

The Silicon Vertex Detector of the Belle II Experiment

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Abstract

The Silicon Vertex Detector (SVD) is a part of the vertex detector in the Belle II experiment at the SuperKEKB collider (KEK, Japan). Since the start of data taking in spring 2019, the SVD has been operating stably and reliably with a high signal-to-noise ratio and hit efficiency, achieving good spatial resolution and high track reconstruction efficiency. The hit occupancy, which mostly comes from the beam-related background, is currently about 0.5% in the innermost layer, causing no impact on the SVD performance. In anticipation of the operation at higher luminosity in the next years, two strategies to sustain the tracking performance in future high beam background conditions have been developed and tested on data. One is to reduce the number of signal waveform samples to decrease dead time, data size, and occupancy. The other is to utilize the good hit-time resolution to reject the beam background hits. We also measured the radiation effects on the sensor current, strip noise, and full depletion voltage caused during the first two and a half years of operation. The results show no detrimental effect on the SVD performance.

Keywords: Silicon strip detector, Vertex detector, Tracking detector, Belle II

1. Introduction

The Belle II experiment [1] aims to probe new physics beyond the Standard Model in high-luminosity e^+e^- collisions at the SuperKEKB collider (KEK, Japan) [2]. SuperKEKB consists of the following components: injector LINAC, positron damping ring, and main storage ring with the electron and

positron beamlines. The Belle II detector is located at the interaction point (IP) of the two beamlines. The main collision energy in the center-of-mass system is 10.58 GeV on the $\Upsilon(4S)$ resonance, which enables various physics programs based on the large samples of B mesons, τ leptons, and D mesons. Also, the asymmetric energy of the 7 GeV electron beam and 4 GeV positron beam is adopted for time-dependent CP violation measurements. The target of SuperKEKB is to accumulate an integrated luminosity of 50 ab^{-1} with peak luminosity of about

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16 $6 \times 10^{35} \text{ cm}^{-2}\text{s}^{-1}$. In June 2021, SuperKEKB recorded the 56
 17 world's highest instantaneous luminosity of $3.1 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$. 57
 18 The data accumulated before July 2021 corresponds to an integrated 58
 19 luminosity of 213 fb^{-1} . 59

20 The Vertex Detector (VXD) is the innermost detector in the 60
 21 Belle II detector system. The VXD has six layers: the inner 61
 22 two layers (layers 1 and 2) are the Pixel Detector (PXD), and 62
 23 the outer four layers (layers 3 to 6) are the Silicon Vertex De- 63
 24 tector (SVD). The schematic cross-sectional view of the VXD 64
 25 is shown in Fig. 1. The PXD consists of DEPFET pixel sensors 65
 26 and its innermost radius is 1.4 cm from the IP. A detailed 66
 27 description of the SVD appears in Sec. 2. 67

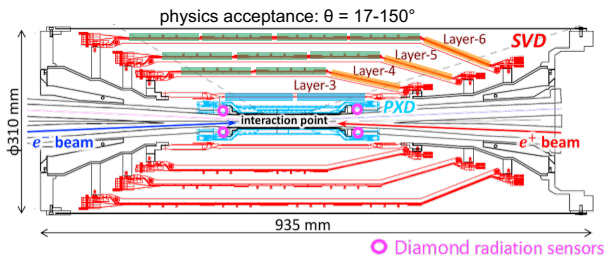


Figure 1: Schematic cross-sectional view of the VXD. The SVD is in red, the PXD in light-blue, and the IP beam pipe diamonds in pink circles. The locations of the three types of DSSDs are indicated by boxes in three colors: blue for small sensors, green for large sensors, and orange for trapezoidal sensors as described in Tab. 1.

