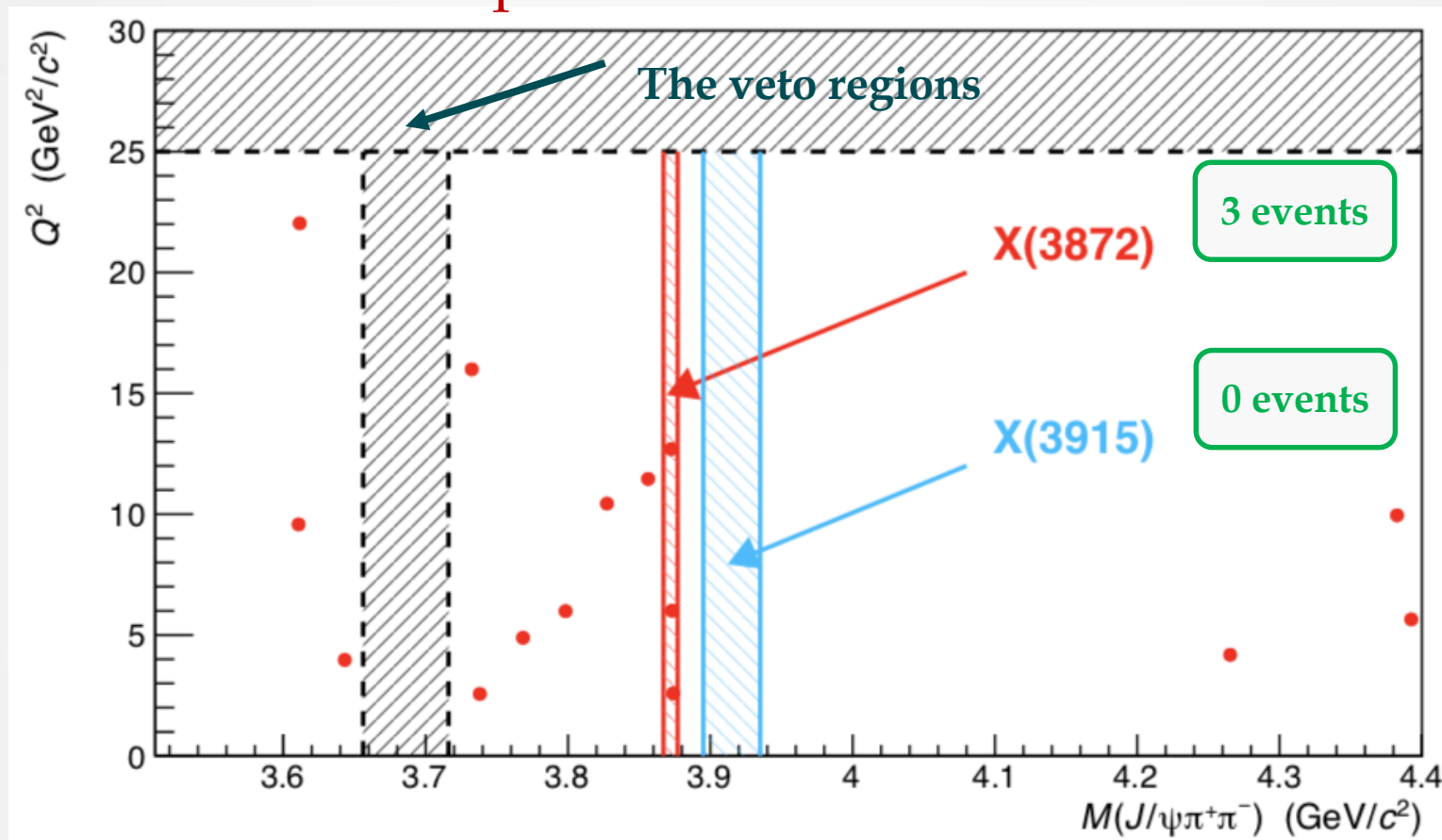
- $X(3872)$ with $J^{PC} = 1^{++}$ is not permitted in $\gamma\gamma \rightarrow X(3872)$, here γ is real.
- $X(3872)$ with $J^{PC} = 1^{++}$ could be produced if one or both photons are highly virtual [Nucl. Phys. B 523, 423 (1998)], i.e. $\gamma\gamma^* \rightarrow X(3872)$, here γ^* is virtual.
- The measurement of $X(3872)$ in two-photon reactions help to understand its internal structure.
- Data sample: 825 fb^{-1} in e^+e^- collisions near 10.6 GeV



$-Q^2$ is the invariant mass-squared of the virtual photon.

If the $X(3872)$ has a molecular component in its structure, it must have a steeper Q^2 dependence than the regular $c \bar{c}$ state. Hence, the single-tag two-photon interactions provide information on the structure of this state. [from discussion with Marek Karliner]

Evidence for $X(3872) \rightarrow \pi^+ \pi^- J/\psi$ produced in single-tag two-photon interactions

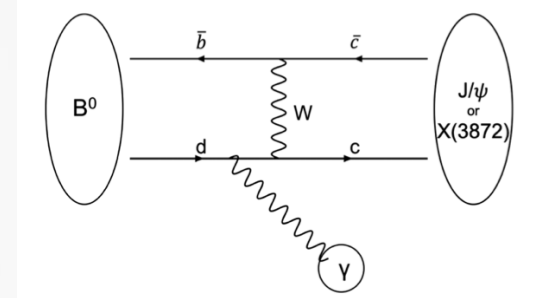


- $M(X(3872)) = (3.8723 \pm 0.0012) \text{ GeV}/c^2$
- With 0.11 ± 0.10 background events, **the number of signal events is $N_{\text{sig}} = 2.9_{-2.0}^{+2.2}(\text{stat.}) \pm 0.1(\text{syst.})$ with a significance of 3.2σ** (Feldman-Cousins method applied [Phys. Rev. D 57, 3873 (1998)]).
- $\tilde{\Gamma}_{\gamma\gamma} \mathcal{B}(X(3872) \rightarrow \pi^+ \pi^- J/\psi) = 5.5_{-3.8}^{+4.1}(\text{stat.}) \pm 0.7(\text{syst.}) \text{ eV}$ using the Q^2 dependence expected from a $c\bar{c}$ meson model.

Search for $B^0 \rightarrow X(3872)\gamma$

[PRD 100, 012002 (2019)]

- In the SM, the decay $B^0 \rightarrow c\bar{c}\gamma$ proceeds dominantly through an exchange of a W boson and the radiation of a photon from the d quark of the B meson.
- Currently, the upper limit for $B^0 \rightarrow J/\psi\gamma$ is 1.5×10^{-6} at 90% confidence level.
- Considering $X(3872)$ may be not a pure $c\bar{c}$ state, the branching fraction of $B^0 \rightarrow X(3872)\gamma$ is larger?



To suppress generic BB spherical events and the jet-like $q\bar{q}$ continuum events, we do

(1) **multivariate analysis** based on the neural network package named NEUROBAYES [Nucl. Instrum. Methods Phys. Res., Sect. A 559, 190

(2006)] to distinguish the signal and background with 33 input variables;

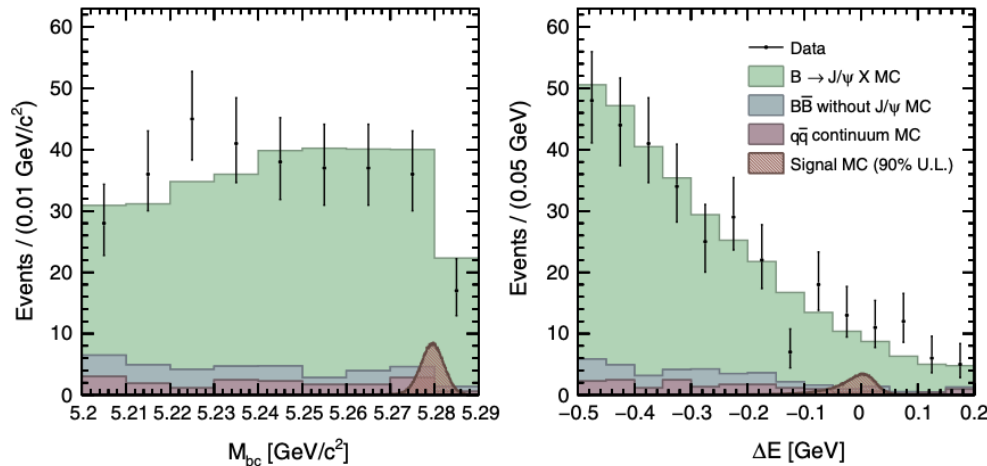
(2) **optimize a figure of merit (FOM).**

$X(3872)$ decays to $J/\psi\pi^+\pi^-$ entirely via $J/\psi\rho$.

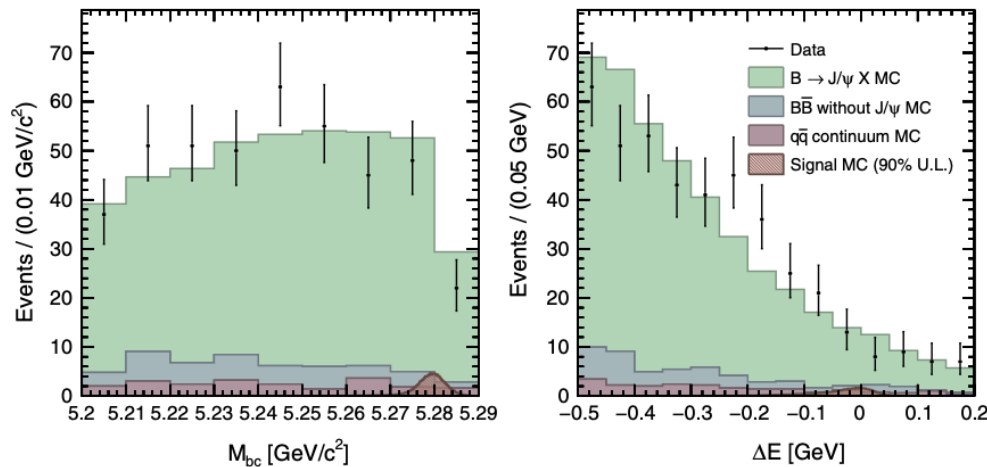
$$\text{FOM} = \frac{\text{efficiency}}{0.5n + \sqrt{N_{\text{bkg}}}}$$

Total luminosity:

$711 \text{ fb}^{-1}; 772 \times 10^6 \text{ B}\bar{\text{B}}$ pairs



(a) Dimuon channel.