28 Besides the VXD, diamond sensors [3] are mounted on the 71
 29 IP beam pipe and the bellows pipes outside of the VXD. The 72
 30 pink circles in Fig. 1 indicate the locations of the diamond sen- 73
 31 sors on the IP beam pipe. They measure the dose rates in these 74
 32 locations. The measured doses are used to estimate the dose in 75
 33 the SVD. They also send beam abort requests to SuperKEKB if 76
 34 the radiation level gets too high to avoid severe damage to the 77
 35 detector. 78

2. Belle II Silicon Vertex Detector

37 The SVD is crucial for extrapolating the tracks to the PXD. 82
 38 This task is essential for measuring the decay vertices with the 83
 39 PXD and pointing at a region-of-interest limiting the PXD read- 84
 40 out volume. Also, the SVD plays a critical role in the decay 85
 41 vertex measurement in the case of long-lived particles like K_S 86
 42 mesons, which decay inside the SVD volume. Other roles of the 87
 43 SVD are the standalone track reconstruction of low-momentum 88
 44 charged particles and their particle identification using ionization 89
 45 energy deposits. 90

46 The SVD [4] consists of four layers of double-sided silicon 91
 47 strip detectors (DSSDs). The material budget of the SVD is 92
 48 about 0.7% of a radiation length per layer. The aluminum read- 93
 49 out strips are AC-coupled to every other n/p-side strips (elec- 94
 50 trodes) on the n-type substrate over the silicon oxide layer. On 95
 51 each DSSD plane, a local coordinate is defined with u and v : 96
 52 u -axis along n-side strips and v -axis perpendicular to u -axis. 97
 53 In other words, p-side strips and n-side strips provide u and v 98
 54 information, respectively. In the cylindrical coordinate, u corre- 99
 55 sponds to $r-\varphi$ information and v corresponds to z information. 100

The SVD consists of three types of sensors: “small” sensors in layer 3, “large” sensors in the barrel region of layers 4, 5, and 6, and “trapezoidal” sensors in the forward region of layers 4, 5, and 6, which is slanted. They are indicated in blue, green, and orange boxes in Fig. 1. The dimensions for these three types of sensors are summarized in Tab. 1. The sensors are manufactured by two companies: the small and large sensors by Hamamatsu and trapezoidal sensors by Micron. The full depletion voltage is 60 V for Hamamatsu sensors, 20 V for Micron sensors; both types of sensors are operated at 100 V. In total, 172 sensors are assembled, corresponding to a total sensor area of 1.2 m^2 and 224,000 readout strips.

	Small	Large	Trapezoidal
No. of u/p-strips	768	768	768
u/p-strip pitch	$50 \mu\text{m}$	$75 \mu\text{m}$	$50\text{--}75 \mu\text{m}$
No. of v/n-strips	768	512	512
v/n-strip pitch	$160 \mu\text{m}$	$240 \mu\text{m}$	$240 \mu\text{m}$
Thickness	$320 \mu\text{m}$	$300 \mu\text{m}$	$300 \mu\text{m}$
Manufacturer	Hamamatsu		Micron

Table 1: Table of dimensions for three types of sensors. Only readout strips are taken into account for number of strips and strip pitch.

The front-end ASIC used in the SVD is APV25 [5], which was originally developed for the CMS silicon tracker. The APV25 is radiation hard for a dose up to 100 Mrad radiation. It has 128 channel inputs and shapers for each channel with a shaping time of about 50 ns. For the SVD, the APV25 is operated in “multi-peak” mode. The mechanism of the data sampling in the multi-peak mode is explained in Fig. 2. The chip samples the height of the signal waveform with the 32 MHz clock and stores each sample’s information in an analog ring buffer. Since the bunch-crossing frequency is eight times faster than the sampling clock, the stored samples are not synchronous to the beam collision, in contrast to CMS, which motivates operation in the multi-peak mode. In the present readout configuration (the six-samples mode), at every reception of the Belle II global Level-1 trigger, the chip reads out six successive samples of the signal waveform stored in the buffers. The six-samples mode offers enough time window ($6/32 \text{ MHz}^{-1} = 187 \text{ ns}$) to accommodate large timing shifts of the trigger. In preparation for operation with higher luminosity, where background occupancy, trigger dead-time, and the data size increase, we developed the three/six-mixed acquisition mode (mixed-mode). The mixed-mode is a new method to read out the signal samples from the APV25, in which the number of the samples changes between three and six in each event, depending on the timing precision of each Level-1 trigger signal in that event. For triggers with good timing precision, three-samples data are read out and the data have half time window and half data size compared to ones of six-samples data, resulting in the reduction of the effect due to higher luminosity. This functionality was already implemented in the running system and confirmed by a few hours of smooth physics data-taking. Before we start to use the mixed-mode, the effect on the performance due to the change of the acquisition mode is to be assessed. As the first

101 step, the effect in the hit efficiency was evaluated as described¹³²
 102 in Sec. 3.

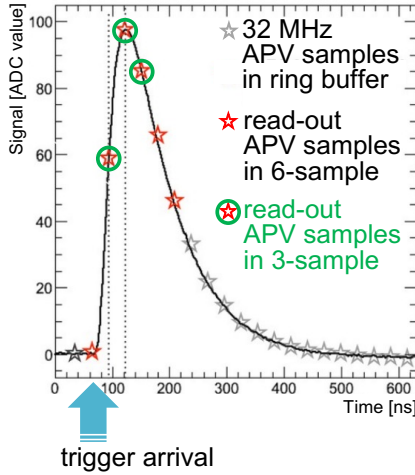


Figure 2: Sampling in the “multi-peak” mode of APV25. The black line shows the signal waveform after the CR-RC shaper circuit. The stars show the sampled signal height recorded in the analog ring buffer according to the 32 MHz sampling clock. The red stars indicate the six successive samples read out at the trigger reception in the six-samples mode. The red stars with a green circle indicate the samples read out in the three-samples acquisition.

103 The APV25 chips are mounted on each middle sensor (chip-
 104 on-sensor concept) with thermal isolation foam in between.
 105 The merit of this concept is shorter signal propagation length,
 106 leading to smaller capacitance of the signal line and hence re-
 107 duced noise level. To minimize the material budget the APV25
 108 chips on the sensor are thinned down to 100 μm . APV25s are
 109 mounted on a single side of the sensor and readout of the signals
 110 is from the other side via wrapped flexible printed circuits. The
 111 power consumption of the APV25 chip is 0.4 W/chip and in total¹⁴⁶
 112 700 W in the entire SVD. The chips are chilled by bi-phase¹⁴⁷
 113 -20°C CO_2 .

114 3. Performance

115 The SVD was combined with the PXD to complete the VXD
 116 assembly in October 2018, and the VXD was installed to the
 117 Belle II detector system in November 2018. Since March 2019,
 118 the SVD has been operating reliably and smoothly for two and
 119 a half years, without any major problems. The total fraction of
 120 masked strips is about 1%. There was only one issue where one
 121 APV25 chip (out of 1,748 chips) was disabled during the spring
 122 of 2019, which was remediated by reconnecting a cable in the
 123 summer of 2019.

124 The SVD has also demonstrated stable and excellent perfor-
 125 mance [6]. The hit efficiency is continuously over 99% in most
 126 of the sensors. The cluster charge distributions are also reason-
 127 able. On the u/p-side, the most probable values agree with the
 128 calculated charge amount induced by MIPs within the uncer-
 129 tainty in calibration. On the v/n-side, 10–30% of the collected
 130 charge is lost compared to MIP due to the smaller inter-strip
 131 capacitance of the floating strips with larger strip pitches than

the u/p-side. The most probable values of the cluster signal-to-
 132 noise ratio distributions range from 13 to 30.

134 We measured the cluster position resolution by analyzing the
 135 $e^+e^- \rightarrow \mu^+\mu^-$ data [7]. The cluster position resolution is esti-
 136 mated from the residual between the cluster position and the
 137 track position not biased by the target cluster after subtracting
 138 the effect of the track extrapolation error. The cluster position
 139 resolutions for different incident angles are shown in Fig. 3. For
 140 normal incident tracks, it agrees with the expectations from the
 141 strip pitch including floating strips. For tracks with an incident
 142 angle, it is expected to get a better resolution, which is indeed
 143 the case in the v/n-side results. However, this effect is not ob-
 144 served on the u/p-side, and the study is still ongoing to improve
 145 the cluster position estimation.