(b) Dielectron channel.

$$\Delta E = E_{\text{recon}}^* - E_{\text{beam}}^*$$

$$M_{bc} = \sqrt{E_{\text{beam}}^2 - \left(\sum_i p_i\right)^2}$$

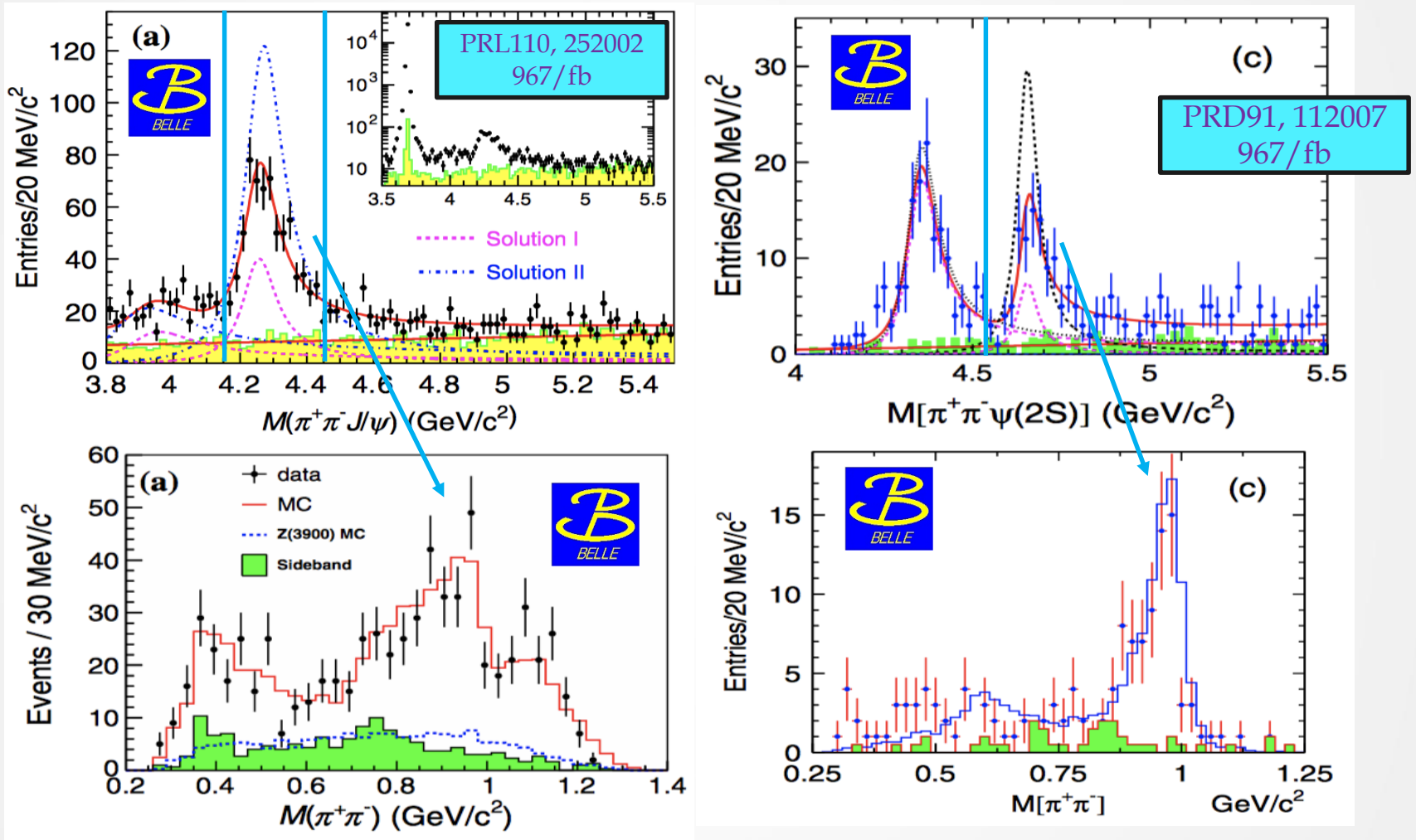
We count the numbers of signal and background events in regions of $M_{bc} > 5.27 \text{ GeV}/c^2$ and $-0.15 < \Delta E < 0.1 \text{ GeV}$.

The upper limit on $\mathcal{B}(B^0 \rightarrow X(3872)\gamma) \times \mathcal{B}(X(3872) \rightarrow J/\psi \pi^+ \pi^-)$ is obtained with the Feldman-Cousins counting method.

Channel	Dimuon	Dielectron	Total
N_{sig}	9	9	18
N_{bkg}	9.3	12.1	21.4
90% U.L.	9.2×10^{-7}	6.8×10^{-7}	5.1×10^{-7}

$$Y(4626): e^+e^- \rightarrow D_s^+ D_{s1}(2536)^- / D_s^+ D_{s2}^*(2573)^- + c.c.$$

Motivation: $Y(4260)$ and $Y(4660)$ with $c\bar{c}s\bar{s}$ component



- $Y(4260) \rightarrow f_0(980)(\rightarrow \pi^+ \pi^-)J/\psi$, $Y(4660) \rightarrow f_0(980)(\rightarrow \pi^+ \pi^-)\psi(2S)$
 $f_0(980)$ has a $s\bar{s}$ component, and ψ has a $c\bar{c}$ component.
- **It is natural to search for such Y states with a quark component of $(c\bar{s})(\bar{c}s)$, e.g., $D_s D_{s1}(2536)$ and $D_s D_{s2}^*(2573)$.**

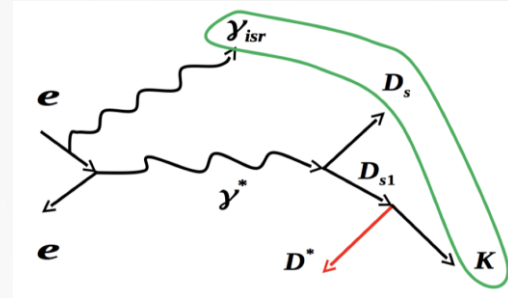
Analysis method

Taking $e^+e^- \rightarrow D_S^+ D_{S1}(2536)^-$ as an example

$$e^+e^- \rightarrow \gamma_{ISR} D_S^+ D_{S1}(2536)^- (\rightarrow \bar{D}^{*0} K^- / D^{*-} K_S^0)$$

We require full reconstruction of the γ_{ISR} , D_S^+ , and K^- / K_S^0 .

- $D_S^+ \rightarrow \phi\pi^+, \bar{K}^{*0}K^+, K_S^0K^+, K^+K^-\pi^+\pi^0, K_S^0\pi^0K^+, K^{*+}K_S^0, \eta\pi^+$, and $\eta'\pi^+$
- For the signals, the spectrum of the mass recoiling against the $D_S^+ K^- \gamma_{ISR}$ system should be accumulated at the \bar{D}^{*0} / D^{*-} nominal mass.



$$M_{\text{rec}}(\gamma_{ISR} D_S^+ K^- / K_S^0) = \sqrt{(E_{\text{c.m.}}^* - E_{\gamma_{ISR} D_S^+ K^- / K_S^0}^*)^2 - (p_{\gamma_{ISR} D_S^+ K^- / K_S^0}^*)^2}$$

- To improve the $M(D_S^+ D_{S1}(2536)^-)$ resolution, $M_{\text{rec}}(\gamma_{ISR} D_S^+ K^- / K_S^0)$ is constrained to be the nominal mass of the \bar{D}^{*0} / D^{*-} .

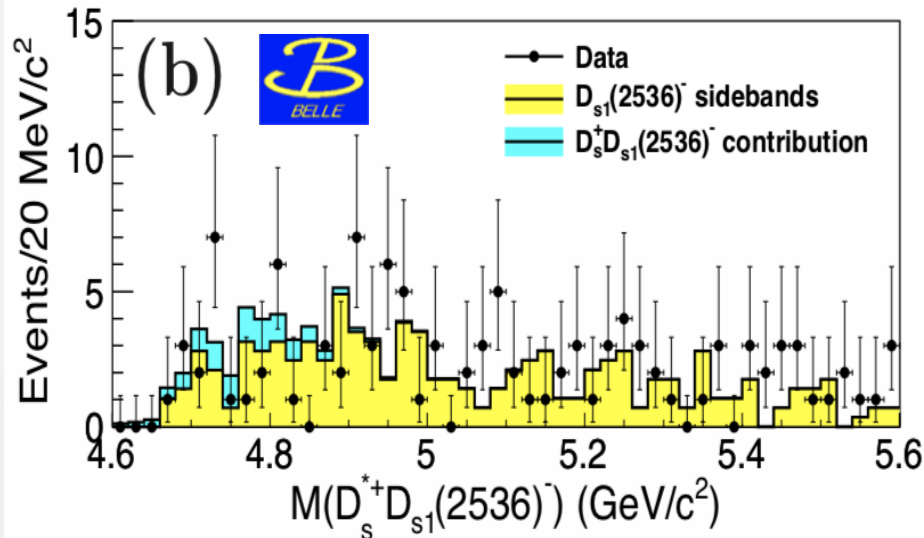
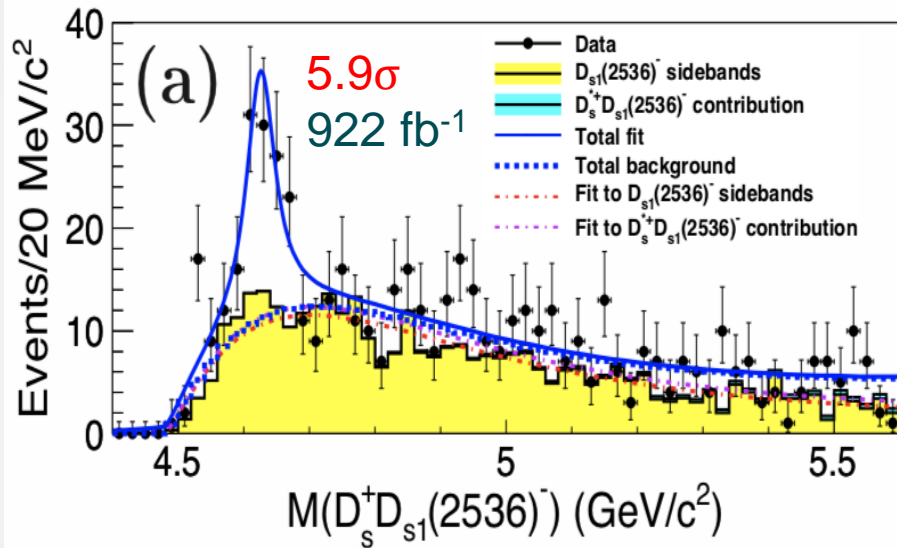
$Y(4626): e^+e^- \rightarrow D_s^+ D_{s1}^-(2536)^-$