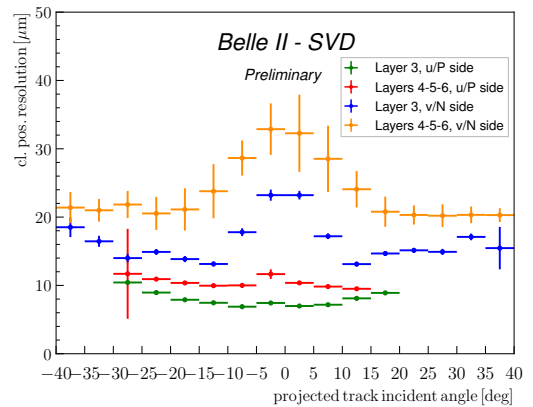


Figure 3: The SVD cluster position resolution depending on the projected track incident angle. The green (blue) plot shows the resolution in the u/p-side (n/v-side) of layer 3 sensors, and the red (yellow) one shows the u/p-side (n/v-side) of layers 4, 5, and 6 sensors.

The cluster hit-time resolution was also evaluated in hadron-
 148 event¹ data using the reference event time estimated by the Cen-
 149 tral Drift Chamber (CDC) outside of the SVD. The error on the
 150 event time, about 0.7 ns, was subtracted to evaluate the intrinsic
 151 SVD hit-time resolution. The resulting resolution is 2.9 ns on
 152 the u/p-side and 2.4 ns on the v/n-side. With such precise hit-
 153 time information, it is possible to reject off-time background
 154 hits efficiently. The hit-time distributions for signal² and off-
 155 time background³ are shown in Fig. 4. The signal distribution
 156 has a narrow peak, while the background hit-time distribution is
 157 broad and almost flat in the signal peak region. The separation
 158 power of the hit-time is high, as expected. For example, if we
 159 reject hits with the hit-time less than -38 ns in this plot, we can
 160 reject 46% of the background hits while keeping 99% of the
 161 signal hits. The background rejection based on the hit-time is
 162 essential to sustain the good tracking performance in the future
 163 high beam background condition.

The performance in three-samples data was compared with
 164 that in six-samples data to evaluate the performance in the

¹The event with more than three good tracks and not like Bhabha scattering.

²The clusters found to be used in the tracks in the hadron events.

³The clusters in events triggered by delayed-Bhabha pseudo-random trigger.

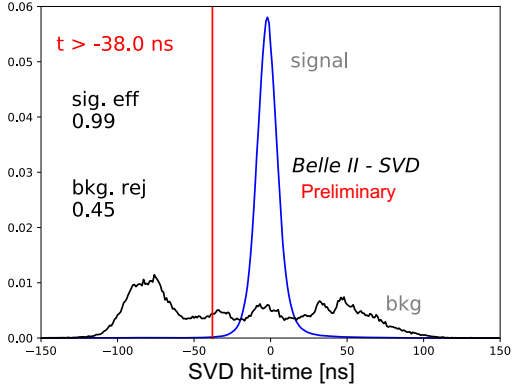


Figure 4: Example of the background hit rejection using hit-time. The blue distribution shows the signal, and the black distribution shows the off-time background. Assuming the hit-time cut at -38 ns, the signal hit efficiency of 99% and the background hit rejection of 46% are achieved.

mixed-mode. If the trigger timing has no deviation, the three-samples data will show comparable performance to the six-samples data because the relevant part of the signal waveform to evaluate the necessary signal properties, which are the signal height and the signal timing, can be accommodated in the three-sample's time window. However, when the trigger has a jitter and the timing shift happens, some part of the signal waveform can be out of the three-sample's time window, and the reconstruction performance deteriorates. We examined the effect on the hit efficiency as a function of the trigger timing shift. The effect is evaluated by the relative hit efficiency, which is defined as the ratio of the hit efficiency in the three-samples data to the one in the six-samples data. For this study, the three-samples data are emulated in the offline analysis from the six-samples data by selecting consecutive three samples at fixed positions in the six samples. The trigger timing shift is evaluated by the CDC event time. The resulting relative efficiencies as a function of the trigger timing shift in the hadron-event data are shown in Fig. 5. The decreasing trend is observed for the shift of the trigger timing, as expected. As a result, the relative efficiency is over 99.9% for the trigger timing shift within ± 30 ns.

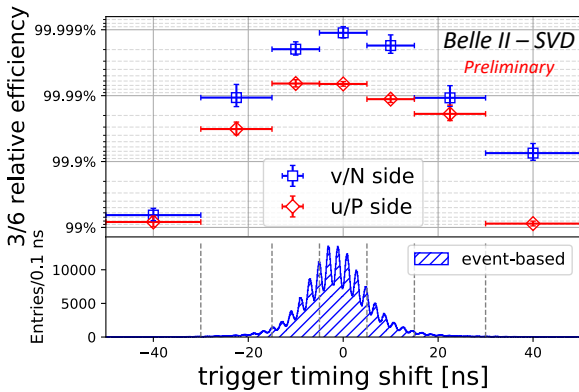


Figure 5: The relative hit efficiencies as a function of the trigger timing shift for v/n-side (blue square) and u/p-side (red diamond). The positive (negative) trigger timing shift corresponds to early (late) trigger timing.

4. Beam-related background effects on SVD

The beam-related background increases the hit occupancy of the SVD, which in turn degrades the tracking performance. Considering this performance degradation, we set the occupancy limit in layer 3 sensors to be about 3%, which will be loosened roughly by a factor of two after we apply the hit-time rejection described in Sec. 3. With the current luminosity, the average hit occupancy in layer 3 sensors is less than 0.5%. However, the projection of the hit occupancy at the luminosity of $8 \times 10^{35} \text{ cm}^{-2}\text{s}^{-1}$ is about 3% in layer 3 sensors. The projected occupancy comes from the Monte Carlo (MC) simulation scaled by the data/MC ratio determined from the beam background data of the current beam optics. The corresponding dose is about 0.2 Mrad/smy, and the equivalent 1-MeV neutron fluence is about $5 \times 10^{11} \text{ n}_{\text{eq}}/\text{cm}^2/\text{smy}$ (smy: Snowmass Year = 10^7 sec). The long-term extrapolation of the beam background is affected by large uncertainties from the optimization of collimator settings in MC and the future evolution of the beam injection background, which is not simulated. This uncertainty motivates the VXD upgrade which improves the tolerance of the hit rates and the radiation damage, and the technology assessment is ongoing for multiple sensor options.

From the measured dose on diamond sensors, the integrated radiation dose in the layer 3 mid-plane sensors, which are the most exposed in the SVD, is estimated to be 70 krad. The estimation is based on the correlation between the SVD occupancy and the diamonds dose. The estimated dose includes uncertainties of about 30% due to the unavailability of the appropriate trigger before December 2020. Assuming the dose/ n_{eq} fluence ratio of $2.3 \times 10^9 \text{ n}_{\text{eq}}/\text{cm}^2/\text{krad}$ from MC, 1-MeV equivalent neutron fluence is evaluated to be about $1.6 \times 10^{11} \text{ n}_{\text{eq}}/\text{cm}^2$ in the first two and a half years.

The effect of the integrated dose on the sensor leakage current is measured, and the results show a clear linear correlation as in the upper plot of Fig. 6. The slopes for all the sensors are summarized in the lower plot of Fig. 6. They are around 2–5 $\mu\text{A}/\text{cm}^2/\text{Mrad}$. The large variations can be explained by temperature effects and the deviation of sensor-by-sensor dose from the average in each layer used in the estimation. The slopes are in the same order of magnitude as previously measured in the BaBar experiment [8], $1 \mu\text{A}/\text{cm}^2/\text{Mrad}$ at 20°C . While the leakage current is increasing, the impact on the strip noise is suppressed by the short shaping time (50 ns) in APV25. It is expected to be comparable to the strip-capacitive noise only after 10 Mrad irradiation and not problematic for ten years where the integrated dose is estimated to be 2 Mrad.