Belle, PRD100, 111103(R) (2019)

After applying the \bar{D}^{*0}/D^{*-} mass constraint

An unbinned simultaneous likelihood fit:

- Signal: a BW convolved with a Gaussian function, then multiplied by an efficiency function
- $D_{s1}(2536)^-$ mass sidebands: a threshold function
- $e^+e^- \rightarrow D_s^{*+} D_{s1}^-(2536)^-$ background contribution: a threshold function
- A non-resonant contribution: a two-body phase space form



$$M = (4625.9^{+6.2}_{-6.0}(\text{stat.}) \pm 0.4(\text{syst.})) \text{ MeV}/c^2$$

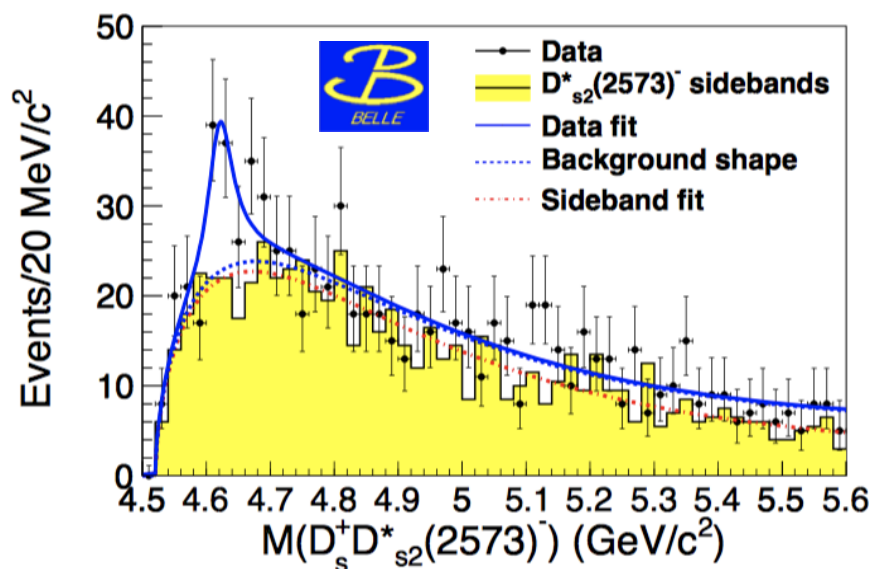
$$\Gamma = (49.8^{+13.9}_{-11.5}(\text{stat.}) \pm 4.0(\text{syst.})) \text{ MeV}$$

$$\Gamma_{ee} \times \mathcal{B}(Y \rightarrow D_s^+ D_{s1}^-(2536)^-) \times \mathcal{B}(D_{s1}^-(2536)^- \rightarrow \bar{D}^{*0} K^-) = (14.3^{+2.8}_{-2.6}(\text{stat.}) \pm 1.5(\text{syst.})) \text{ eV}$$

One possible background is from $e^+e^- \rightarrow D_s^{*+} (\rightarrow D_s^+ \gamma) D_{s1}^-(2536)^-$.
 No obvious structure is observed in the $e^+e^- \rightarrow D_s^{*+} (\rightarrow D_s^+ \gamma) D_{s1}^-(2536)^-$.

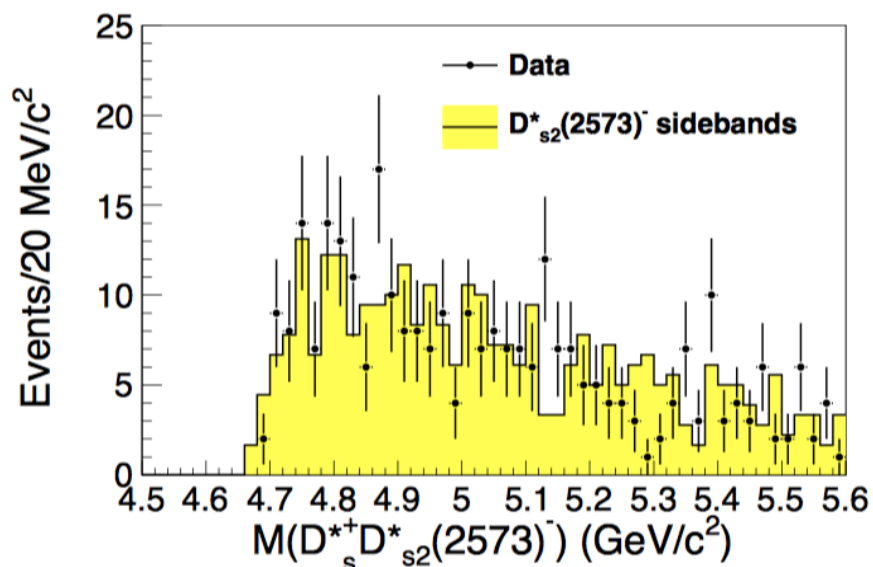
$Y(4626): e^+e^- \rightarrow D_s^+ D_{s2}^*(2573)^-$

To improve the $M_{\text{rec}}(Y_{\text{ISR}})$ resolution, $M_{\text{rec}}(Y_{\text{ISR}} D_s^+ K^-)$ is constrained to the nominal mass of the \bar{D}^0 .



Belle, PRD101, 091101(R) (2020)

- An unbinned simultaneous likelihood fit:
- Signal: a BW convolved with a Gaussian function, then multiplied by an efficiency function
 - $D_{s2}^*(2573)^-$ mass sidebands: a threshold function
 - A non-resonant contribution: a two-body phase space form



$$M = (4619.8_{-8.0}^{+8.9}(\text{stat.}) \pm 2.3(\text{syst.})) \text{ MeV}/c^2$$

$$\Gamma = (47.0_{-14.8}^{+31.3}(\text{stat.}) \pm 4.6(\text{syst.})) \text{ MeV}$$

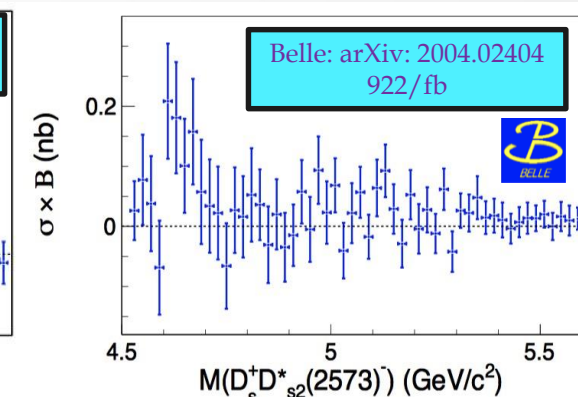
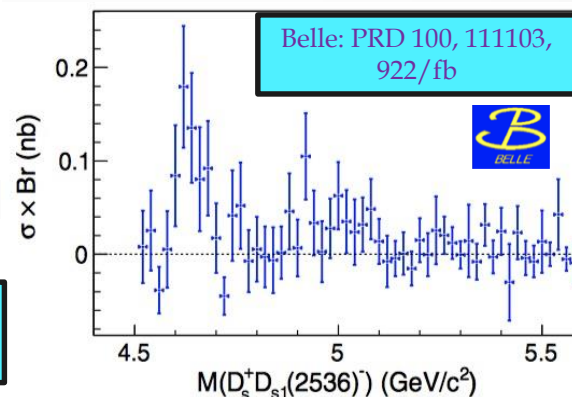
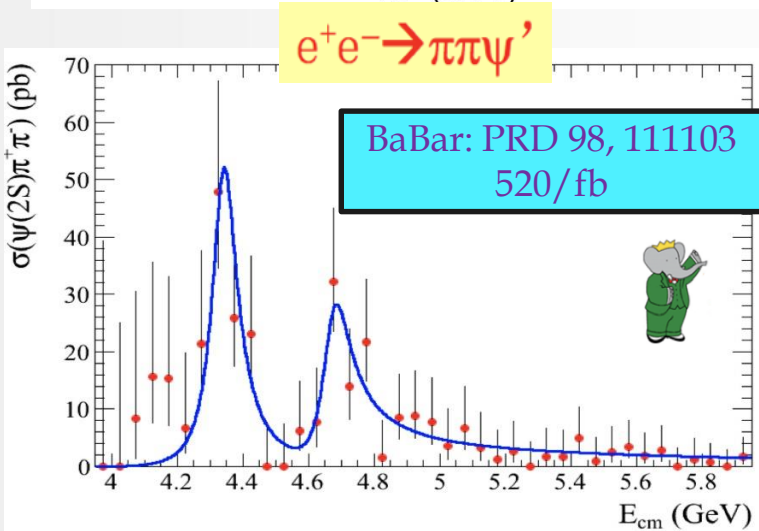
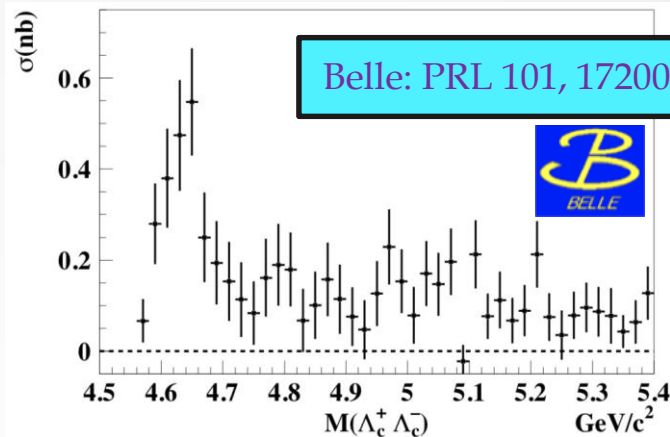
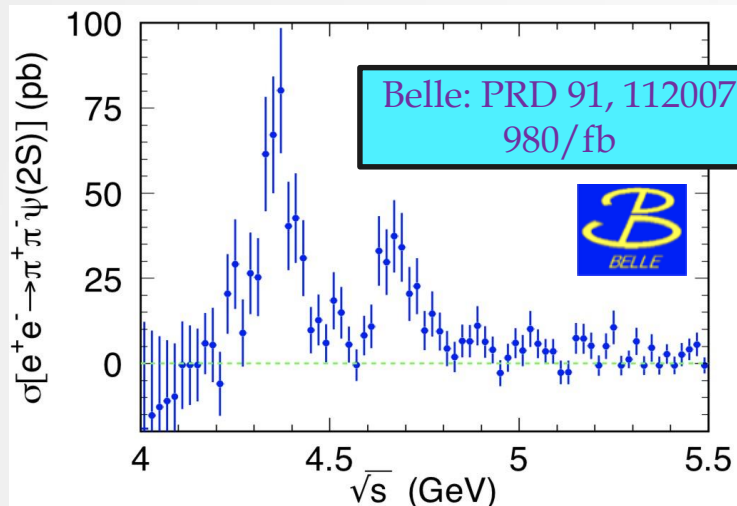
$$\Gamma_{ee} \times \mathcal{B}(Y \rightarrow D_s^+ D_{s2}^*(2573)^-) \times \mathcal{B}(D_{s2}^*(2573)^- \rightarrow \bar{D}^0 K^-) = (14.7_{-4.5}^{+5.9}(\text{stat.}) \pm 3.6(\text{syst.})) \text{ eV}$$

Reminder: considering BESIII has had data in this energy region, we need BESIII to cross check. The current error at Belle is large.

Interpretations of $Y(4626)$

- A **tetraquark state** in a chiral constituent quark model with a scaling method [Y.Tan and J. L. Ping, PRD101, 054010 (2020)].
- A **P-wave tetraquark state** $[cs][\bar{c}\bar{s}]$ with 1^{--} in the multiquark color flux-tube model [C. R. Deng, H. Cheng and J.L. Ping, PRD 101, 054039 (2020)].
- A hidden-strange **molecular state from $\Lambda_c^+ \Lambda_c^-$ interaction** [J. T. Zhu, Y. Liu, D. Y. Chen, L. Y. Jiang, and J. He, arXiv:1911.03706 (2020)].
- A **molecular state from interaction $D_s^* \bar{D}_{s1}(2536) - D_s \bar{D}_{s1}(2536)$** [J. He, J. T. Zhu, and D. Y. Chen, EPJC 80, 246 (2020)].
- A tetraquark and etc instead of $D_s^* \bar{D}_{s1}(2536)$ molecular within the Bethe-Salpeter framework [H.W.Ke, X.H.Liu, and X.Q.Li, arXiv:2004.03167 (2020)].
- A higher charmonium [J.Z.Wang, R.Q.Qian, X. Liu, and T. Matsuki, PRD 101, 034001 (2020)].
- A hidden-charm exotic mesons in the diquark model [Z. G. He, B. A. Kniehl, and X.P.Wang, PRD 101, 074032 (2020); J.F.Giron, R.F.Lebed arxiv:2005.07100].

$Y(4630) = Y(4660)?$



Experiment	Mass (MeV)	Width (MeV)
Belle, $\Lambda_c^+ \Lambda_c^-$	$4634^{+8}_{-7} {}^{+5}_{-8}$	$92^{+40}_{-24} {}^{+10}_{-21}$
Belle, $\pi^+ \pi^- \psi(2S)$	$4652 \pm 10 \pm 8$	$68 \pm 11 \pm 1$
BaBar, $\pi^+ \pi^- \psi(2S)$	$4669 \pm 21 \pm 3$	$104 \pm 48 \pm 10$
Belle, $D_s^+ D_{s1}(2536)^-$	$4626^{+7}_{-7} \pm 1$	$49.8^{+14}_{-12} \pm 4$
Belle, $D_s^+ D_{s2}^*(2573)^-$	$4620^{+9}_{-8} \pm 3$	$47.0^{+32}_{-15} \pm 5$

- These states may be the same
- Need improved precision

Search for $R^{++} \rightarrow D^+ D_s^{*+}$

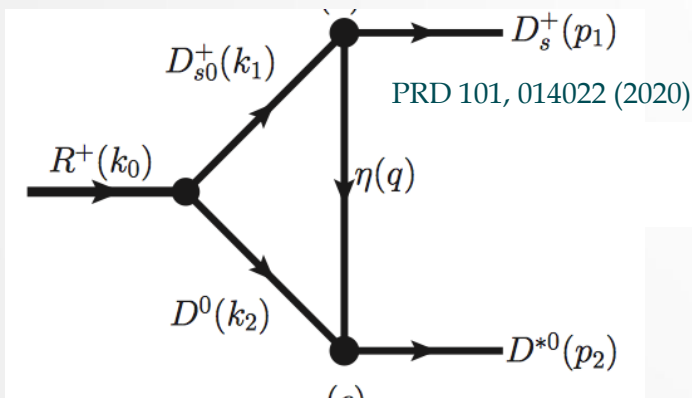
Phys. Rev. D 102, 112001 (2020)

- A doubly-charged and doubly-charmed molecule R^{++} decays to $D^+ D_s^{*+}$ with modest rates according to Refs. [PRD 99, 076017 (2019), PRD 101, 014022 (2020)].
- The mass of R^{++} is predicted to be in the range of 4.13 to 4.17 GeV/c^2 ; the width is (2.30–2.49) MeV.
- A state decaying to $D^+ D_s^{*+}$ is also a good candidate for a doubly-charged tetraquark according to Ref. [PRL 119, 202001 (2017)].

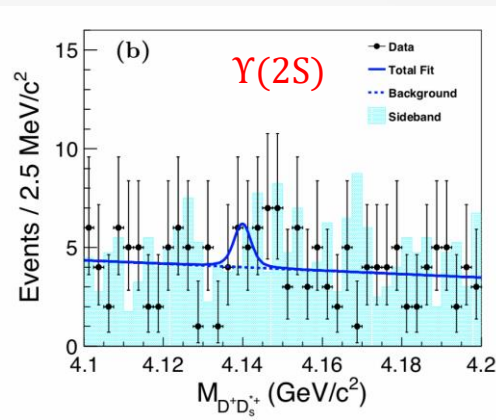
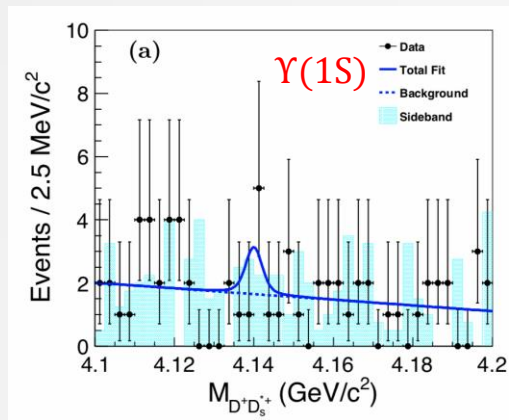
Data samples:

\sqrt{s} (GeV)	Luminosity (fb^{-1})	Events
9.46 [Y(1S)]	5.74 ± 0.09	(102±3) million
10.023 [Y(2S)]	24.91 ± 0.35	(158±4) million
10.52	89.5 ± 1.3	-
10.58 [Y(4S)]	711 ± 10	-
10.867 [Y(5S)]	121.4 ± 1.7	-

The Punzi parameter $S/(3/2 + \sqrt{B})$ [arXiv:physics/0308063] is applied to optimize the mass windows of intermediate states.

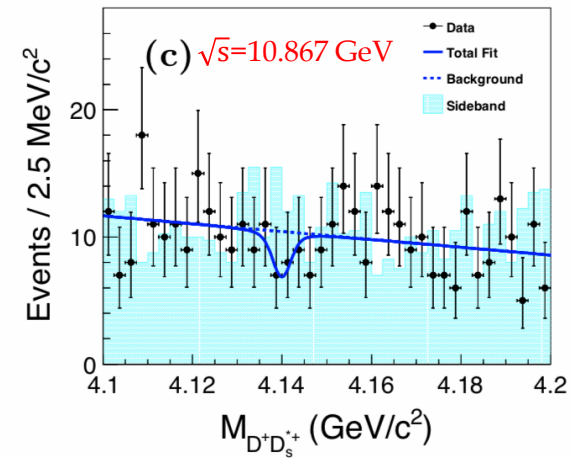
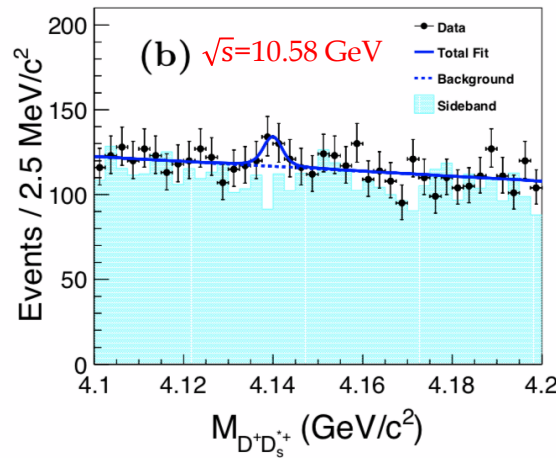
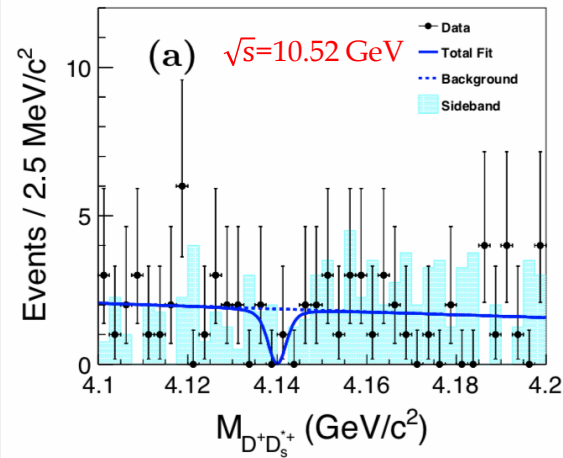


- $D^+ \rightarrow K^- \pi^+ \pi^- / K_s^0 (\rightarrow \pi^+ \pi^-) \pi^+$
- $D_s^{*-} \rightarrow D_s^- \gamma$
- $D_s^- \rightarrow \phi \pi^- / \bar{K}^{*0} K^+$



The fitted results with the R^{++} mass fixed at $4.14 \text{ GeV}/c^2$ and width fixed at 2 MeV .

No R^{++} signals are observed.