The relation between the noise and the integrated dose is shown in Fig. 7. The noise increase of 20–25% is observed in layer 3, but this does not affect the performance of SVD. This noise increase is likely due to the radiation effects on the sensor surface. Fixed oxide charges on sensor surface increase non-linearly, enlarging inter-strip capacitance. The noise saturation is observed on the v/n-side and also starts to be seen on the u/p-side. This behavior agrees with the increase of fixed oxide charges.

The full depletion voltage of the sensor is also a key property

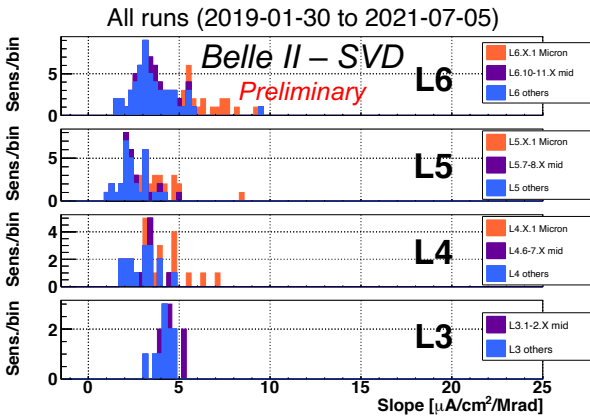
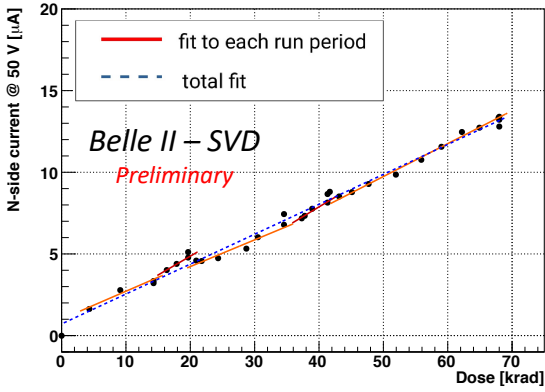


Figure 6: (upper) Effect of the integrated dose on the leakage current in the n/v-side of one sensor in layer 3. The slope is fitted for each run period (solid red line) and for all the runs (dashed blue line). (lower) The fit results of all the sensors for all runs. The sensors are classified as trapezoidal sensors in the forward region (Micron), sensors around the midplane, and the others.

that can be affected by the radiation damage. It can be measured from the v/n-side strip noise, which suddenly decreases at the full depletion voltage because the sensor substrate is n-type and thus the v/n-side strips can be fully isolated at the full depletion. From this measurement, reasonable full depletion voltages, which are consistent with the values mentioned in Sec. 2, were confirmed, and so far no change in full depletion voltage is observed in the first two and a half years of operation, which is consistent with the expectation from low integrated neutron fluence of $1.6 \times 10^{11} \text{ n}_{\text{eq}}/\text{cm}^2$.

5. Conclusions

The SVD has been taking data in Belle II since March 2019 smoothly and reliably. The detector performance is excellent and agrees with expectations. We are ready to cope with the increased background during higher luminosity running by rejecting the off-time background hits using hit-time and operating in the three/six-mixed acquisition mode. In the recent study, the efficiency loss in the three-samples data is confirmed to be less than 0.1% for the trigger timing shift within ± 30 ns. The observed first effects of radiation damage are also within expectation and do not affect the detector performance.

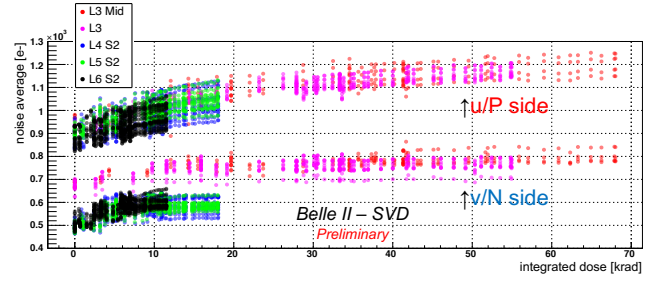


Figure 7: Effect of the integrated dose on the noise average in electron. The clear increase is observed and saturated (or start to be saturated) for layer 3 sensors.

Acknowledgments

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