90% C. L. Upper limits [$M(R^{++})$ varying from 4.13 to $4.17 \text{ GeV}/c^2$, $\Gamma(R^{++})$ varying from 0 to 5 MeV]

$$\mathcal{B}(Y(1S) \rightarrow R^{++} + \text{anything})\mathcal{B}(R^{++} \rightarrow D^+D_s^{*+}) = (11.8 - 54.5) \times 10^{-5}$$

$$\mathcal{B}(Y(2S) \rightarrow R^{++} + \text{anything})\mathcal{B}(R^{++} \rightarrow D^+D_s^{*+}) = (16.3 - 68.6) \times 10^{-5}$$

$$\sigma(e^+e^- \rightarrow R^{++} + \text{anything})\mathcal{B}(R^{++} \rightarrow D^+D_s^{*+}) = (202.8 - 880.4) \text{ fb at } \sqrt{s} = 10.52 \text{ GeV}$$

$$\sigma(e^+e^- \rightarrow R^{++} + \text{anything})\mathcal{B}(R^{++} \rightarrow D^+D_s^{*+}) = (218.9 - 1054.0) \text{ fb at } \sqrt{s} = 10.58 \text{ GeV}$$

$$\sigma(e^+e^- \rightarrow R^{++} + \text{anything})\mathcal{B}(R^{++} \rightarrow D^+D_s^{*+}) = (346.6 - 1517.6) \text{ fb at } \sqrt{s} = 10.867 \text{ GeV}$$



$Z_{cs}(4000)^+$ at LHCb

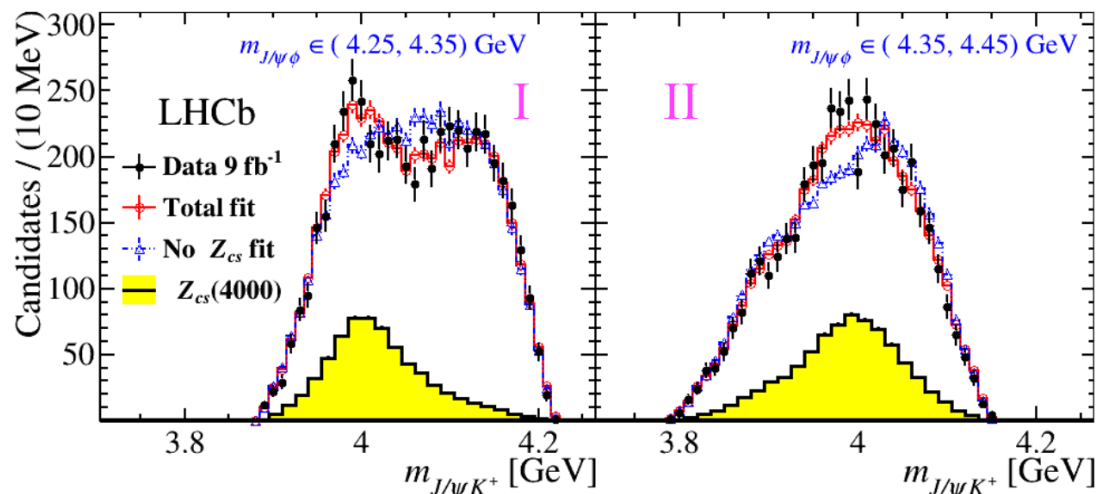
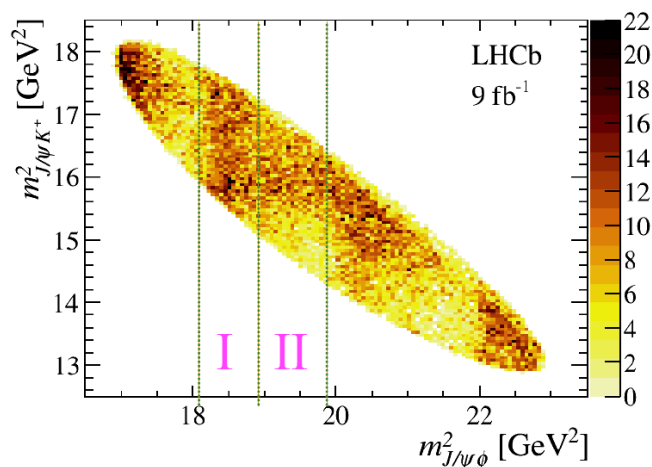
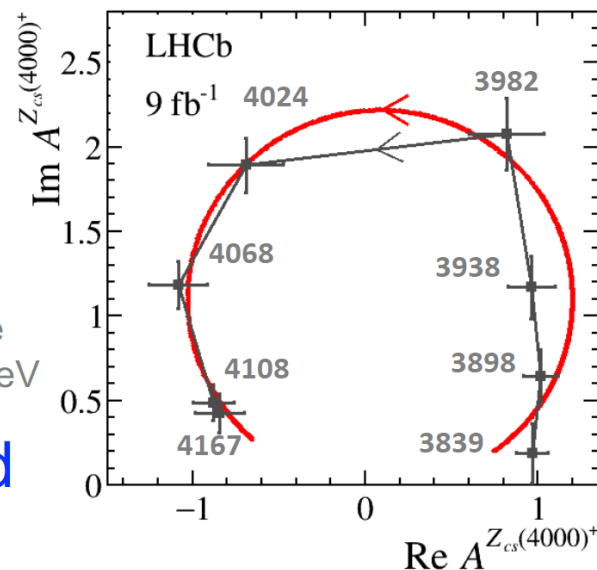
[arXiv:2103.01803]

$B^+ \rightarrow K^+ \Phi J/\psi$: Run2 (4*Run1)+Run1 data

- Argand diagram gives further evidence of resonant character
 - Magnitude and phase evolved in the counter-clockwise direction

- $Z_{cs}(4000)^+$ can be clearly viewed in the two slices of $m_{J/\psi\phi}$

numbers are $m_{J/\psi K^+}$ in MeV



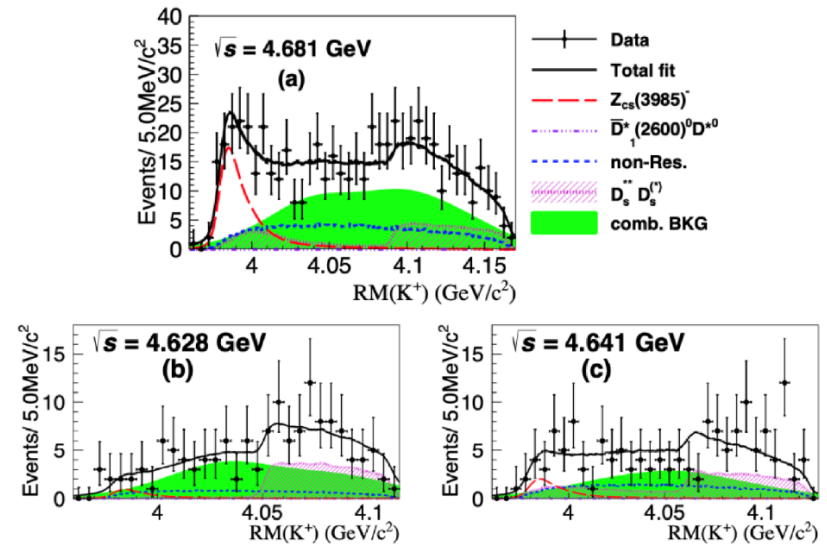


Comparison with BESIII

- BESIII experiment recently reported 5.3σ observation of a very narrow Z_{CS}^- in $D_S^- D^* + DD_S^{*-}$ mass distributions
- Their masses are close, but $Z_{CS}(4000)^+$ is $\sim 10\times$ broader
- Tests are applied:
 - Fix $Z_{CS}(4000)^+$ to BESIII's result; $2\ln L$ is worse by 160
 - Adding on top of the default model almost doesn't improve the fit likelihood
- No evidence that $Z_{CS}(4000)^+$ is the same as $Z_{CS}(3985)^-$ seen by BESIII

From Liming Zhang

PRL126, 102001 (2021)

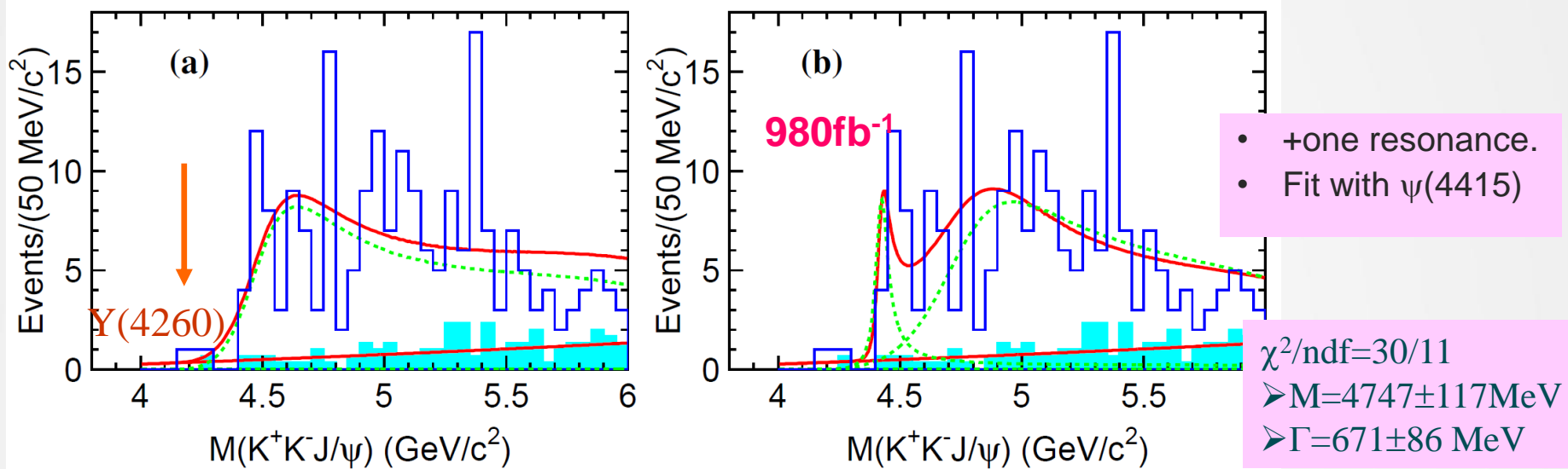


$$m_{\text{pole}}(Z_{CS}(3985)^-) = (3982.5_{-2.6}^{+1.8} \pm 2.1) \text{ MeV}/c^2,$$

$$\Gamma_{\text{pole}}(Z_{CS}(3985)^-) = (12.8_{-4.4}^{+5.3} \pm 3.0) \text{ MeV}.$$

$e^+e^- \rightarrow K^+K^-J/\psi$

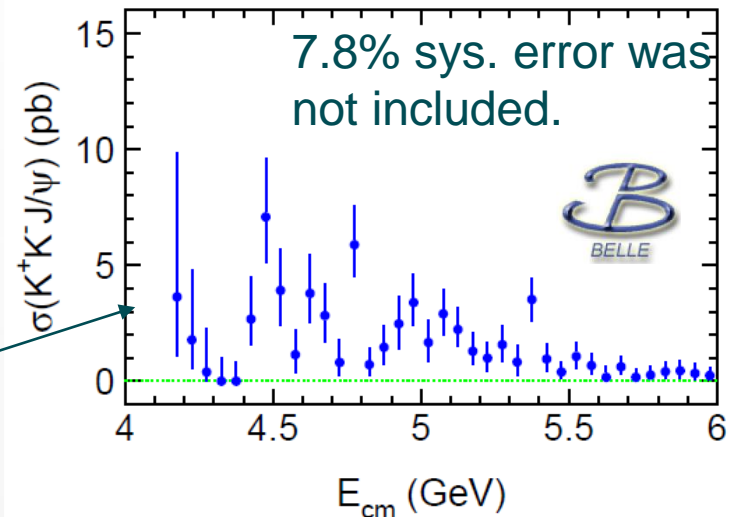
PRD 89,072015(2014)



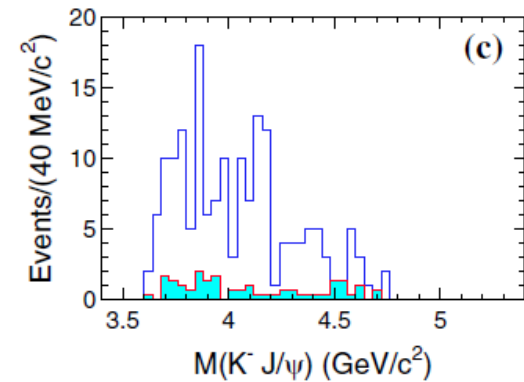
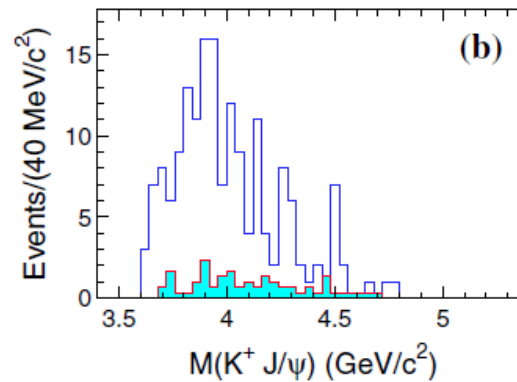
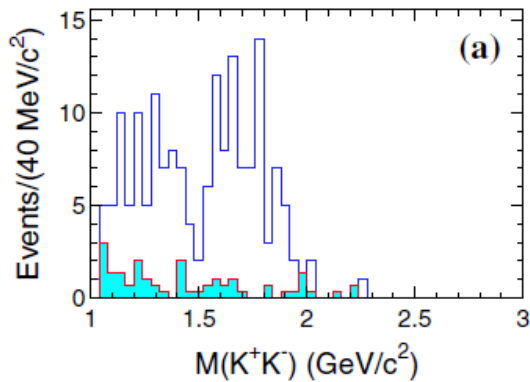
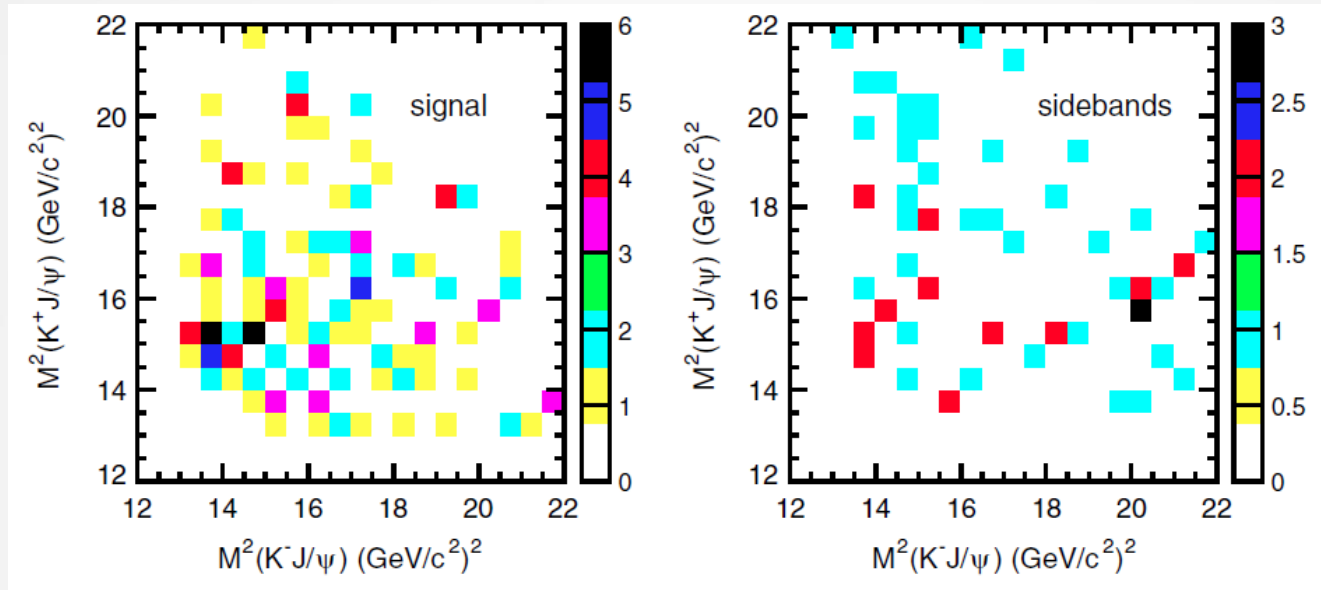
Shaded hist.: J/ψ mass sidebands

4-6 GeV: 213 events
 35 bkg, 178 ± 16 signal

$$\sigma_i = \frac{n_i^{\text{obs}} - f \times n_i^{\text{bkg}}}{\mathcal{L}_i \cdot \epsilon_i \cdot \mathcal{B}(J/\psi \rightarrow l^+l^-)}$$



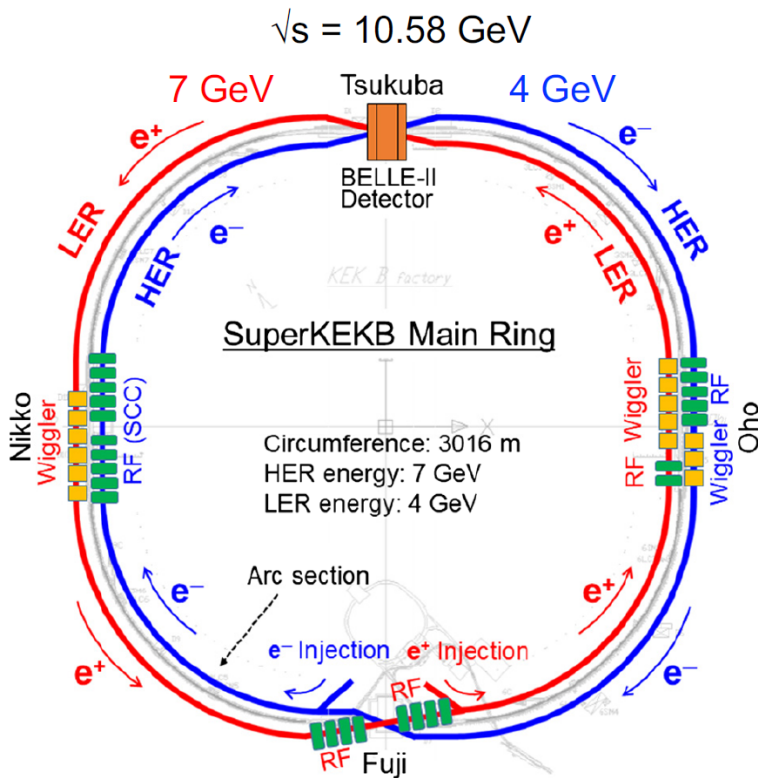
Search for $Z_{cs} \rightarrow KJ/\psi$ states in $e^+e^- \rightarrow K^+K^-J/\psi$



No evident structure in K^+J/ψ mass distribution under current statistics

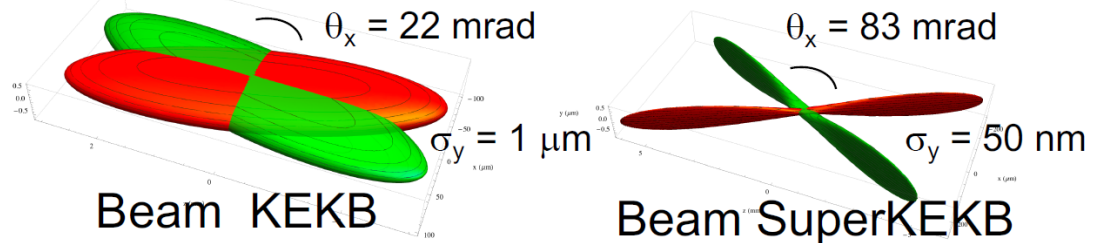
SuperKEKB Collider

SuperKEKB is a new e^+e^- collider located at KEK (Tsukuba, Japan), it operates in the **intensity frontier** region with a target instantaneous luminosity of $6 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$ which is **30 times** larger than that of the previous KEKB collider.



	Instantaneous luminosity ($\text{cm}^{-2} \text{ s}^{-1}$)	Integrated recorded luminosity (ab^{-1})	
Babar PEP-II	$1 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$	0.55	0.43 Y(4S)
Belle KEKB	$2.11 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$	1	0.71 (Y4S)
Belle II SuperKEKB	$6 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$	50	

$\mathcal{L} \times 30$: x1.5 current increase
x 20 β_y^* vertical beta function decrease



Current integrated luminosity

We kept SuperKEKB and Belle II running in 2020/2021 during the COVID-19 crisis, with extra effort from the local crew and the help of remote shifters

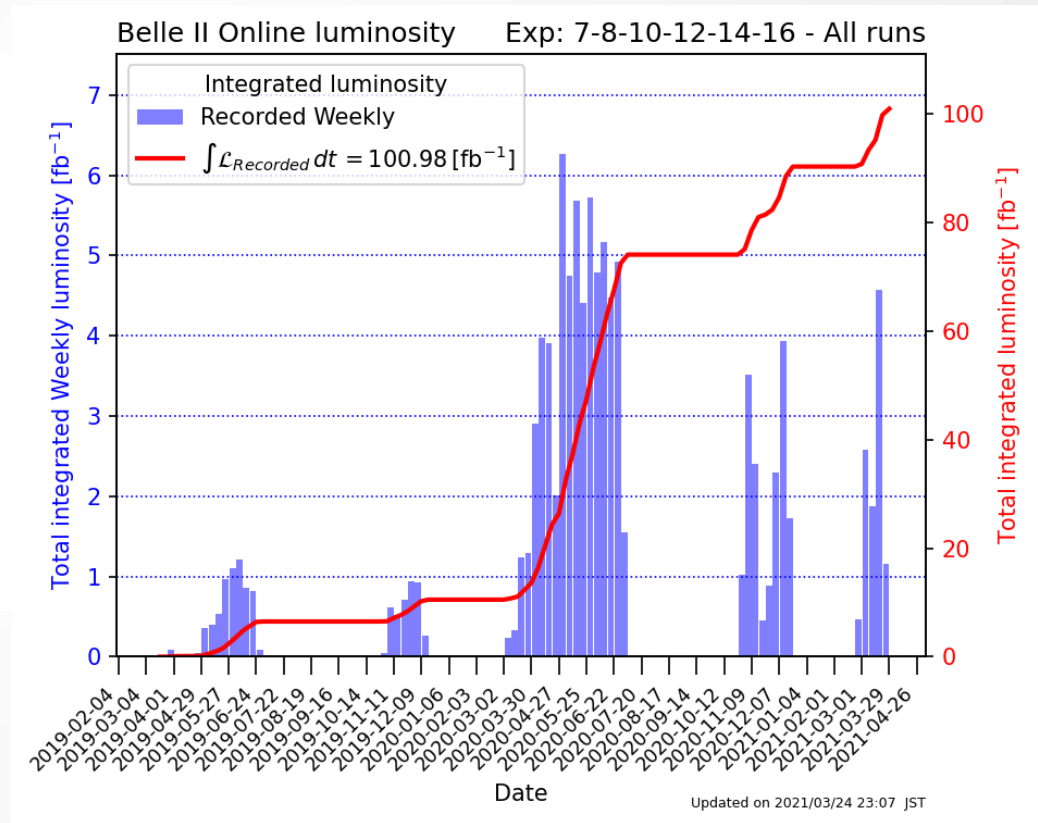
Luminosity world record

$2.11 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$
(KEK June 2009)

$2.14 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$
(LHC May 2018)

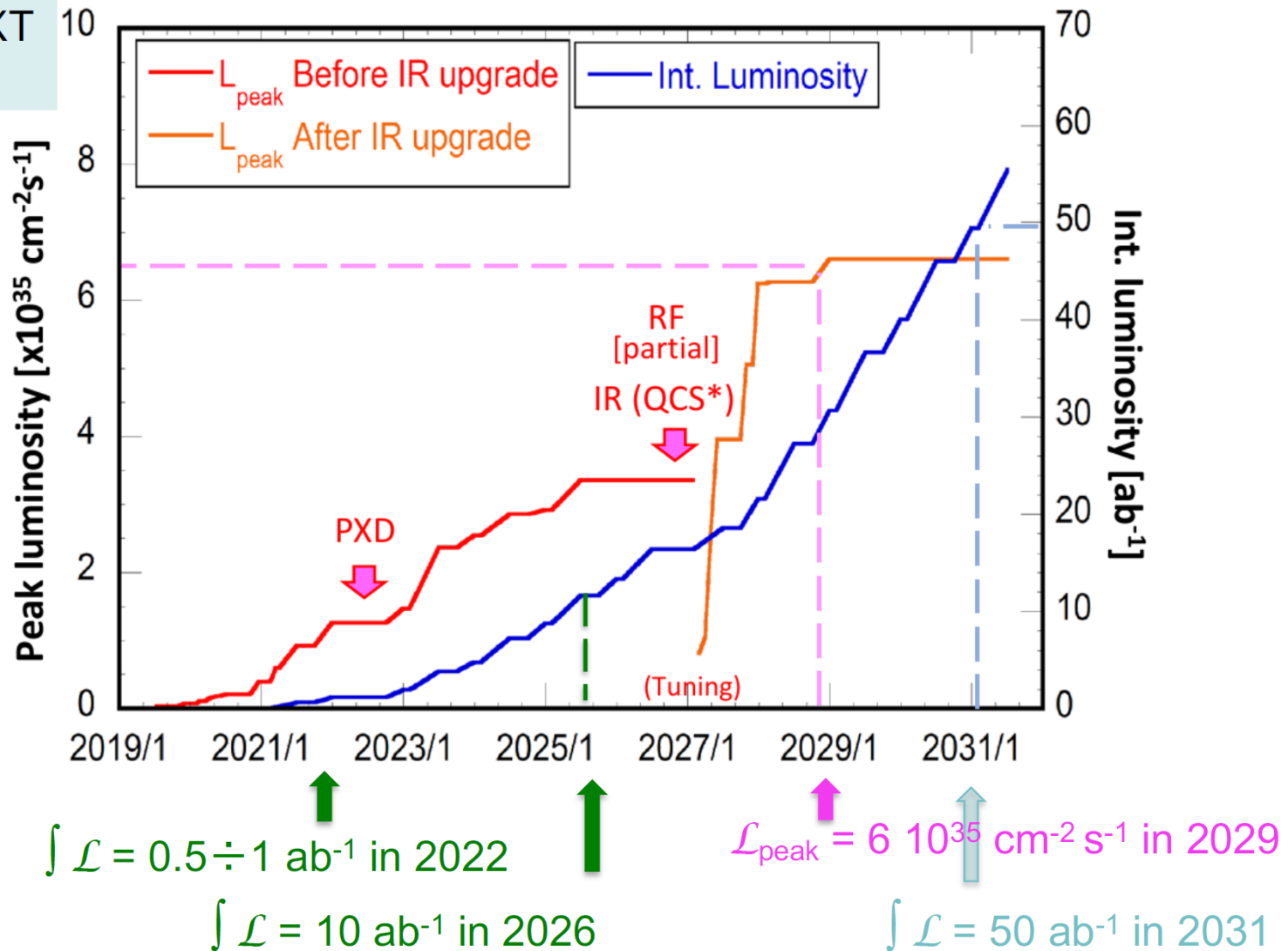
$2.4 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$
(SuperKEKB June 2020)

$$\int \mathcal{L}_{\text{Recorded}} dt = 100.98 [\text{fb}^{-1}]$$



Luminosity Plan

Submitted to MEXT
roadmap 2020



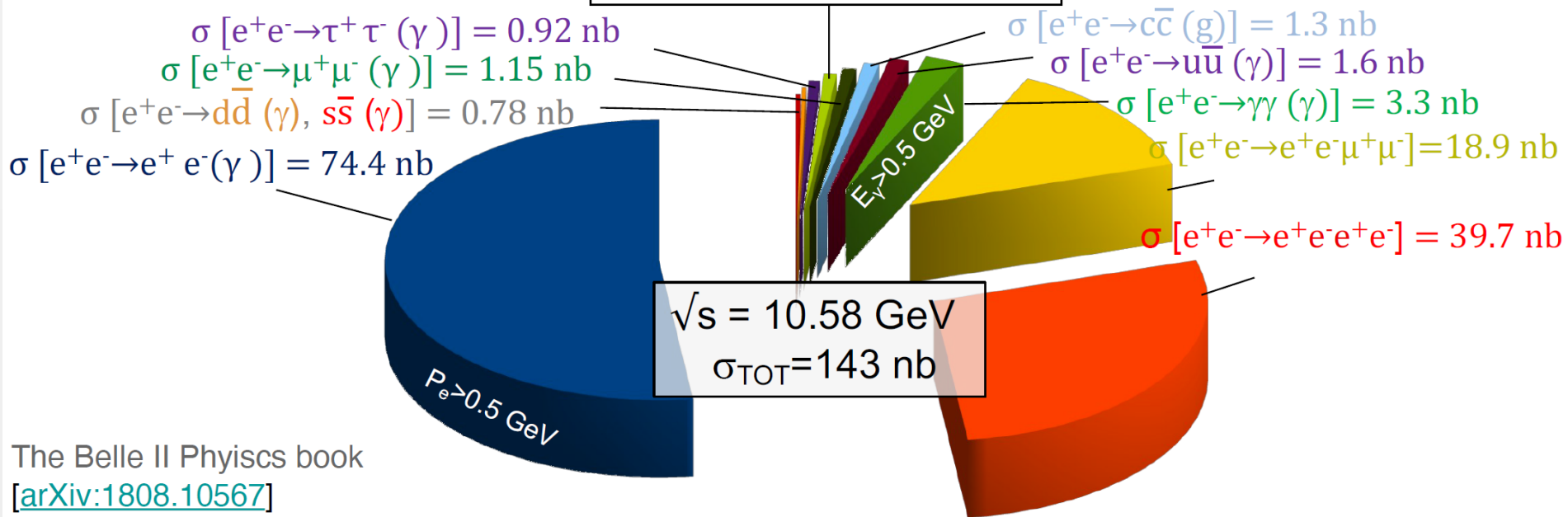
Belle II energy points

Energy scan

Y(1s)	Y(2s)	Y(3s)	Y(4s)	Y(5s)	Y(6s)
9.46 GeV	10.02 GeV	10.35 GeV	10.58 GeV	10.86 GeV	11.02 GeV

B^+B^- (51.4 ± 0.6)% $B^0\bar{B}^0$ (48.6 ± 0.6)%

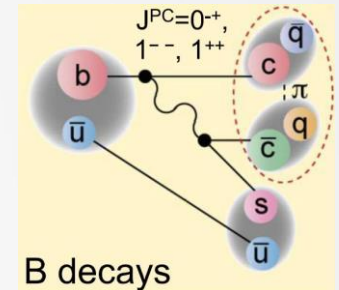
$$\sigma [e^+e^- \rightarrow Y(4S)] = 1.11 \text{ nb}$$



Potential XYZ results in charm sector at Belle II

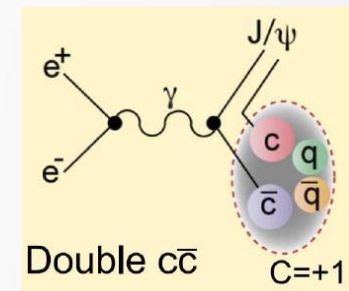
B decays: Competition from LHCb, advantages for modes with neutrals

- Confirm Z_c states and search for neutral partners
- Absolute branching fractions $B \rightarrow X(3872, 3915) K$
- Confirmation of $X(3872)$ width measurement with $D^0 \bar{D}^0 \pi^0$ mode
- **Absolute branching fractions are unique for Belle II**



ISR processes:

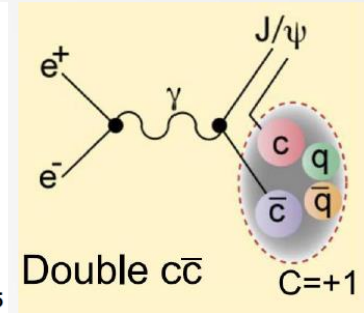
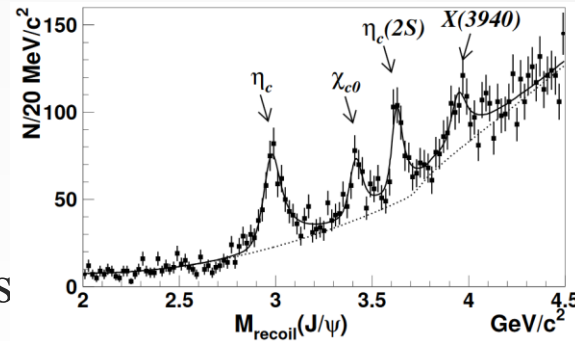
- Continuous mass range to investigate fine structures
- Higher mass region
- Confirm Z_c states and search for neutral partners
- **Higher mass region (>4.7 GeV) is unique for Belle II**



Potential XYZ results in charm sector at Belle II

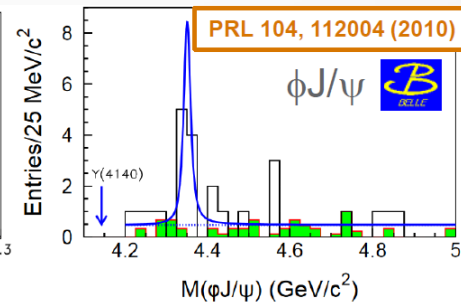
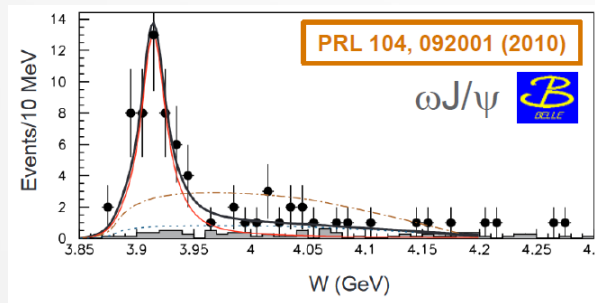
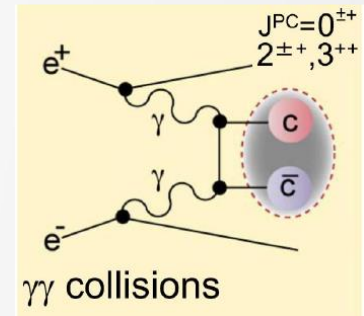
Double charmonium processes:

- $e^+e^- \rightarrow (c\bar{c})_{J=1} (c\bar{c})_{J=0}$ production rule
- Rediscovery of X(3940, 4160)
- Expand to other $c\bar{c}$, search for new states



Two-photon processes:

- Determine J^P values for some confirmed states, like X(3915)
- Higher mass region
- Confirm some states with evidence, like X(4350)
- Check more modes, like $D^{(*)}\bar{D}^{(*)}(n)\pi$



Rediscover the X(3872)

- Reconstruct final states:

- $B^\pm \rightarrow \pi^+\pi^- J/\psi(\ell^+\ell^-) K^\pm$
- $B^0 \rightarrow \pi^+\pi^- J/\psi(\ell^+\ell^-) K_S$

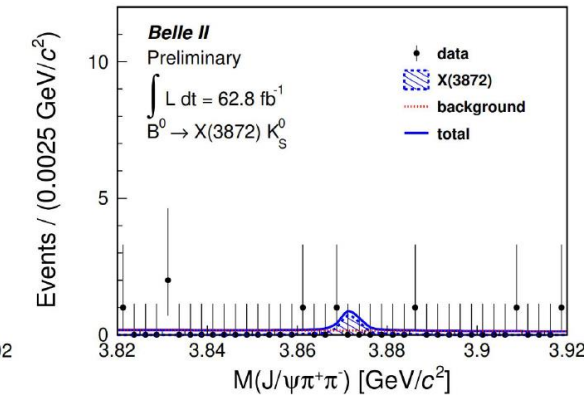
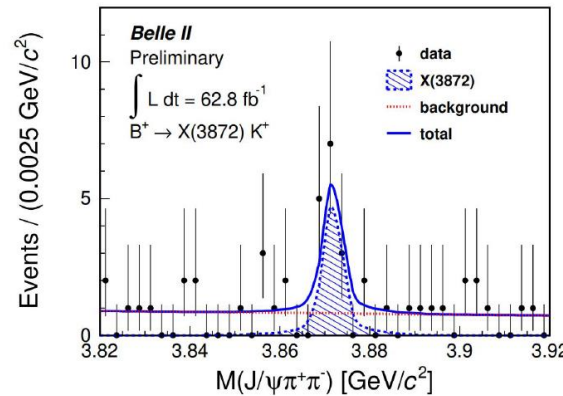
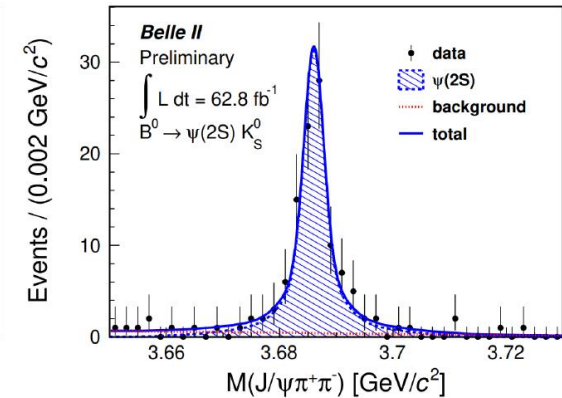
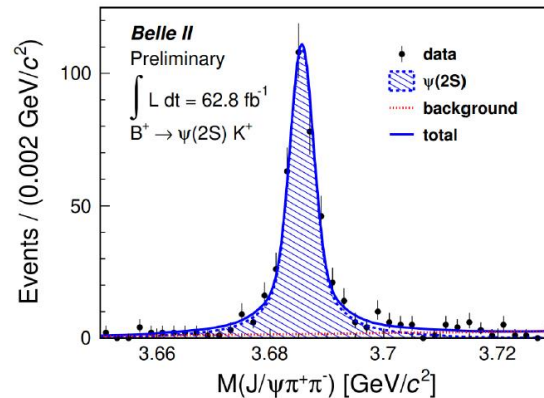
- “Standard” selection criteria

- Particle identification
- Continuum: nTracks, R_2
- Kinematics: $M_{\pi^+\pi^-}$, M_{BC} , $|\Delta E|$

- Observe $B \rightarrow \psi(2S) K$

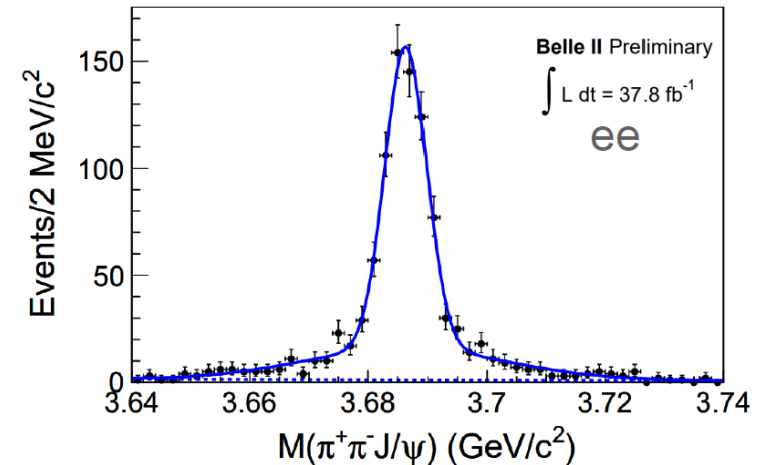
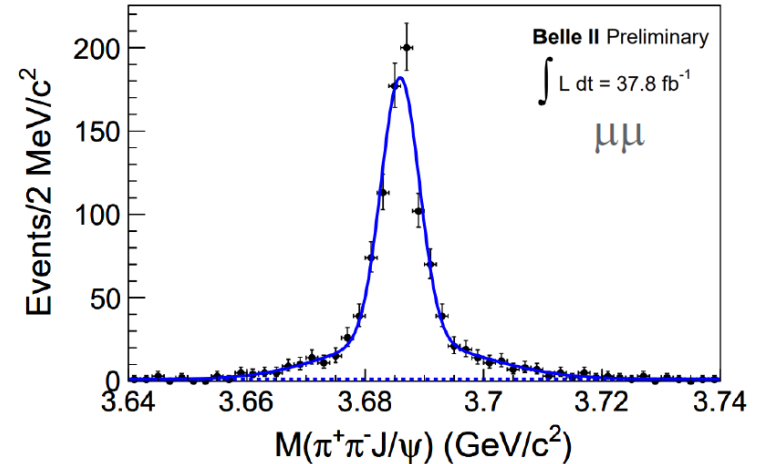
- First X(3872) at Belle II

- 14.4 ± 4.6 events (4.6σ)



ISR preliminary studies at Belle II

- $e^+e^-\gamma_{\text{ISR}} \rightarrow \pi^+\pi^-J/\psi(\ell^+\ell^-)$ final states
 - Nominal PID requirements
 - $|M(J/\psi)-M(\text{PDG})| < 75 \text{ MeV}$
 - ISR photon not required (high efficiency)
 - $|MM^2(\pi^+\pi^-J/\psi)| < 2 \text{ GeV}^2$
- Clear observation of ISR $\psi(2S)$ signals
- Next step: “Y(4260)” rediscovery
 - Expect ~ 60 total events per 100 fb^{-1}



Summary

- We are still producing interesting XYZ results using Belle data
- The expected Belle II data sample of 50 ab^{-1} will provide a lot of new opportunities for physics analyses
- Some of them, for example, double charmonium production, charmonium in two-photon processes, are unique for Belle II.
- Several quarkonium states and exclusive B decays to charmonium and other particles were “rediscovered” using the currently available data.

Thanks a lot!



Thanks for your attention

沈成平

shencp@buaa.edu.cn

